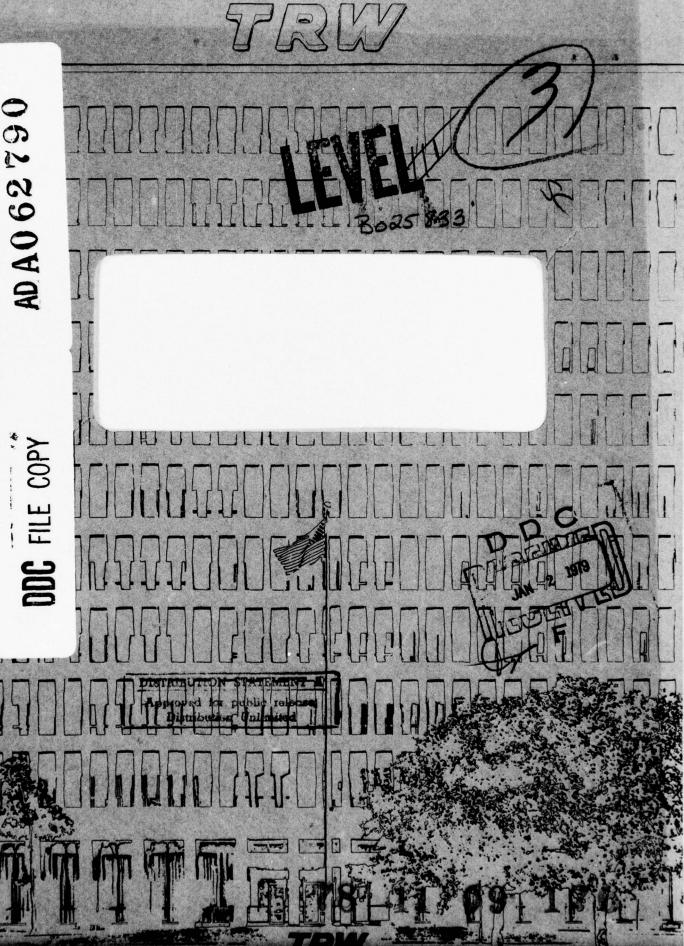


AD AO 62 790



TRW- 78.4734.7-023- PT-2-VOL-2 062 A062 PROCESS-ORIENTED, HIGH-INJECTION CIRCUIT MODELS FOR INTEGRATED BIPOLAR JUNCTION TRANSISTORS. PART II. - VOLUME II. THERE JAN 2 1979

SOLD JAN 2 1979 PRINCIPAL INVESTIGATOR Dr. John/Choma, Jr Electronic Systems & Technology Department

January 1978

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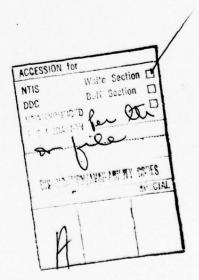
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### APPENDIX A

# EQSOLVE (EQUATION SOLVER) PROGRAM

- A.O EQSOLVE Program Documentation
- A.1 Introduction Program Overview
- A.2 Detailed Main Program Structure
- A.3 Program Calculations
- A.4 Detailed subroutine structure GRTM

### A.O EQSOLVE PROGRAM DOCUMENTATION

This appendix describes a software program called EQSOLVE, which was developed at TRW Defense and Space Systems Group for the ONR in conjunction with this contract. EQSOLVE accepts parameters of the Choma Bipolar Transistor model and a tabular list of base current values, and produces device performance curves such as Beta vs. Ic. The performance curves are presented in families, allowing parametric variation with a third parameter which may be a model parameter or the collector-emitter voltage. This allows a sensitivity analysis of transistor operating characteristics with respect to model parameters and allows a "fine tuning" of the model based on easily observed device performance curves. The EQSOLVE code is made to accept model data from a companion code called PAREV. PAREV generates from electrical and processing data the model parameters used as input to EQSOLVE.

#### A.1 INTRODUCTION - PROGRAM OVERVIEW

### A.1.1 Objectives

The EQSOLVE program is a dual-purpose code. In the early stages of contract performance, the code provided insight into development of test methodology for specific model parameters. By examination of Calcomp plot families of, for example, Beta vs. I with extrinsic emitter resistance R as a parameter, the effect of model parameters on test-generated curves is made evident, and insight is gained into testing methods for various model parameters. In the later stages of contract performance, after parameter testing methodology has been established, the EQSOLVE is code is used to fine tune and to verify the performance of a given device model prior to model release for CAD activity. An additional future goal for the EQSOLVE is in the area of computer-assisted device design, i.e., the process of using an ideal set of performance characteristics to direct the device fabrication. EQSOLVE provides a link allowing software comparison of ideal and predicted device performance in an iterative improvement feedback loop that can produce masking and diffusion information for device fabrication based upon desired performance characteristics.

## A.1.2 Program Structure

The EQSOLVE program is written in FORTRAN IV for use on the CDC 6000 or CYBER series machines. The core requirements are less than 60K (octal) to insure interactivity capability. Expected interactive terminal time for a single graph output is less than five minutes with a total cost of less than \$5.

A multiple graph output capability exists, allowing a user to generate several performance curve families on which a different model parameter is varied for every family. Additionally, different performance curves for a given model parameter (e.g.,  $f_T$  vs.  $I_c$  with  $R_c$  varying and  $\beta$  vs.  $I_c$  with  $R_c$  varying) can be generated. Each family request generates a list of common device values in addition to the Calcomp plot, namely:

- 1. I<sub>B</sub> Base current
- 2. I<sub>C</sub> Collector current
- 3.  $h_{\rm FE}$  DC current gain
- 4. B<sub>F</sub> Base pushout factor
- 5.  $f_T$  Gain bandwidth product
- 6.  $V_{
  m BE}$  Base-emitter voltage

Program input-output interaction and user interaction are summarized in figure A-1. User interaction consists of real time requests, wherein the details and contents of the Calcomp families of curves are stipulated.

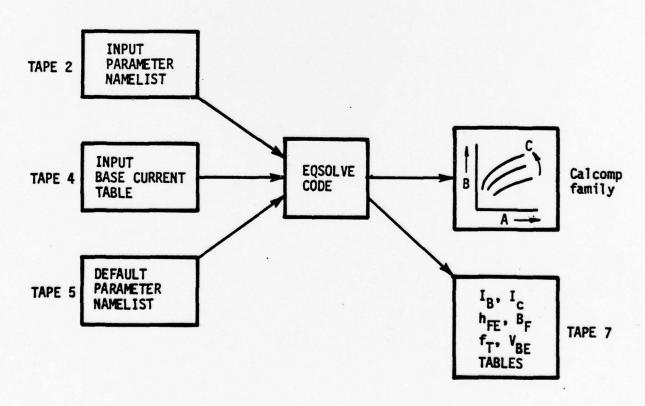


Figure A-1 EQSOLVE Input-Output Interaction

#### A.2 DETAILED MAIN PROGRAM STRUCTURE

### A.2.1 Namelist Input

The namelist input works in a default override fashion. Default parameters corresponding to the model of a 2 Gigahertz small geometry bipolar transistor are used on input TAPE5. The program reads the default values and then reads in any user-supplied update values on input TAPE2. The TAPE2 values replace the default values for the duration of the user session. If the PAREV companion code is used prior to EQSOLVE, TAPE2 is available with a complete set of updated parameters, and it totally overwrites the TAPE5 default set. This is the expected mode of usage. An example TAPE2 for PAREV output and EQSOLVE input is shown in figure A-2. Detailed parameter explanation is found in Section 2 of this report, and a list of input parameters with default values is supplied in table A-1.

### A.2.2 Base Current File Input

Base current file TAPE4 contains the list of base current points at which the bipolar device performance is to be calculated. The user thus has flexibility in selecting the range and coarseness of Calcomp output by tailoring TAPE4 to his own specifications. Input format is free-format as follows:

N < integer value = number of base current points >

IB1 First base current point in amperes

IB2 Second base current point (IB2 > IB1)

IBN Nth base current point

An example of a TAPE4 file is given in Figure A-3. Base current points are from .luA to 1.0mA.

### [ LIST, TAPE2

```
1STRANS
RBO
         = 1.8,
         = 1.5E+02,
RBB
         = 5.0E-01,
RE
         = 0.,
RCC
BS
         = 3.33333333E-01,
BC
         = 3.3333333E-01,
BE
         = 5.0E-01,
CSO
         = 2.0E-12,
CEO
         = 8.998858595E-13,
CCO
         = 7.999348551E-13,
PHISS
         = 9.0E-01,
 PHICO
         = 9.004004474E-01,
PHIEO
         = 1.200517562,
         = 3.0E-12,
ISR
IS
         = 3.2E-16,
IER
         = 3.0E-14,
 IBR
         = 3.0E-13,
VER
         = 2.196833685E+01,
VCR
         = 6.056999515E+01,
         = 1.8,
 NC
VT
         = 2.587500647E-02,
VTS
         = 3.0E-02,
 QFN
         = 6.178E-16,
 GRN
         = 2.261411812E-12,
         = 1.275024046E-10,
 TAUFO
         = 4.667132469E-07,
 TAUR
 BETAO
         = 5.5E+01,
 BETAR
         = 1.5,
 AE
         = 9.8E-07,
 XO
         = 2.7E-04,
 WCPRIME = 3.0E-04,
         = 1.5,
 NE
         = 5.0E-01,
 KR
 BF2
         = 1.4E+01,
         = 1.368120028E+01,
 VB
 VF
         = 2.587500647E-02,
 FF
         = 1.027800526E-14,
 VCE
         = 5.0,
         = 1.380000345E-03,
 IBB
         = 3.25E-03,
 PC
 SEND
```

Figure A-2 PAREV Output - EQSOLVE Input TAPE2

TABLE A-1 INPUT PARAMETER LIST - SECMENT 1 OF 3

2

4

Program Name	Symbol.	Description	TAPE 5 Default	Units
RHOB	p <sub>B</sub>	Extrinsic Base Resistivity	600.	ohms
TE	t E	Emitter Depth	$1 \times 10^{-4}$	СШ
AB	A <sub>B</sub>	Base Area	$5 \times 10^{-7}$	cm <sup>2</sup>
RBO	R <sub>BO</sub>	Extrinsic Base Resistance		ohms
1	1	Emitter Width	$6 \times 10^{-5}$	cm
Ħ	,c	Emitter Length	6 x 10 <sup>-5</sup>	CB
3	3	Base Width	$6 \times 10^{-5}$	E C
RBB	RBB	Intrinsic Base Resistance	0.	ohms
RE	ᄯ	Extrinsic Emitter Resistance	0.5	ohms
RHOC	o <sub>o</sub>	Collector Resistivity	.00325	ohm-cm
WCPRIME	W,	Distance from Collector-Base Junction to Buried Layer	3 × 10 <sup>-4</sup>	8
OX	v°	$x_0/w_C^2$ is a parametric ratio	2.7 × 10 <sup>-4</sup>	E
AE	A E	Area of Emitter	9.8 × 10 <sup>-7</sup>	cm <sup>2</sup>
RCC	R <sub>CC</sub>	Low Injection Extrinsic Collector Response	7.0	ohms

TABLE A-1 INPUT PARAMETER LIST - SEGMENT 2 OF 3

Program Name	Symbol T	Description	TAPE 5 Default	اب ی	Units
	v TS	Substrate Junction Voltage Parameter	.03		volts
	ISR	Substrate Saturation Current	3 x 10 <sup>-12</sup>	-15	amps
	os <sub>o</sub>	Substrate Zero Bias Capacitance	$2 \times 10^{-12}$	-15	farads
	\$S\$	Substrate Built-In Junction Voltage	6.0		volts
	og <sub>o</sub>	Emitter Junction Zero-Bias Capacitance	9 x 10 <sup>-13</sup>	-13	farads
	φ E0	Emitter Junction Built-In Junction Voltage	٠.		volts
_	တ္ဗ	Collector Junction Zero-Bias Capacitance	8 x 10 <sup>-13</sup>	-13	farads
	00	Collector Junction Built-In Junction Voltage	1.2		volts
	LER	Nonideal Emitter Diode Satur- ation Current	3.2 x 10 <sup>-14</sup> amps	10_14	amps
	E E	Nonideal Emitter Exponent Multiplier	1.5		

TABLE A-1 INPUT PARAMETER LIST — SEGMENT 3 OF 3

17 .1E-6 . 2E-6 .4E-6 .7E-6 1.0E-6 2.0E-6 4.0E-6 7.0E-6 10.E-6 20.E-6 40.E-6 70.5-6 .1E-3 . 2E-3 .4E-3 .7E-3 1.E-3

Figure A-3 TAPE4 - Base Current Input File Example

# A.2.3 Tabular Output

Local file TAPE7 is written at every  $I_B$  point with the variables  $V_C$ ,  $I_B$ ,  $I_C$ ,  $h_{FE}$ ,  $B_F$ ,  $f_T$ , and  $V_{BE}$ . TAPE7 may be listed at the termination of an interactive session to supplement the Calcomp printer plots. An example TAPE7 partial listing is shown in Figure A-4.

### A.2.4 Calcomp Output

Examples of Calcomp output graphs are shown in the sensitivity curve sets of Appendix B.

### A.2.5 User Interaction - Sample Session

The interactive session using EQSOLVE requires user input in response to terminal prompters for specification of Calcomp output details. The first request identifies the independent (x-axis), dependent (y-axis), and parametric variables by name and index number. Index number is the relative position of the variables within a COMMON block called PRAMTRS, and is used for internal program identification only. Parameter name is used for axis labeling and may be specified by alphanumeric string less than eight characters without altering graphical output. Index values are tabulated, including default values, in Figure A-5.

#### [ LIST, TAPE7

VCPRIME = -4.514E+00 IB = 1.000E-06 IC = 2.491E-05 HFE = 2.491E+01 BF = 1.009E+00 FT = 7.939E+07 VBE = 6.485E-01

Figure A-4 TAPE7 Partial Output List

Index	Parameter	Default	Index	Parameter	Default
1	RBO	9.	35	NDC	1.E16
2	RBB	150	36	PHIC	.7
3	RE	.5	37	DNC	48.
4	RCO	5.	38	XO	1.2E-4
5	PC	.00325	39	KR	.4
5	BS	.33333	40	TD	20.E-12
7	BC	.33333	41	M	2.0
8	BE	.5 .	42	NE	1.5
9	CSO	2E-12	43	IB	
10	CEO	.9E-12	44	VCPRIME	
11	CCO	.8E-12	45	IBB ·	
12	PHISS	.9	46	QF	
13	PHICO	.9	47	QR	
14	PHIEO	1.2	48	LAMBDAE	
15	ISR	3.E-12	49	ISS	
16	IS	3.2E-16	50	IN	
17	IER	3.2E-14	51	II	
18	IBR	3E-13	52	IO	
19	VER	20.	53	ILIM	
20	VCR	60.	54	BF	
21	NC	1.8	55	QBN	
22	VT	.0259	56	VBE	
23	VTS	.03	57	IC	
24	QFN	1.5E-16	58	HFE	
25	QRN	1.5E-12	59	FT	
26	TAUFO	159.E-12	60	VCE	5.0
27	TAUR	1270E-12	61	BETAE	
28	BETAO	55.	62	BF2	13.64
29	BETAR	1.5	63	VB	17.793
30	W	.6E-4	64	VF	2.59E-3
31	AE	9.8E-7	65	FF	1.336E-14
32	VLIM	6.E6	66	F	
33	DNB	22.6	67	BFO	
34	WCPRIME	3.E-4			

Figure A-5 Common Parameter Indices

A typical prompter-response dialogue is as follows:

ENTER, IDNME, INDEX, DPNME, INDEX, PARNME, INDEX

? | \*VBE\*, 56, \*IC\*, 57, \*RBO\*, 1 | user supplied

The dialogue requested an IC versus VBE plot with resistor RBO as the varying parameter.

The next dialogue requests Calcomp type, number of curves per plot, and the varying parameter values. Curve type is specified by the following index list:

X-Axis	Y-Axis	Index
LINEAR	LINEAR	1
LINEAR	LOG	2
LOG	LINEAR	3
LOG	LOG	4

A prompter-response dialogue requesting a linear  $(x) - \log (Y)$  plot with three values of a varying parameter equal to 90, 9, and .9 (RBO values) is:  $(NCURVES \le 5)$ 

ENTER, TYPE, NCURVES, PARAM VALUES

?[2,3,90.,9.,.9] user supplied

The third and final dialogue requests adjustments in size of the x-axis and y-axis, and number of points between sumbols on a curve. A sample dialogue requesting a 6 inch x-axis, 7 inch y-axis, and plot symbols every 3 points (default values) is:

UPDATE XINCH, YINCH, MARKS (5) IF REQD SCURV

YINCH=7,XINCH=6,MARKS=3,3,3\$|+ user supplied

Note that three entries of three (3) are required since three curves are requested by NCURVES prompter.

An actual teletype session for Calcomp output using the above dialogues is given in Figure A-6. The default x-axis, y-axis size and symbol density are used by supplying a terminating dollar sign (\$). An actual Calcomp output for this session is given in Figure A-7.

07/13/77. 10.37.10. TRW/TSS C174 - 07/11/77.

USER NUMBER: JOB NAME PLRG. PORT 220. [ GET, TAPE4, EQSOLVE, TAPE5 = EQDATA1 [ RX, I=EQSOLVE [ LGO ENTER INDNNE, INDEX, DPNNE, INDEX, PARNNE, INDEX ? \*VBE\*,56,\*IC\*,57,\*RBO\*,1 ENTER TYPE, NCURVES, PARAN VALUES 7 2,3,90,.9,9 UPDATE YINCH, XINCH, HARKS (5) IF REQD SCURV TYPE 1 FOR ANOTHER CASE ENTER INDNNE, INDEX, DPNNE, INDEX, PARNNE, INDEX (etc.)

Figure A-6 Interactive Session Printout

# -- RBO = 9.000E+01 + -- RBO = 9.000E-01 0 -- RBO = 9.000E+00

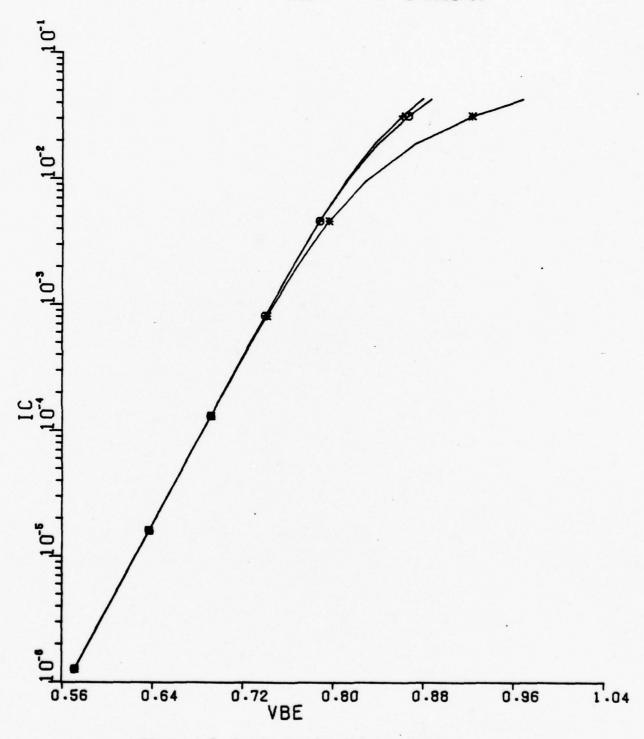


Figure A-7 Sample Session Calcomp Output

## A.3 Program Calculations

The EQSOLVE program performs calculations of internal and external parameters of a given modeled transistor at a set of base current values and at a fixed VCE in the normal (forward) regime. Due to the dependencies of the model equations, a nonlinear root finder routine is required to calculate internal currents and voltages, and an iteration path is used to converge the internal currents and junction voltages to values which satisfy both the model equations and the external constraints of a fixed base current IB and collector-emitter voltage VCE. An overview of the program flow is shown in Figure A-8, illustrating the internal variables Z, V<sub>E</sub> and I<sub>C</sub> generated by three separate calls to a root finder routine called GRTM. An auxiliary subroutine is used in each call, which describes the controlling equation containing the variable to be found. Details of the GRTM and auxiliary subroutines are explained in other sections of this document.

The set of calculations begins with a requested current point,  $\mathbf{I}_{B}$ , and model parameters (user supplied or default). An initial guess of collector junction intrinsic voltage is given as

$$V_{C} = -V_{CE} \tag{A-1}$$

and GRTM is called to evaluate Z from the equation

$$ZTan(Z) = I_{R}/I_{RR}$$
 (A-2)

Given Z,

$$R_{B} = \frac{R_{BB}[Tan(Z) - Z]}{4 \cdot Z \cdot Tan(Z)}$$
 (A-3)

and

$$A_{EF} = A_{E} \cdot \sin(2Z)/2Z \tag{A-4}$$

The estimate  $V_C$  is now required to compute  $V_E$  from

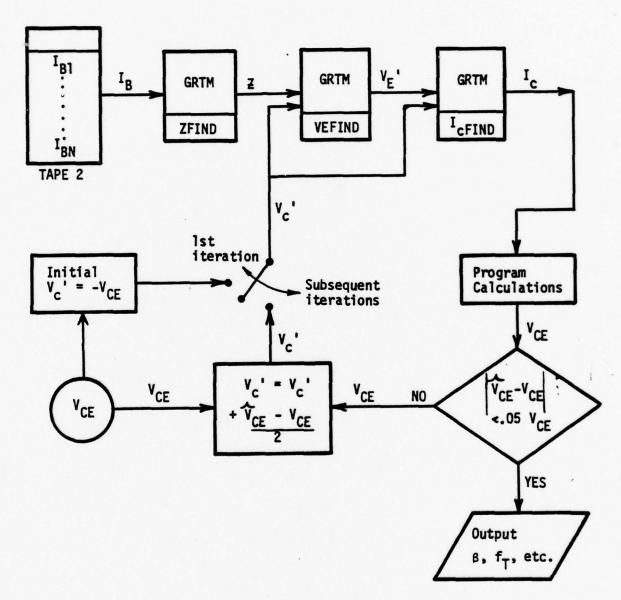


Figure A-8 EQSOLVE Program Flow Diagram

$$I_{B} = I_{ER} \left[ \exp(V_{E}^{2}/N_{E}V_{T}^{2}) - 1 \right] + \frac{I_{S}}{\beta_{o}} \left[ \exp(V_{E}^{2}/V_{T}^{2}) - 1 \right] + \frac{I_{S}}{\beta_{r}} \left[ \exp(V_{C}^{2}/V_{T}^{2}) - 1 \right] + I_{BR} \left[ \exp(V_{C}^{2}/N_{C}V_{T}^{2}) - 1 \right]$$
(A-5)

by another call to the root finder routine. The intrinsic junction voltages are now known, and the current and charge terms are calculated. Prior to this, if  $R_{\rm CC}$  is not given as input, the calculation

$$R_{CC} = \frac{\rho_C W_C'}{A_{EF}}$$
 (A-6)

is performed. Then,

$$I_{N} = I_{S} \left[ \exp(V_{E}^{\prime}/V_{T}) - 1 \right] \tag{A-7}$$

$$I_{I} = I_{S}[\exp(V_{C}/V_{T}) - 1]$$
 (A-8)

$$Q_F = Q_{FN} \left[ \exp(V_F'/V_T) - 1 \right] \tag{A-9}$$

$$Q_{R} = Q_{RN} \left[ \exp \left( V_{C} / V_{T} \right) - 1 \right] \tag{A-10}$$

$$I_{SS} = I_{SR} [exp(V_C/V_{TS}) - 1]$$
 (A-11)

The next block of calculations is Early effect and base pushout parameters.

$$\lambda_{E} = 1 + V_{E}/V_{ER} + V_{C}/V_{CR}$$
, (A-12)

$$F = FF[\exp(V_F^2/V_F) - 1], \qquad (A-13)$$

$$B_{FO} = 1 + B_{F2} [\exp(V_C/V_R)],$$
 (A-14)

$$B_{F} = B_{FO} + (1 - B_{FO})[(\exp(-F)]]$$
 (A-15)

The third call to the root finder routine is made to evaluate  $I_{\hat{C}}$  from the equations

$$Q_{BN} = \frac{\lambda_E}{2} + \left[ \left( \frac{\lambda_E}{2} \right)^2 + B_F Q_F + Q_R \right]^{1/2}$$
 (A-16)

$$I_{C} = I_{SS} - I_{BR} \left[ \exp(V_{C}^{\prime}/N_{C}V_{T}^{\prime}) - 1 \right] - \frac{I_{S}}{\beta_{r}} \left[ \exp(V_{C}^{\prime}/V_{T}^{\prime}) - 1 \right] + \frac{I_{N} - I_{I}}{Q_{DN}}$$
(A-17)

Additionally, the program evaluates

$$\lambda_{\rm CC} = \frac{X_{\rm o}}{W_{\rm C}'} \left(\frac{B_{\rm F} - 1}{B_{\rm F2}}\right)^{1/2}$$
 (A-18)

$$R_{C} = R_{CC} (1 - \lambda_{CC}) \tag{A-19}$$

$$I_{BE} = \frac{I_S}{\beta_0} \left[ \exp(V_E / V_T) - 1 \right]$$
 (A-20)

$$I_{BC} = \frac{I_S}{\beta_r} \left[ \exp(V_C / V_T) - 1 \right] \tag{A-21}$$

$$V_{BE} = I_B \cdot R_{BO} + (I_{BE} + I_{BC}) \cdot R_B + V_E' + (I_B + I_C) \cdot R_E$$
 (A-22)

Finally,

$$\hat{v}_{CE} = v_{X} = i_{C} \cdot R_{CO} + (i_{C} - i_{SS}) \cdot R_{C} - v_{C}' + v_{E}' + (i_{B} + i_{C}) \cdot R_{E}$$
 (A-23)

The comparison of  $\hat{\mathbf{v}}_{\text{CE}}$  to  $\mathbf{v}_{\text{CE}}$  is now performed.

If  $|\hat{v}_{CE} - v_{CE}| \le .05 \ v_{CE}$ , the calculations proceed. Otherwise, a new estimate of  $v_C^2$  is generated by

$$\hat{v}_{C} = v_{C} + (\hat{v}_{CE} - v_{CE})/2$$
 (A-24)

and program control s transferred to the second call to GRTM following equation (A-5).

Once a "suitably close" value for  $\hat{V}_{CE}$  is found, the program calculates several parameters of interest at the operating point ( $I_B$ ,  $V_{CE}$ ). The first quantities are the base pushout and forward base charge sensitivities,

$$S_{BE} = \frac{(B_{FO} - 1) \cdot FF \cdot exp(V_E/V_F - F)}{B_F * V_F}$$
 (A-25)

$$S_{BC} = \frac{B_{F2} \cdot \exp(V_C^2/V_B) \cdot [1 - \exp(-F)]}{B_F \cdot V_B}$$
 (A-26)

$$S_{QE} = \frac{\left[I_{C} \cdot B_{F} \cdot Q_{FN} / (I_{S} \cdot V_{T})\right] (1 + V_{T} \cdot S_{BE})}{2 \cdot Q_{BN} - \lambda_{E}}$$
(A-27)

$$S_{QC} = \left(\frac{1}{2 \cdot V_{CR} \cdot Q_{BN}} + \frac{1}{2 \cdot V_{CR} / (2 \cdot Q_{BN} - \lambda_E)}\right) \left(\frac{\lambda_E}{Q_{BN}} + 2 \cdot V_{CR} \cdot B_F \cdot Q_{FN} \cdot I_C \cdot S_{BC} / I_S\right) (A-28)$$

The forward and reverse incremental transconductances are given by

$$gmf = (I_C/V_T) \cdot (1 - S_{OE}/V_T)$$
 (A-29)

$$gmr = I_C \cdot S_{OC} \tag{A-30}$$

The program now checks if base pushout is in evidence  $(B_{\overline{F}} > 1)$  and calculates

$$R_{\text{OC}} = \begin{cases} R_{\text{CC}} & B_{\text{F}} = 1 \\ R_{\text{C}} - \frac{I_{\text{C}} \cdot R_{\text{CC}} \cdot (X_{\text{O}} / W_{\text{C}}^{2}) \cdot (\frac{B_{\text{F}} \cdot S_{\text{BE}}}{\text{gmf}} - \frac{B_{\text{F}} \cdot S_{\text{BC}}}{\text{gmr}})}{2 \cdot B_{\text{F}} \cdot [(B_{\text{F}} - 1) / B_{\text{F2}}]^{1/2}} & B_{\text{F}} > 1 \end{cases}$$
(A-31)

and

1

$$R_{OB} = 2 \cdot_{VT} / (3 \cdot I_{BB}) \tag{A-32}$$

The junction capacitances are assigned values by the equations

$$C_{jE} = \begin{cases} \frac{C_{EO}}{(1 - V_E'/\phi_{EO})^{B_E}} & V_E' < \phi_{EO} \\ C_{EO} \cdot (1 + \frac{B_E \cdot V_E'}{\phi_{EO}}) & V_E' > \phi_{EO} \end{cases}$$
(A-33)

$$C_{jC} = \begin{cases} \frac{C_{CO}}{(1 - V_C^{\prime}/\phi_{CO})^{B_C}} & V_C^{\prime} < 0 \\ C_{CO} \cdot (1 + \frac{B_C \cdot V_C^{\prime}}{\phi_{CO}}) & V_C^{\prime} > 0 \end{cases}$$
(A-34)

Diffusion capacitances are

$$C_{DE} = B_F \cdot \tau_{FO} \cdot (gmf + I_C \cdot S_{BE})$$
 (A-35)

$$^{C}_{DC} = ^{B}_{F} \cdot ^{\tau}_{FO} \cdot (^{I}_{C} \cdot ^{S}_{BC} - gmr)$$
 (A-36)

yielding total capacitances given by

$$C_{TE} = C_{DE} + C_{1E}$$
 (A-37)

$$c_{TC} = c_{DC} + c_{1C} \tag{A-38}$$

Substrate voltage is assumed equal to collector junction voltage, and

$$c_{SS} = \begin{cases} \frac{c_{SO}}{(1 - v_C^{\prime}/\phi_{SS})^B S} & v_C^{\prime} < 0 \\ c_{SO} \cdot (1 + b_S \cdot v_C^{\prime}/\phi_{SS}) & v_C^{\prime} > 0 \end{cases}$$
 (A-39)

The time moment calculations proceed with

$$g_{BE} = \frac{I_{BE}}{V_{T}} + \frac{I_{ER}}{n_{E} \cdot V_{T}} \cdot \left[ \exp(V_{E} / n_{E} \cdot V_{T}) - 1 \right]$$
 (A-40)

$$F_E = 1 + (gmf + g_{BE}) \cdot R_E$$
 (A-41)

$$g_B = g_{BE}/F_E \tag{A-42}$$

$$C_{E} = C_{TE}/F_{E}$$
 (A-43)

$$go = gmr/F_{g}$$
 (A-44)

$$gm = (gmf - gmr)/F_E$$
 (A-45)

$$R_{L}^{\prime} = \frac{R_{OC} + R_{CO}}{1 + go \cdot (R_{OC} + R_{CO})}$$
 (A-46)

$$r_{11} = 1/g_B$$
 (A-47)

$$r_{22} = R_L^2 + (1 + gm \cdot R_L^2) \cdot r_{11}$$
 (A-48)

$$r_{33} = R_{0B} + r_{11} + (1 + gm \cdot r_{11}) \cdot R_L$$
 (A-49)

$$r_{44} = \frac{(1 + go \cdot R_{OC}) \cdot R_{CO}}{1 + go \cdot (R_{OC} + R_{CO})}$$
 (A-50)

and

$$T_1 = r_{11} \cdot c_{TE} + k_R \cdot r_{22} \cdot c_{TC} \cdot (1 - k_R) \cdot r_{33} \cdot c_{TC} + r_{44} \cdot c_{SS}.$$
 (A-51)

The final parameters of interest are calculated as

$$\beta_{F} = (\frac{gm}{g_{B}}) [1 + go(R_{OC} + R_{CO})]$$
 (A-52)

$$f_{T} = \frac{\beta_{F}}{2 \cdot \pi \cdot T_{1}} \tag{A-53}$$

$$h_{FE} = I_C/I_B \tag{A-54}$$

After output of  $I_B$ ,  $I_C$ ,  $h_{FE}$ ,  $B_F$ ,  $f_T$ ,  $V_{BE}$ , and any other user-requested values, the program control is returned to the beginning of the calculations, and the calculations are performed again with a new value of  $I_B$ .

## A.4 Detailed Subroutine Structure - GRTM

## A.4.1 Module Usage and Interface

GRTM is a general root finder routine which is called three times during a pass through the EQSOLVE equations. It is called to determine the onset parameter Z,  $V_{\rm E}$ , and  $I_{\rm C}$ . Linkage into GRTM is accomplished by providing auxiliary subroutines which contain an equation that evaluates to zero for the proper value of the variable to be found. The auxiliary subroutine names and equations used in EQSOLVE calls are as follows:

Name	Equation
ZFIND	(A-2)
VEFIND	(A-5)
ICFIND	(A-17)

## A.4.2 Subroutine Input

Input to GRTM is the auxiliary function name, an initial guess of roots requested, and an iteration convergence value described in section 4.1.4.4. The maximum iteration value is set at 50 and a convergence criteria of  $1 \times 10^{-6}$  is used for EQSOLVE.

### A.4.3 Subroutine Output

The onset parameter Z, intrinsic base-emitter voltage  $V_{\underline{E}}$ , and collector current  $I_{\underline{C}}$ , are outputs of GRTM when called from EQSOLVE.

## A.4.4 Subroutine Calculations

GRTM uses Muller's method to find the roots of f(x). This method is detailed in Section 4.3.2.

A.5 EQSOLVE CODE LISTING

+ TAPES-INPUT)	TAPES-INPUT)
PROGRAM WILL SOLV	PROGRAM WILL SOLVE A SERIES OF EQUATIONS TO DETERMINE THE EFFECTS
OF PROCESS PARAME	TER VARIATION ON SEMICONDUCTOR DEVICES.
CONNON /DATA/ FX	500,51, X(500,51, CGVAR(5), CGLBL, DPEND,
+ INDPEND, TYPE, NPTS,	PTS. NCURVES. YINCH, XINCH, MARKS(5)
COMMON /PRANTRS/	COMMON IPPAMTRS/ RBD. RBD. RE. RCO. PC. BS. BC. BE. CSO. CEO.
+ CCO, PHISS, PHIC	O. PHIED, ISR, IS. IER, IBR, VER, VCR, NC. VI.
+ VIS, OFN, ORN,	AUFO, TAUR, BETAD, BETAR, M. AE, VLIM, DNB,
WCPRIME, NOC. PI	MCPRIME, NOC. PHIC. ONC. XO, KR. TO, M. NE. 18, VCPRIME, 188,
THE THE LABORE	1550 IN 110 IOS ILINO BES OBNO VBES ICS HEES
ENSTON	TINK(67)
	/CURV/ YINCH, XINCH, MARKS
	, IS, IER, IBR, KR, P. NC, NE, IO, ILIM, NDC,
115 1556	THE BOLD KIN KED LMBDAED ICO LMBDACCO LASTED
. LASTVE, LINK	
EQUIVALENCE (PBD, LINK(1))	LINK(1))
NAMEL IST /TRANS/	
. CCO. PHISS, PHIC	CCO. PHISS, PHICO, PHIED, ISP, IS, IER, IBR, VER, VCR, NC, VI,
+ VIS, OFM, ORN, 1	VISS OFNS ORNS TAUFOS TAURS BETADS BETARS AES
+ MCPRIME, XO, NE,	MCPRIME, XO, NE, KR, BF2, VB, VF, FF, VCE, 188
FORMATI / / SX + VC	PRIME . * FELO.305X24IB . * FELO.305X24IC . *
+ E10.3s/s5xseHFE	SOMPOHIE & COELOSSOMPHE & CARLOSSOMPHE & CO
THIERED TYPE YOUR YOUR CONTRACTOR	FIG. 39289 9289 4 9FIG. 31
EXTERNAL ZFIND. VEFIND.	FERNO LOFTED
CALL NALEDF	
CALL PLOT (0., 1., 23)	
16 REWIND 2	
TE (FOF.2) 4.4	
S XINCH-6. S YINCH-	7. S HARKS(1) - MARKS(2) - MARKS(3) - 3
DISPLAY •EN	Y GENTER INDNME, INDEX, DPNME, INDEX, PARNME, INDEX+
ACCEPT IND	END, XPTR, OPEND, YPTR, CGLBL, CGPNTR
DISPLAY OFF	TENTER TYPE-NCURVES-PARAY VALUES+
OTCOL AV CLD	PERSONAL VINCEL

1.602E-19 1.602E-19 1.002E-19 1.002E-12 1.002E-19 1.002E		The state of the s		AND THE REAL PROPERTY OF THE P		2) + (TAN(Z)++2))							.0) .0) RINE + (18 + IC) + RE FE + VEPRINE +
	-	SEAPES . O NPASS . O RELIND 4	ACCEPT(4) NPTS GRAPHS + 1 LINK(GPNTR)-CGVAR(GRAPHS)	ACCEPT(4) 18 IF (EDF, 4) 7,0 VALUES - 1 AST7	Z	1 - 21) / ((4.0 +	ZE.	C + UCPRIME) / AEF + (EXP(VEPRIME/VI) - 1 + (EXP(VCPRIME/VI) - 1	GFN + (EXP(VEPRINE/VI) - 1 GRN + (EXP(VCPRINE/VI) - 1  • ISR + (EXP(VCPRINE/VIS) -	ME/VER + VC ME/VF) - 1.0 XP(VCPRIME/V	MALUES = 10 CALL GRIM11, VALUES, 0, 50, ICFIND, 0, 1, E-6)	1C * VALUES. LMBDACC *(XO/WCPRIME)*((8F-1.)/8F2)**.5 RC * RCC * (1.0 - LMBDACC)	- (15/8ETAD) + (15/8ETAR) + (16 + 16 + 16 + (16 + 16)

	1.0 + VCR + GBN)) + ((1.0 / (2.0 + VCR) / HBDAE)) + ((HBDAE/GBN) + (2.0 + VCR + BF + GFN   S)))	### 1.0 PROCERC-ICERCCE(XD/WCPRINE) # (BF * SBE/GMF-BF * SEC/GMR)/ #### 1.1	CJE  CJC  LE 0.01 CSS = CSG/((1.0 - VCPRIME/PHISS)**BS)  AT + IER/(NE + VI) + (EXP(VEPRIME/CNE + VI) - 1.0)  (GMF + GBE) + RE	JMF - GMR)/FE  - (ROC + RCD)/((11.0 + 60 + (RGC + RCD)))  - (AOF)  - (ROC + RCD)/((11.0 + 6M + RIPRIME) + R1]  - (B + R11 + (1.0 + 6M + R1)) + RLPRIME  - (B + R11 + (1.0 + 6M + R1)) + RLPRIME  - (B + R11 + (1.0 + 6M + R1)) + RPRIME  - (B + R0 + R0C) + RCD)/(1.0 + 60 + (ROC + RCD))  - (GM/GB)/(1.0 + 60 + (ROC + RCD))  - (GM/GB)/(1.0 + 60 + (ROC + RCD))  - (GM/GB)/(1.0 + 61 + 11)  - (GM/GB)/(1.0 + 61 + 11)
SCVX-VCE) LE	/ (2.0 + - LHBDA / IS)) 17 + (1.		CTE = CDE + CJE CTC = CDC + CJC IF(VCPRIME GE, 0.01 CSS = CSO 1F(VCPRIME GE, 0.01 CSS = CSO GBE = 18E/VT + 1ER/(ME + VT) + FE = 1.0 + (GMF + GBE) + RE GB = GBE/FE GD = GMR/FE	GN = (GNF - GNR)/FE  RIPRIME = (ROC + RCO)/((1).0 + GO + (RGC  R11 = 1.0/60  R22 = RLPFIME + (1).0 + GM + RIPRIME) + R  R33 = ROB + R11 + (1).0 + GM + R11) + RLP  R44 = ((1).0 + GO + ROC)/(1).0 + GO  T1 = R11 + CTE + KR + R22 + CTC + (1).0  + R44 + CSS  8ETAF = (GM/GB)/(1).0 + GO + (ROC + RCO)  FT = BETAF/(2.0 + P1 + T1)  HFE = IC/IB

-	WRITE (7.1) VCPRIME. IB. IC. HFE. BF. FT. VBE FX(NPASS.GRAPHS) - LINK(YPTR) X(NPASS.GRAPHS) - LINK(XPTR) IF (NPASS. LT. NPTS) GO TO 5 IF (GRAPHS. LT. VCURVES) GO TO 9 GALL PLEKEC DISPLAYPE 1 FOR ANCTHER CASE*	
	IF(III.Eq.1) GD TO 16 CALL PLOT(0,0,999) STOP END SUBROUTINE ZEIND(2,FZ)	
U	ZFIND CALULATES Z AS A FUNCTION OF THE TANGENT OF ITSELF. COMMON /PPAMTRS/ SKIP1(42), 18, SKIP2, 188, SKIP3(22) REAL 18, 188 FZ ~ Z + TAN(2) - (18 / 188) FEIUR N	
uu	COLTINE VEF IND CALCULA RIME. NOW 7PPANTR IP3(5), BET I IER, IS,	
	TERN2 = (IS / BFTAD) + (EXP(VEPRINE/VI) - 1.0)  TERN4 = 18	
<b>.</b>	ICFIND CALCULATES THE VALUE FIC WHICH IS DEPENDENT UPON IC. THE VALLE  IC AND FIC MUST BE EQUAL LPGN CONVERGENCE.  COMMON PRANTRS/ SKIP1(15), IS, SKIP2, IBR, SKIP3(2), NC, VI,  SKIP4(6), BETAR, M. SKIP3(2), DNB, LCPRIME, SKIP6(2), DNC,  SKIP7(6), VCPRIME, SKIP8, QF, QR, LMBDAE, IS, IN, II, IG, ILIH,  REA GBN, SKIP9(12)  REAL IC, IO, ILIM, IS, ISS, IBR, NC, II, LMBDAE, IN  TERM: BR + (FXP(VCPRIME/(NC+VI)) - 1.0)	

SUBROUTINE PLIFKEC PLIEKEC SELECTS THE ROLTINES TO BE TRANSFERRED TO SETUP DEPENDING UPON WHICH TYPE OF PLOT IS DESIRED. FITYPE . Eq. 4) CALL SETUPISCALG, LGAXS, SCALG, LGAXS, LGLIN, LDGTYPE) FITYPE . EG. 21 CALL SETUPISCALG, LGAXS, SCALE, AXIS, LGLIN, LGGTYPE) FITYPE . EQ. 1) CALL SETUPISCALE, AXIS, SCALE, AXIS, LINE, LOGTYPE) SUBROUTINE SETUP (\*SCALE, FAXIS, ASCALE, XAXIS, CURVES, LOGTYPE)
SETUP IS THE EXECUTIVE FOR GENENRATING THE PLOTS.
COMMON /DATA/ FX(2500), X(2500), CGVAR(5), CGLBL, DPEND,
INDPEND, TYPE, NPTS, NCURVES, YINCH, XINCH, HARKS(5)
DIMENSION TITLE(4), CHARS(5), CHRINOX(5)
INTEGER TYPE, BASE
DATA CHARS /4HP -->4H+ -->4HD -->4HS -->4HP --/
DATA CHARS /4HP -->4H+ -->4HD -->4HS -->4HP --/ COMMON /DATA/ SKIP1(50C8), TYPE, SKIP2(5) Integer type External scalf, axis, line, scalg, Lgaxs, Iglin TERN2 - (IS/BETAR) + (EXPLVCPRIME/VI) - 1.6)
TERN3 - (IN - II) / 08N
FIC - ISS - TERN1 - TERN2 + TERN3 - IC
RETURN IS LINEAR/LINEAR IS LOG/LINEAR(Y/X) IS LINEAR/LOG(Y/X) TYPE = 1 PLOT IS LINEAR/LINE TYPE = 3 PLOT IS LINEAR/LINEA CALL PLOT(8.5,0.5-3) T AXIS PREPARATION .DGTYPE . -2 DGTYPE . 0 LOGIYPE . 1 BASE - 1 000

GENERATE CURVES AND PLOT LEGEND

C. GENERATE CURVES AND PLOT LEGEND

C. LINEAR/LINEAR PLOT

TELLOGTYPE .GT. -21 GD TO 3

TELLOGTYPE .GT. -21 GD TO 4

TELLOGTYPE

, 1 i	1 1	1 1		1 1	
		!			
FINANCE INVALOBELTA MAXVAL AXENGTH)  TERMINES THE MAX VALUE OF THE INFORMATION JUST SUBMITTED  ALTING ROUTINE.  ALTING MAXVALL LGMIN  LOGIOCHINVAL)  1.0 + 10.0**(AXLNGTH + DELTA + LGMIN)					
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5					
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IX (NINVAL, DELTA, MAXVAL, AXLNGTH) ES THE MAX VALUE OF THE INFORMA ROUTINE. KVAL, LGMIN MINVAL) 10.0**(AXLNGTH * DELTA * LGMIN)					
1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1			il		
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A B					
12 E					1 1
NG KE					
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## APPENDIX B

## SENSITIVITY CURVE SETS

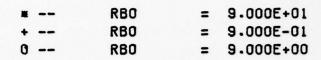
- B.O Introduction
- B.1  $I_B$  vs.  $V_{BE}$  Curves
- B.2 I<sub>C</sub> vs. V<sub>BE</sub> Curves
- B.3 H<sub>FE</sub> vs. I<sub>C</sub> Curves
- B.4 F<sub>t</sub> vs. I<sub>C</sub> Curves
- B.5 B<sub>F</sub> vs. I<sub>C</sub> Curves

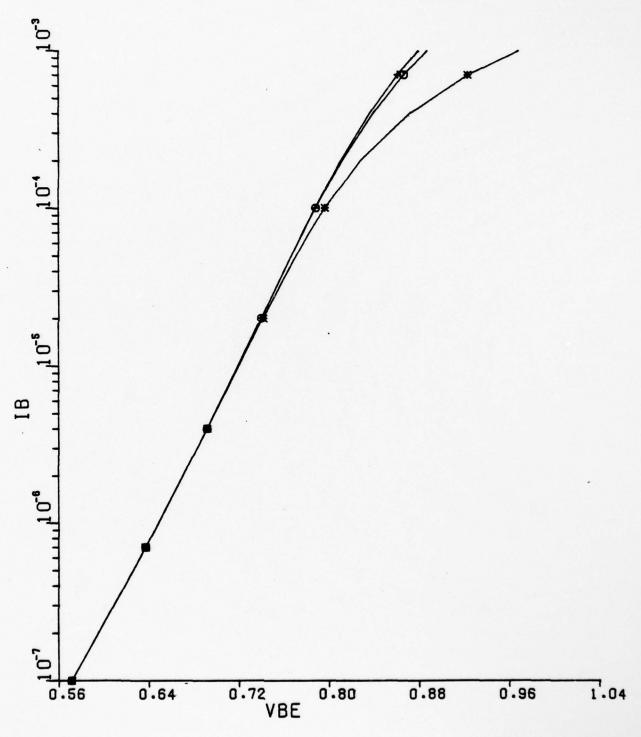
## B.O INTRODUCTION

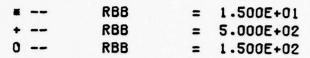
The sets of parametric curves comprising this appendix are Calcomp output from the EQSOLVE software code documented in Appendix A. The curve sets show the change in functional relationship of two-parameters of the Choma LSBJT model of a 2 GHz transistor under perturbation in value of a third parameter. The five sets of curves are selected, with the exception of the  $\rm B_F$  vs.  $\rm I_c$  set, as useful design-type curves which provide the essential device performance characteristics at a glance. Hence, the effects of the "third" parameter variation on these characteristic curves are produced by varying the third parameter above and below the nominal value. The  $\rm B_F$  vs.  $\rm I_c$  curves are included to provide insight into the effects of base pushout parameter values on the complex mathematical model of base pushout, an effect which is vital to understand for precise, large signal, high injection modeling.

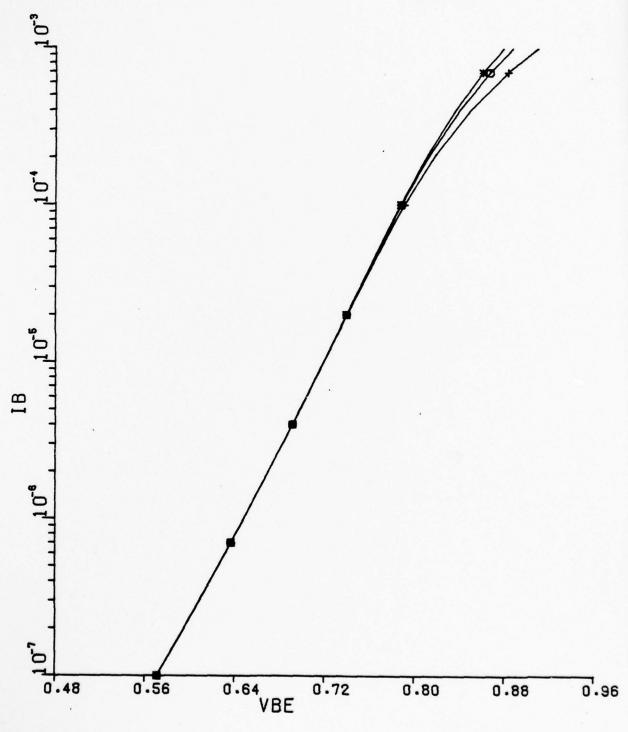
The design curves are included as a last step, or "fine-tuning" capability for a user who has dettermined a set of model parameters for a specific device. Comparison of device test data curves or other CAD values to mathematical performance curves generated by EQSOLVE or any general modeling software package (ECAP, SCEPTRE, SPICE, etc.) will produce curve differences similar to the curve differences on one or several of the curve sets contained in this appendix. The user then has insight on the erroneous parameter(s).

B.1  $I_B$  vs.  $V_{BE}$  Curves

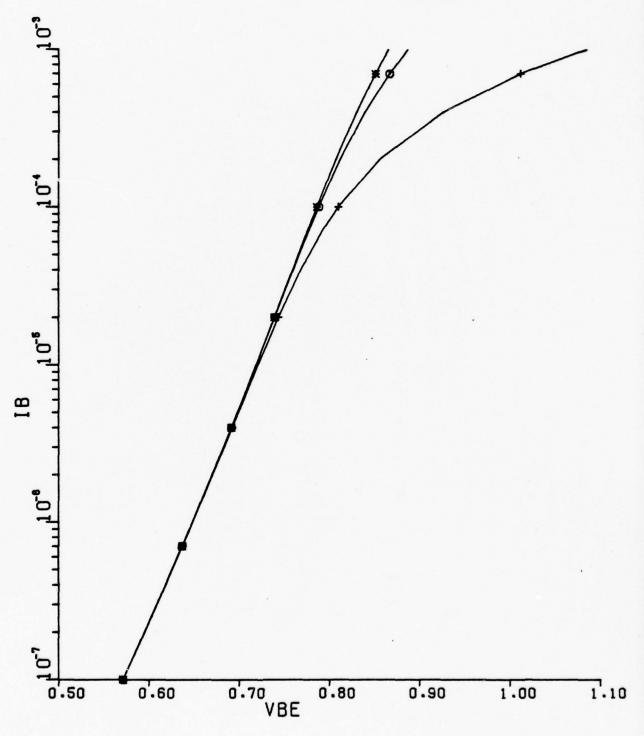


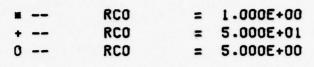


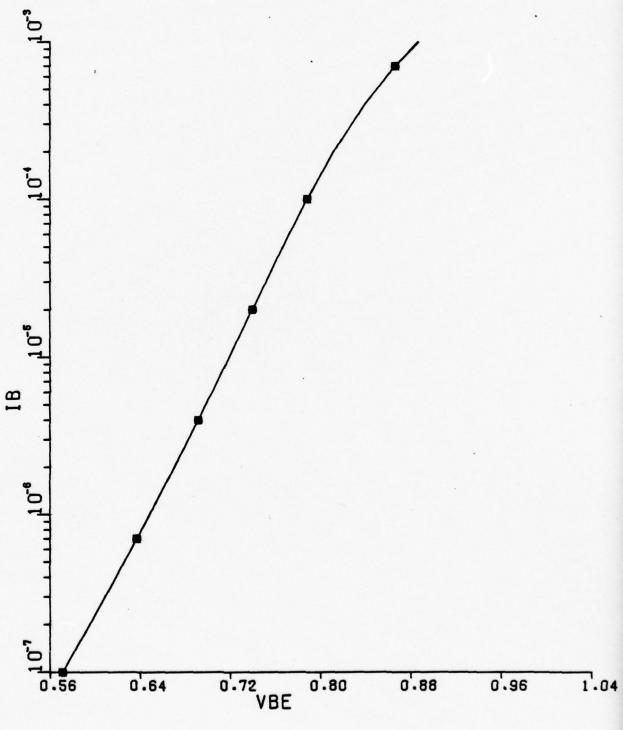




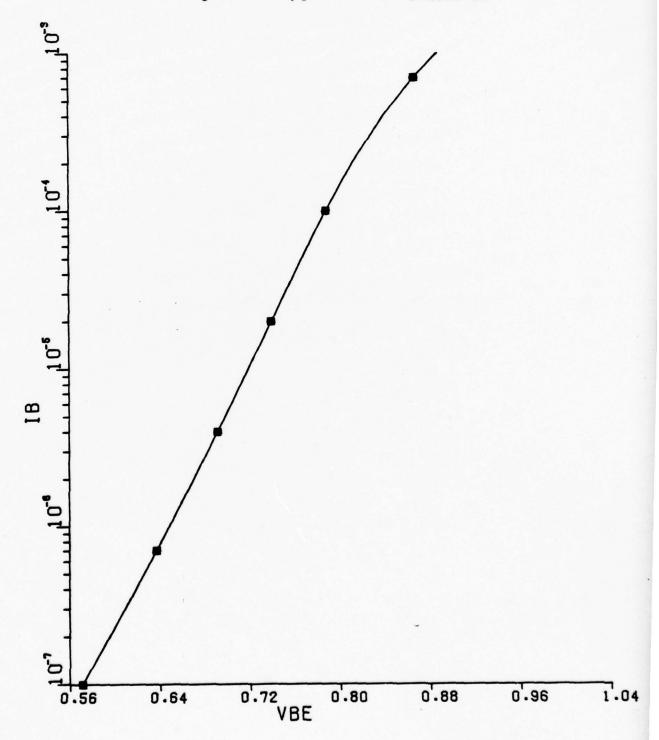




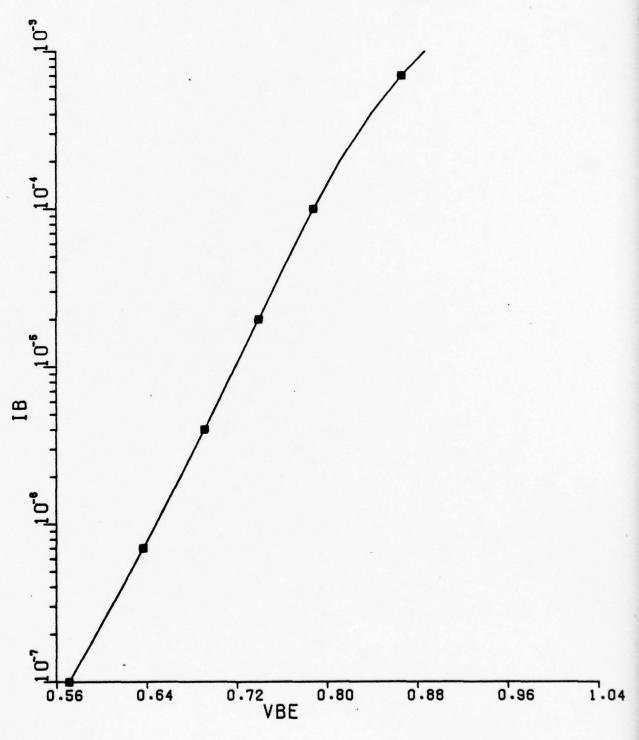




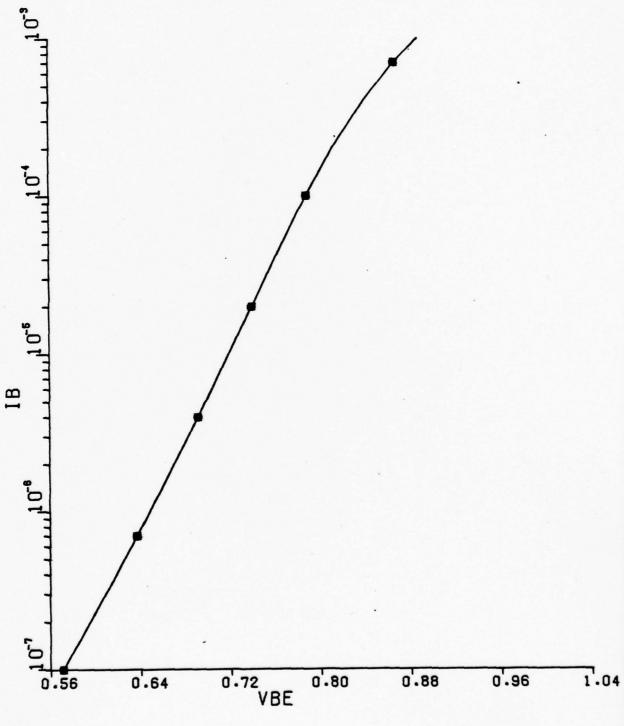
# -- PC = 6.000E-03 + -- PC = 1.000E-03 0 -- PC = 3.250E-03



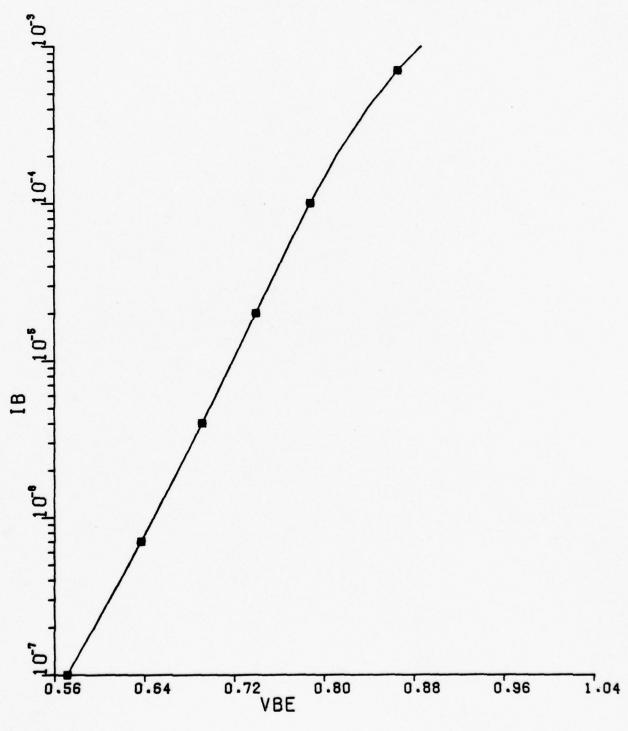


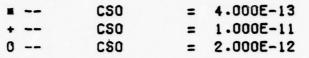


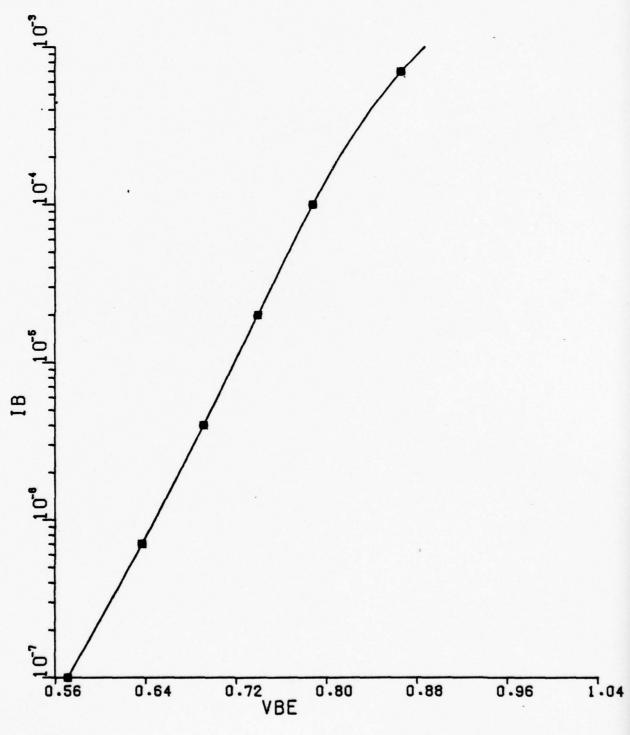




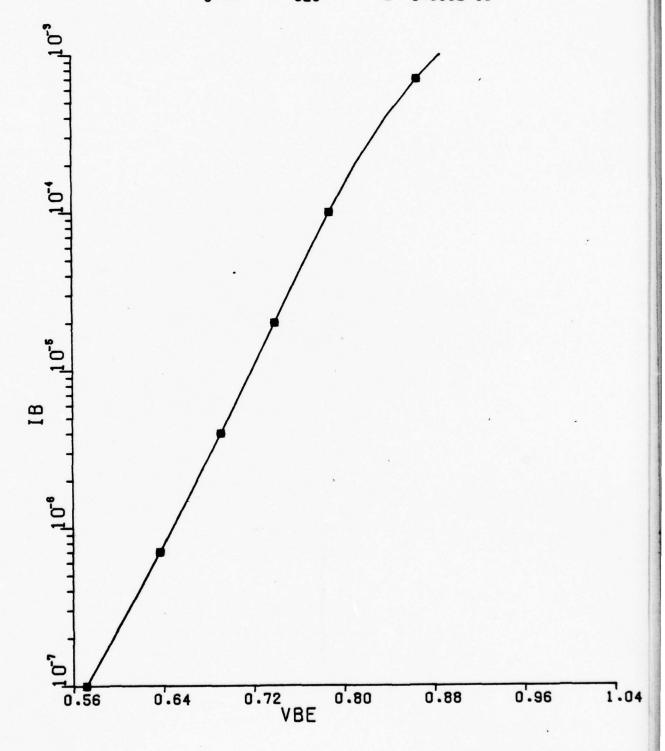




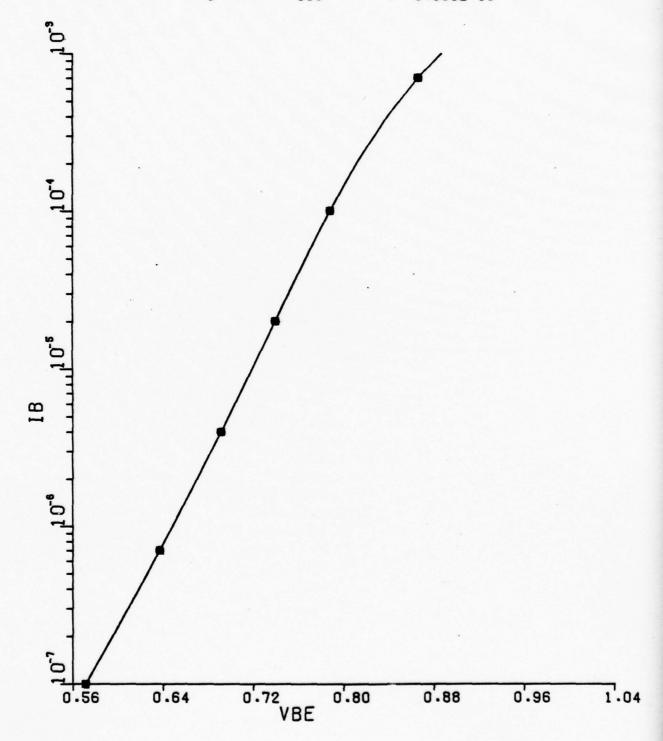




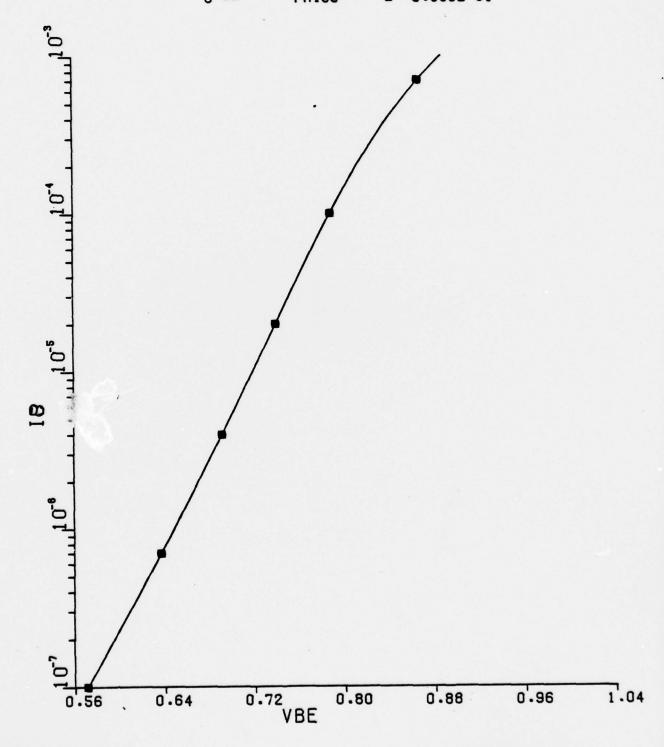
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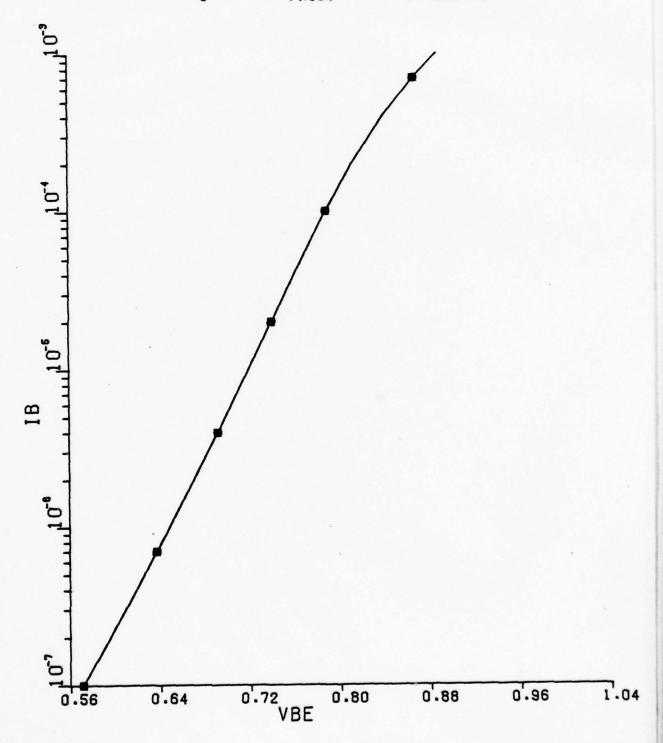
\* -- CCO = 4.000E-13 + -- CCO = 1.000E-10 0 -- CCO = 8.000E-13



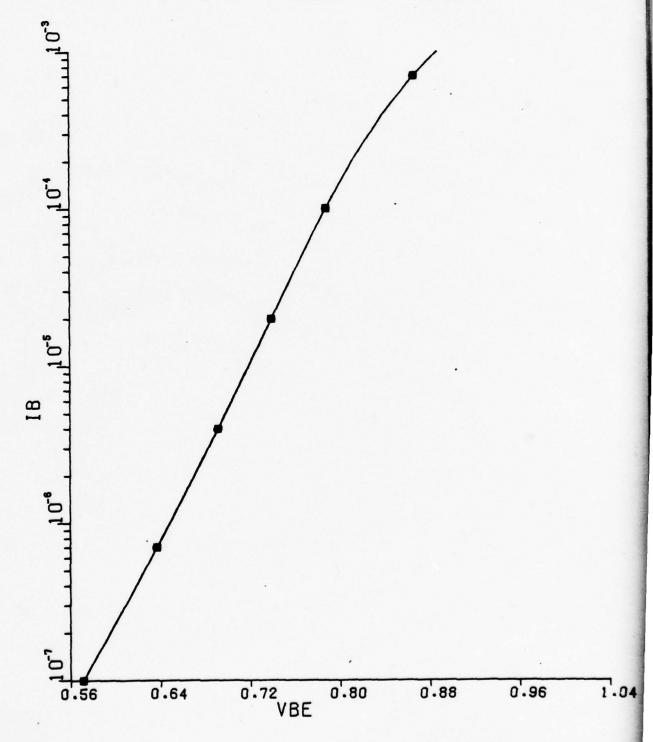
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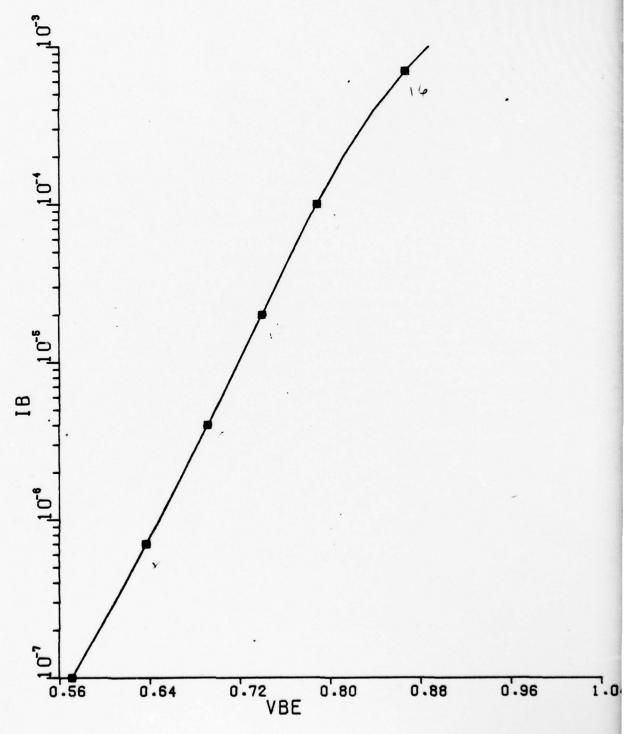
# -- PHICO = 5.000E-01 + -- PHICO = 1.200E+00 0 -- PHICO = 9.000E-01



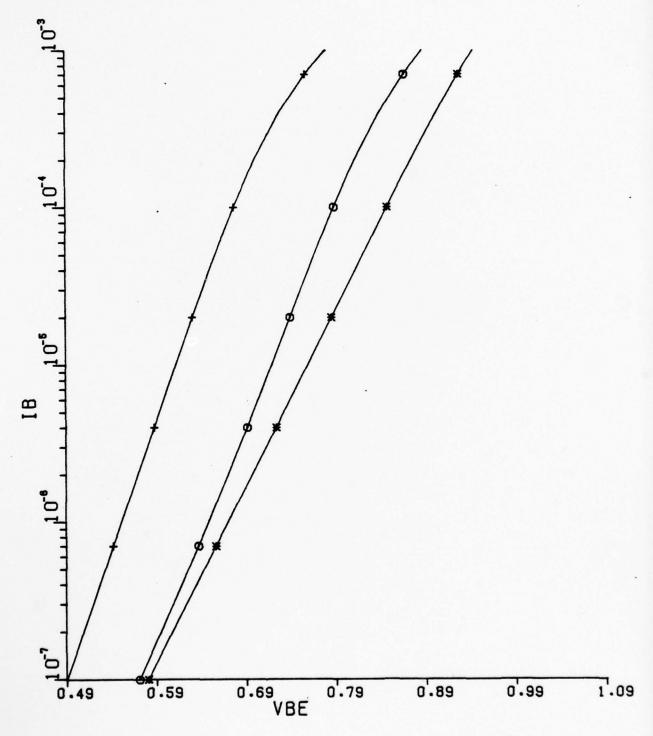
+ -- PHIEO = 5.000E-01 + -- PHIEO = 1.500E+00 0 -- PHIEO = 1.200E+00

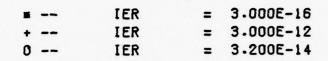


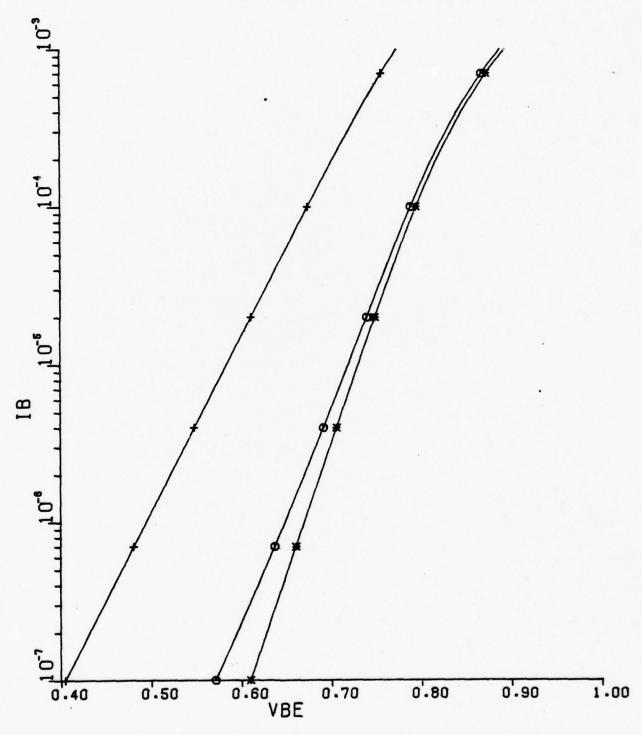


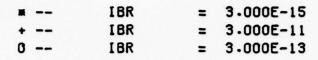


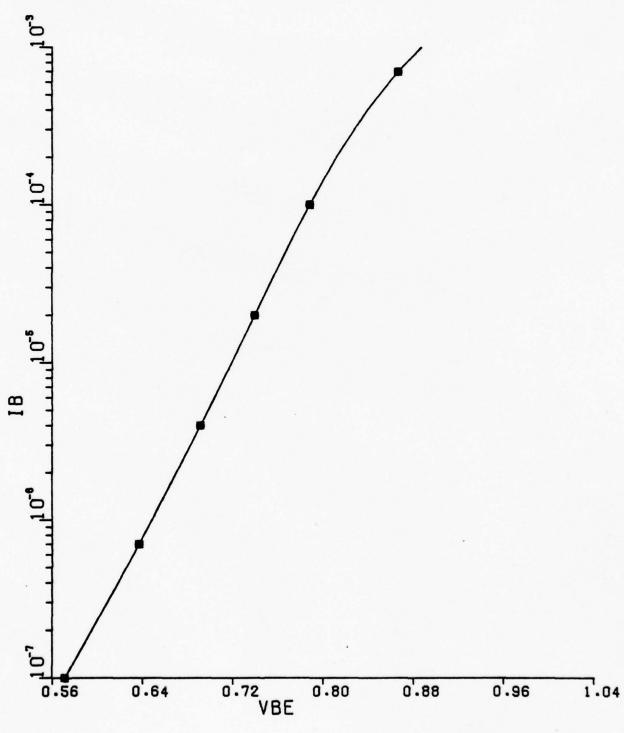


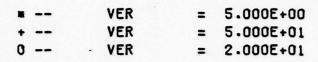


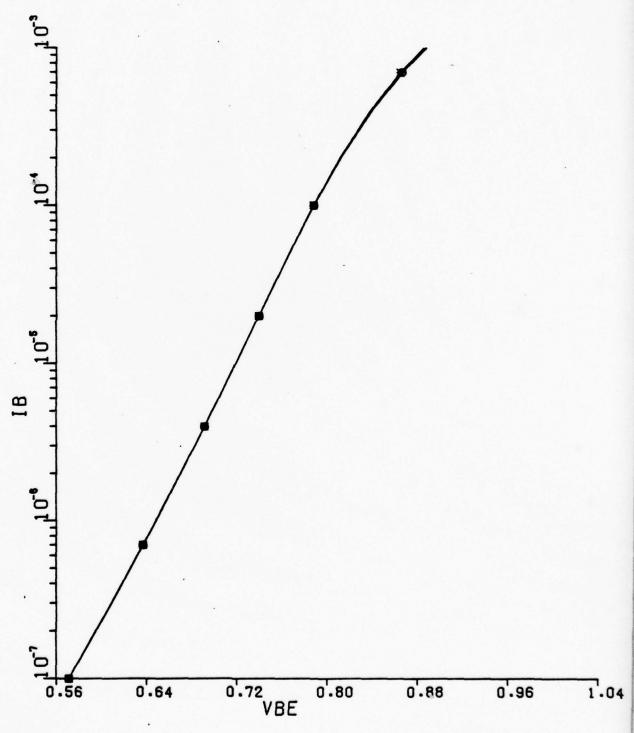


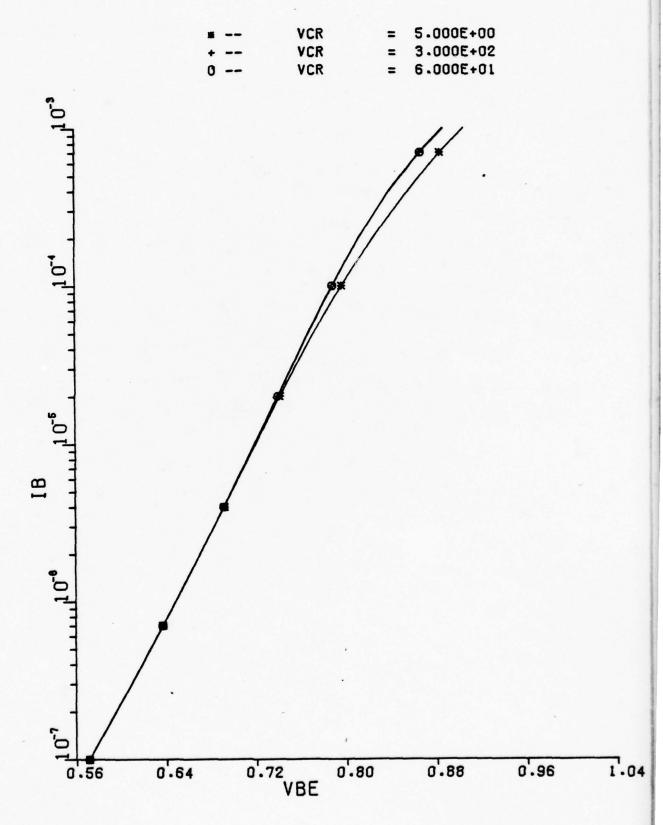




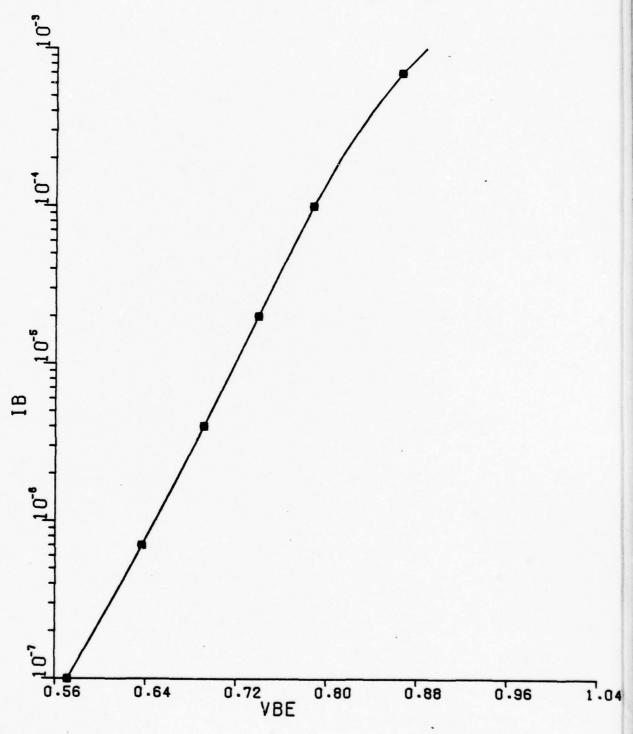




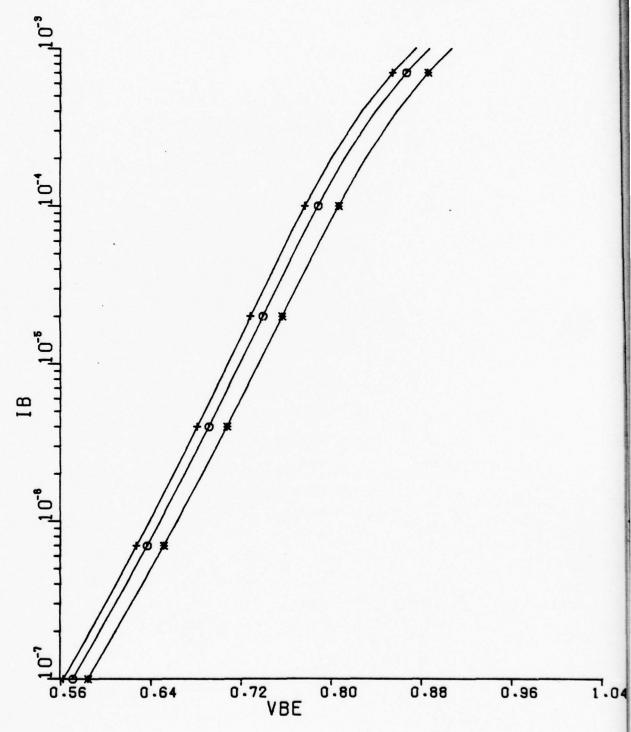


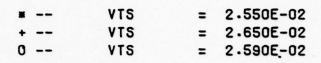


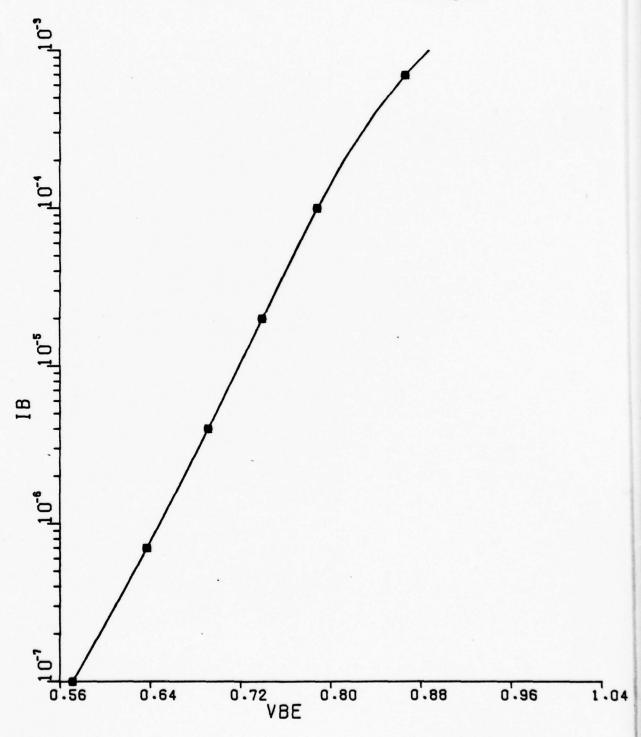




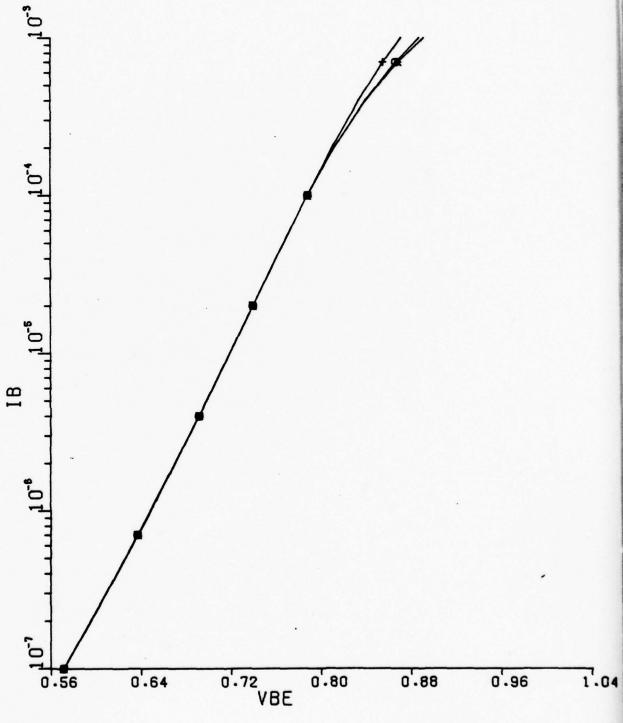




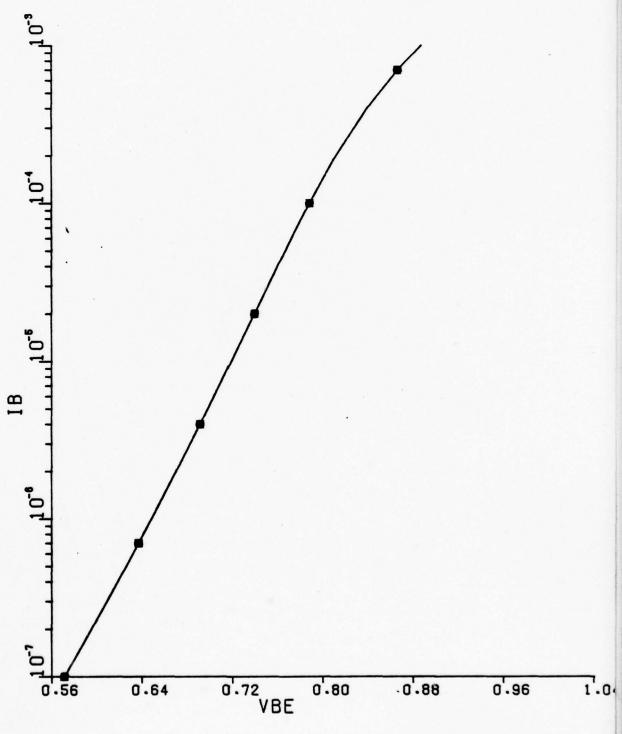




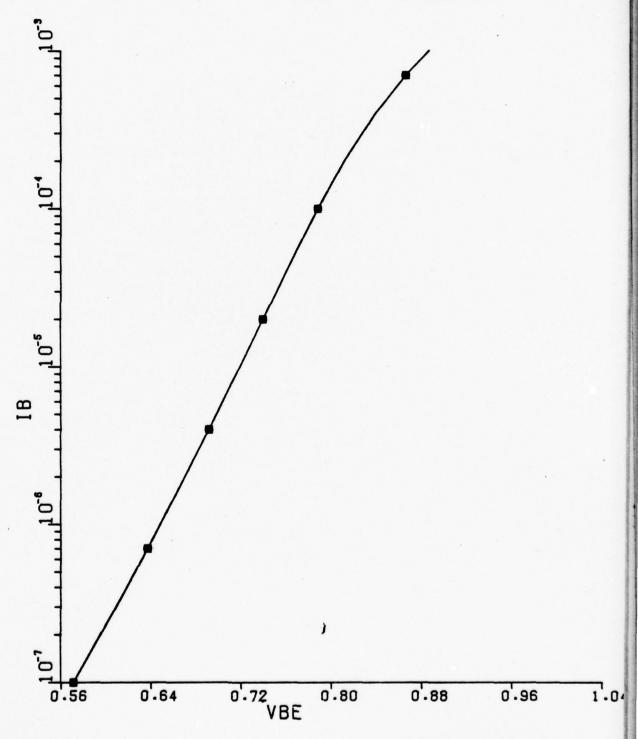




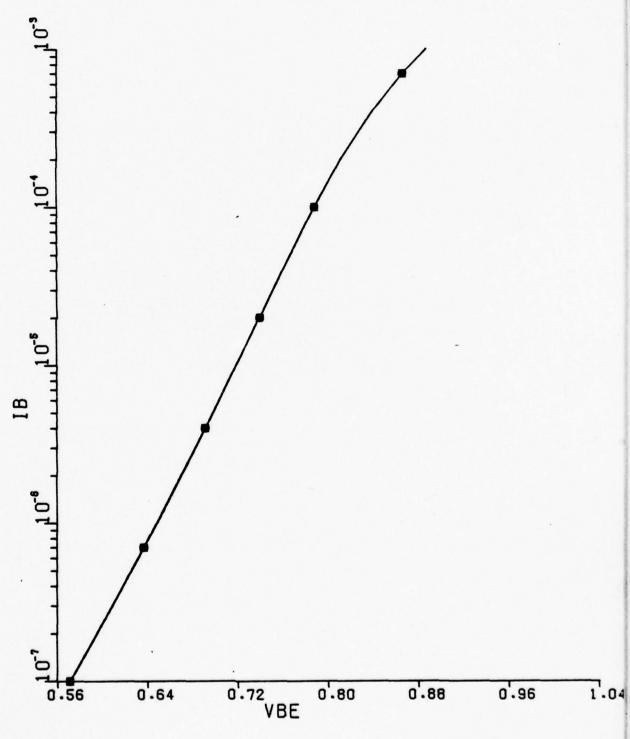




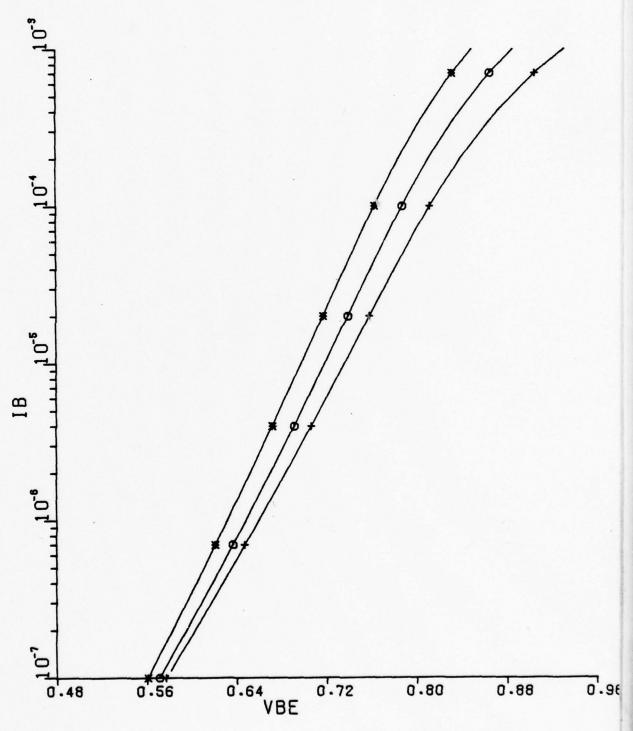




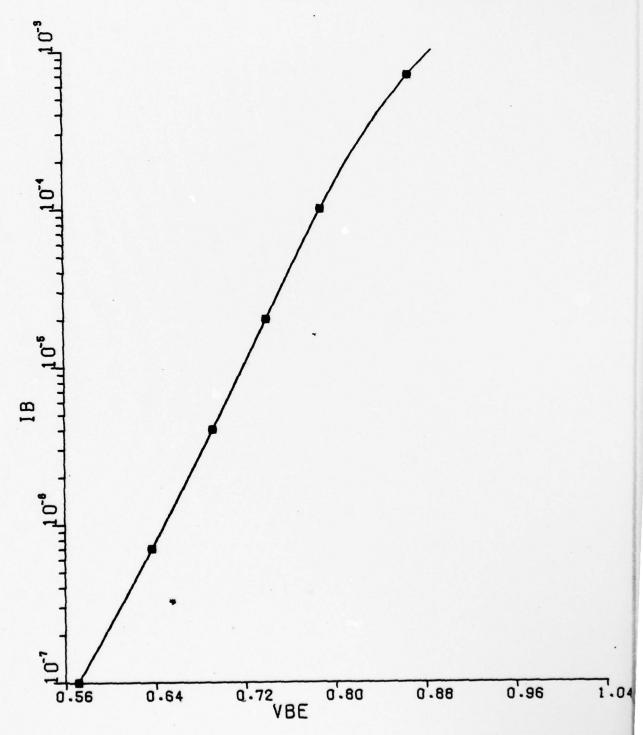




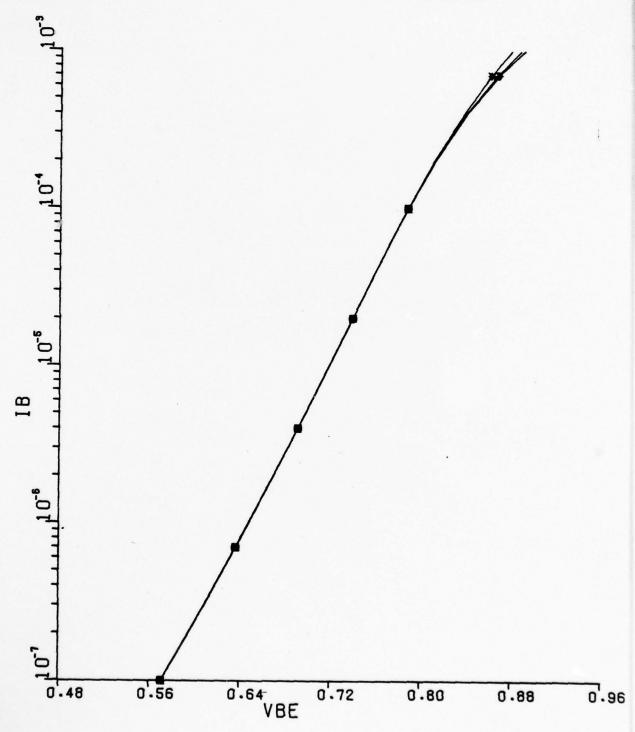


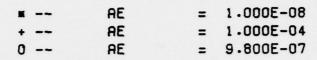


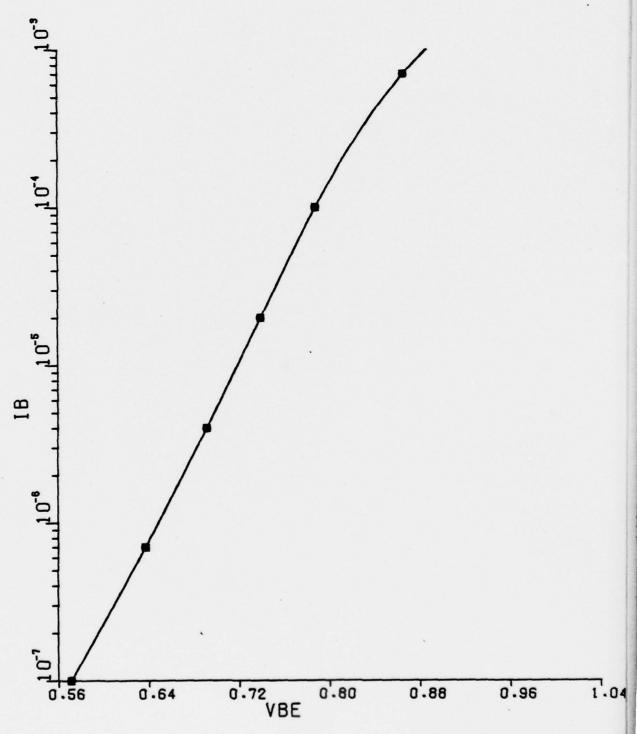




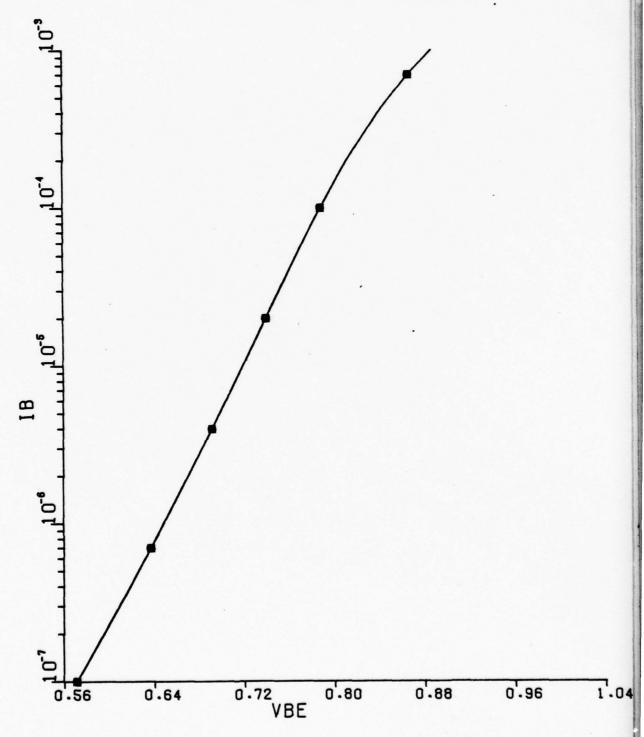






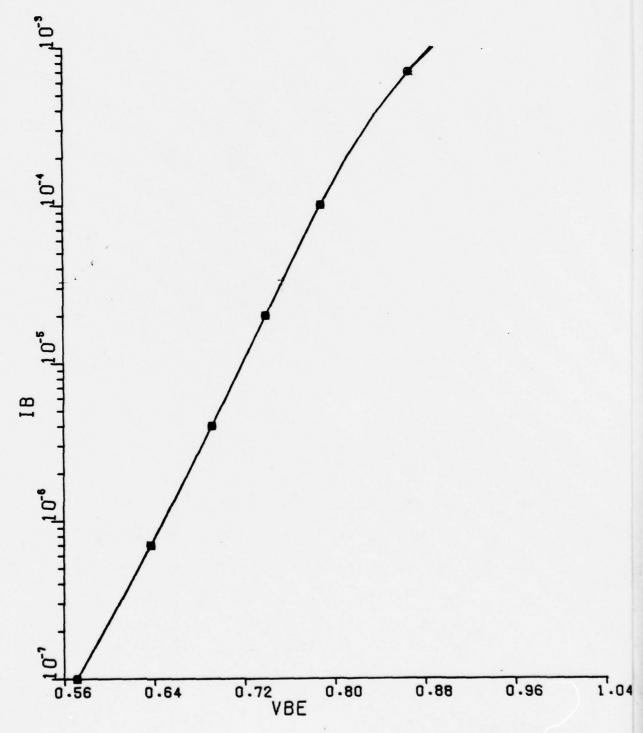


\* -- VLIM = 1.000E+06 + -- VLIM = 1.500E+07 0 -- VLIM = 6.000E+06

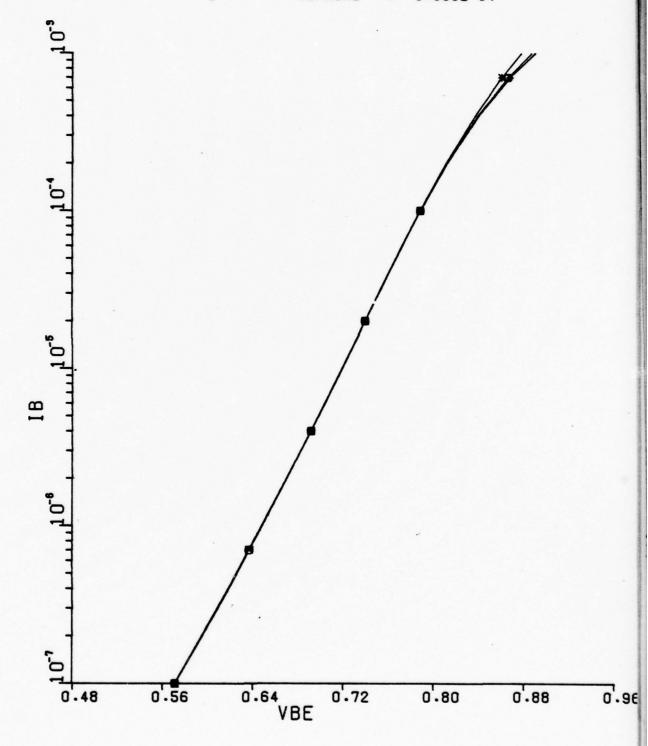


1

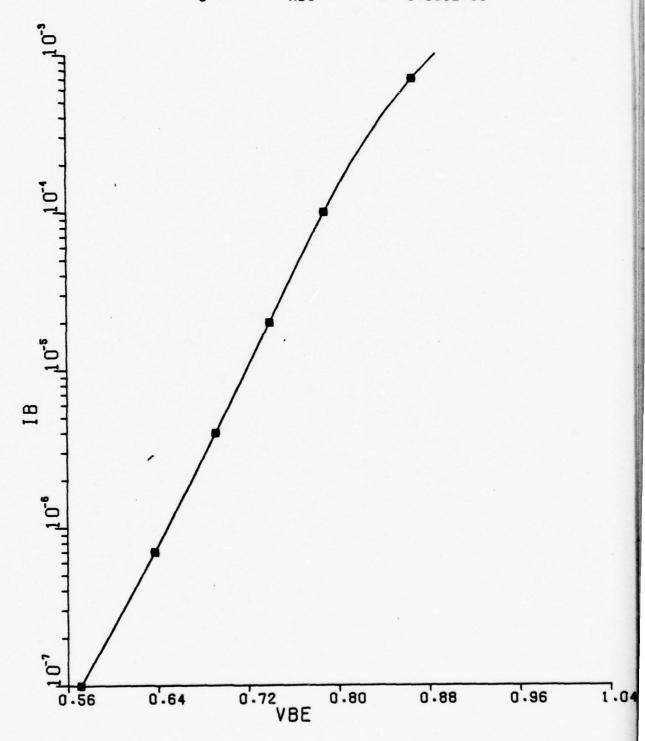




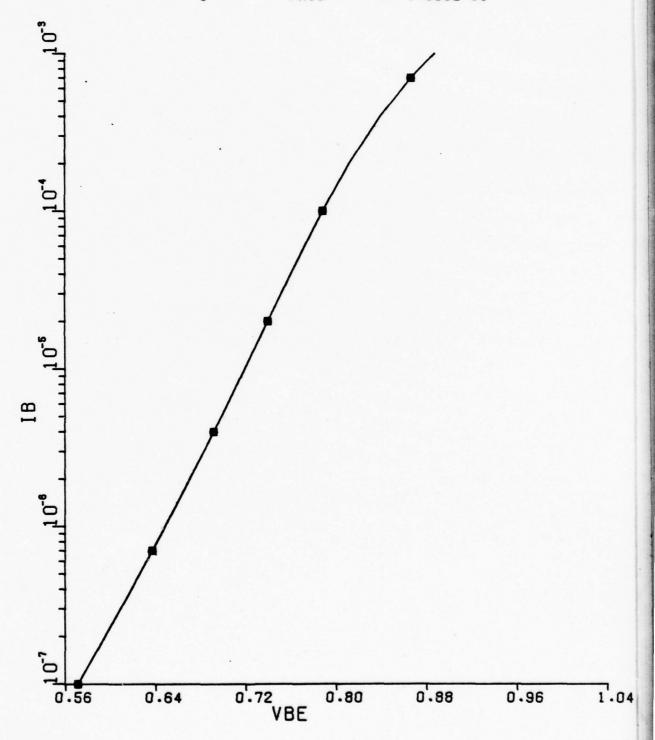
# -- WCPRIME = 1.000E-03 + -- WCPRIME = 1.000E-04 0 -- WCPRIME = 3.000E-04

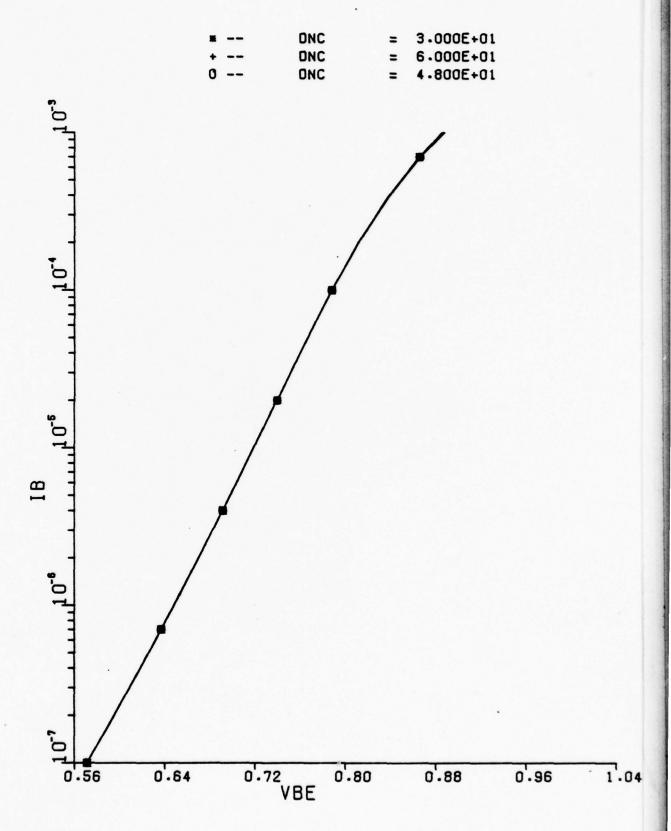


# -- NDC = 1.000E+20 + -- NDC = 1.000E+13 0 -- NDC = 1.000E+16

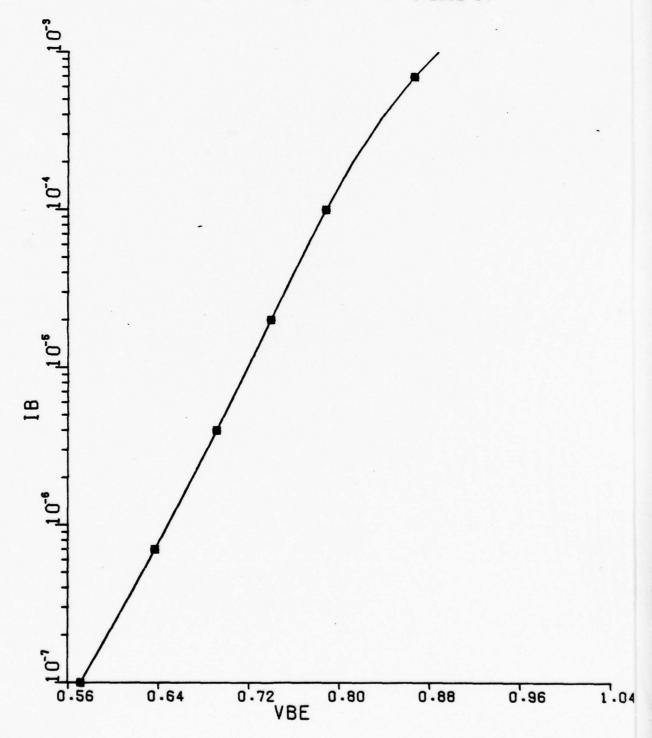


\* -- PHIC = 4.000E-01 + -- PHIC = 1.200E+00 0 -- PHIC = 7.000E-01

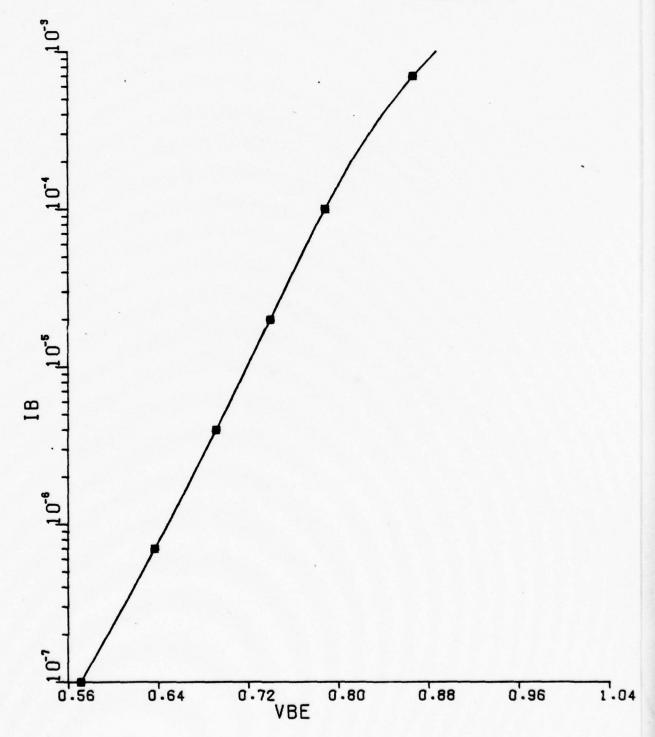




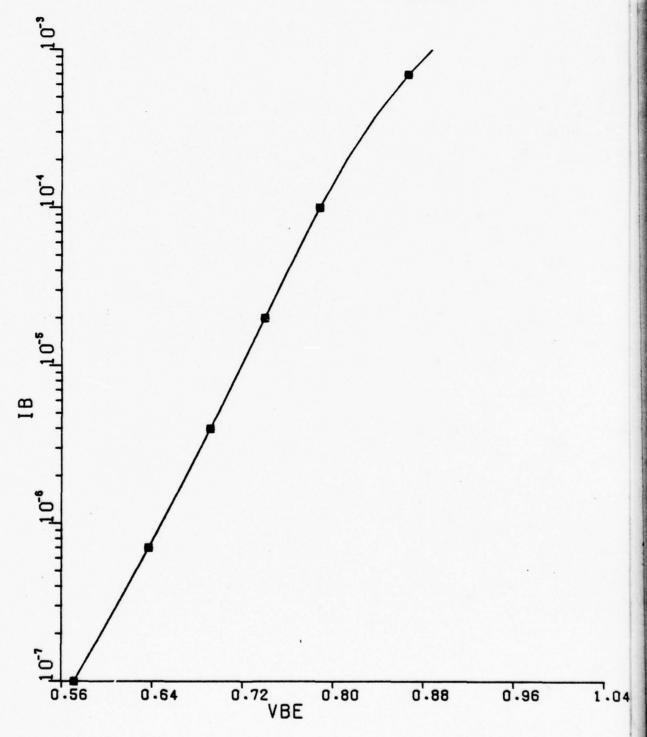
x -- X0 = 4.000E-05 + -- X0 = 2.400E-04 0 -- X0 = 1.200E-04



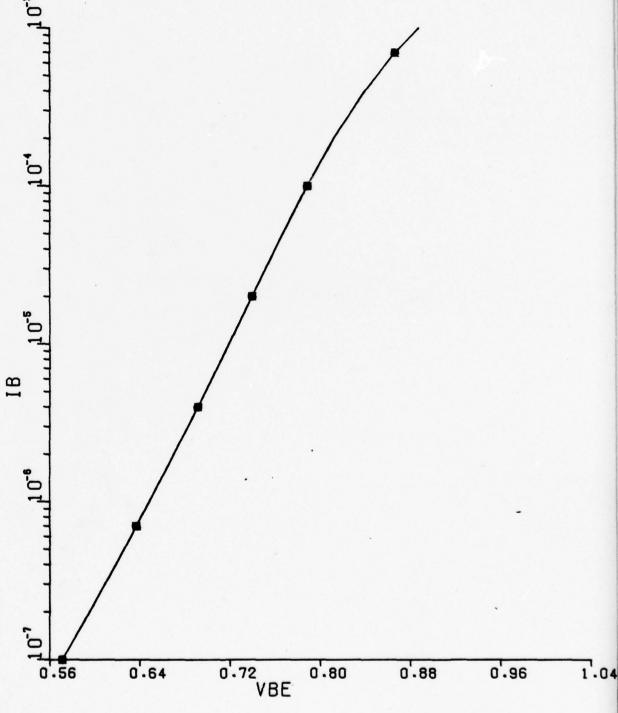
# -- KR = 1.000E-01 + -- KR = 9.000E-01 0 -- KR = 4.000E-01



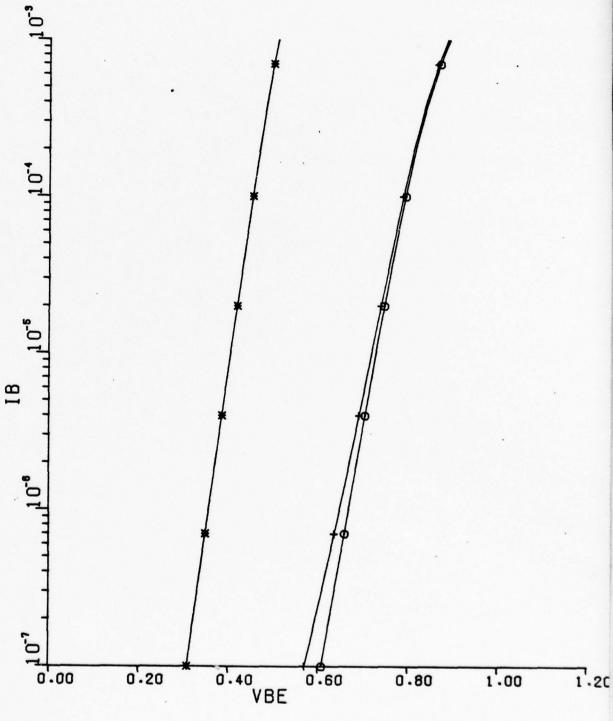
TO = 2.000E-12 +-- TO = 2.000E-10 0-- TO = 2.000E-11



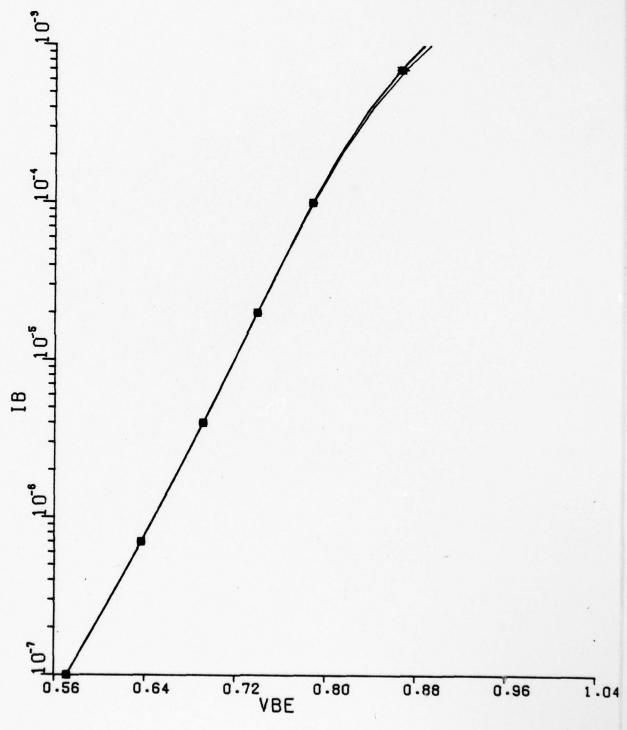


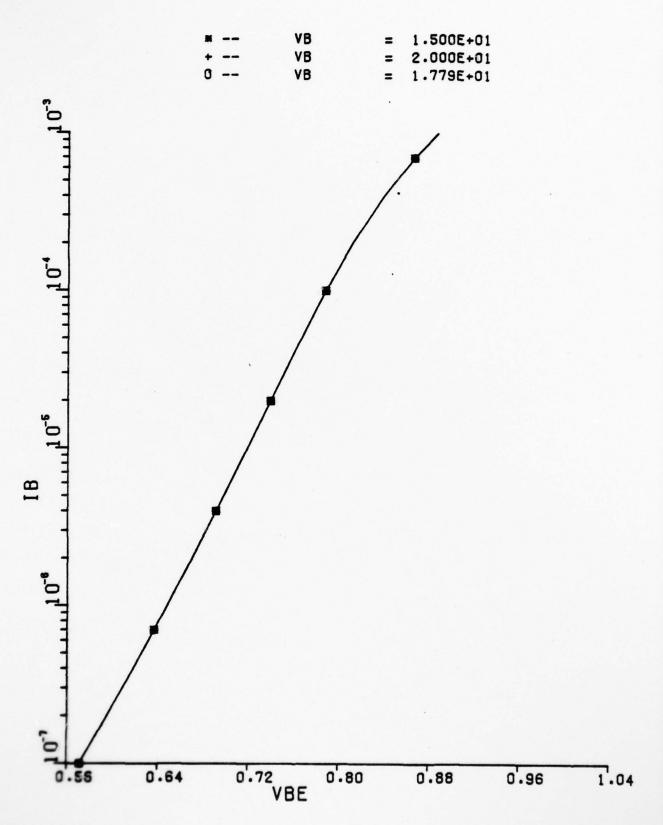


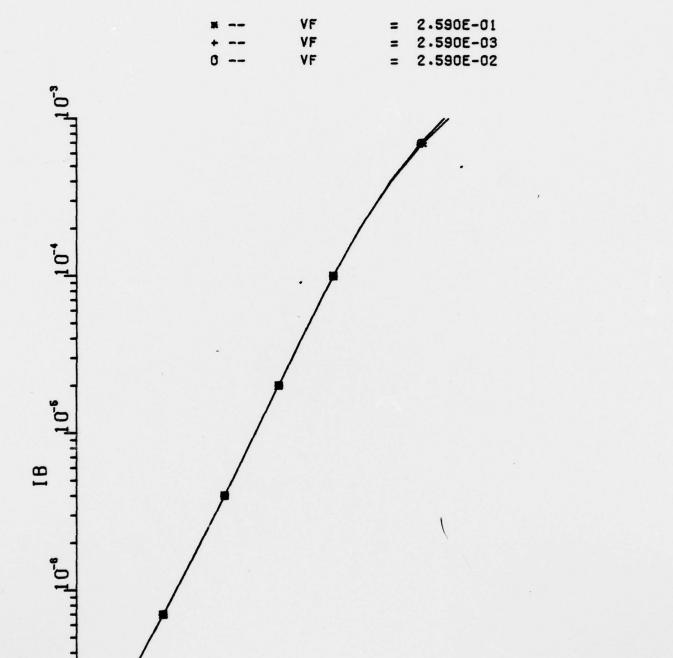












0.80

0.88

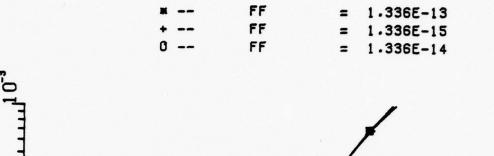
1.04

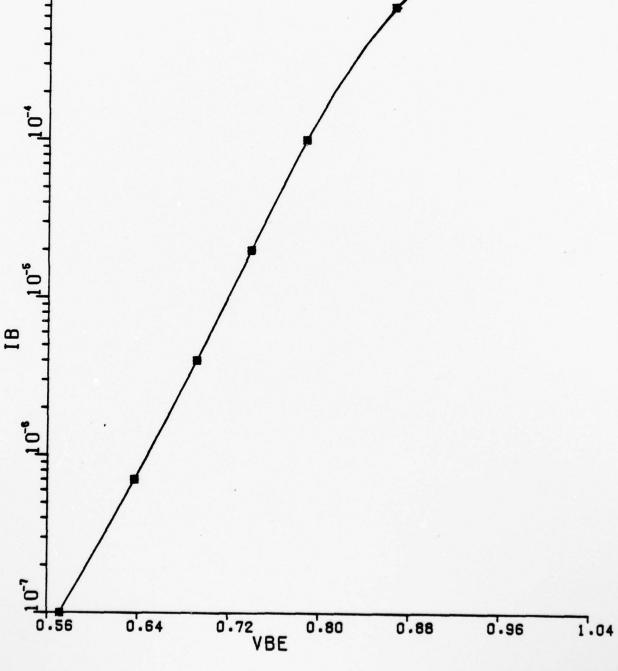
0.96

0'.72 VBE

0.64

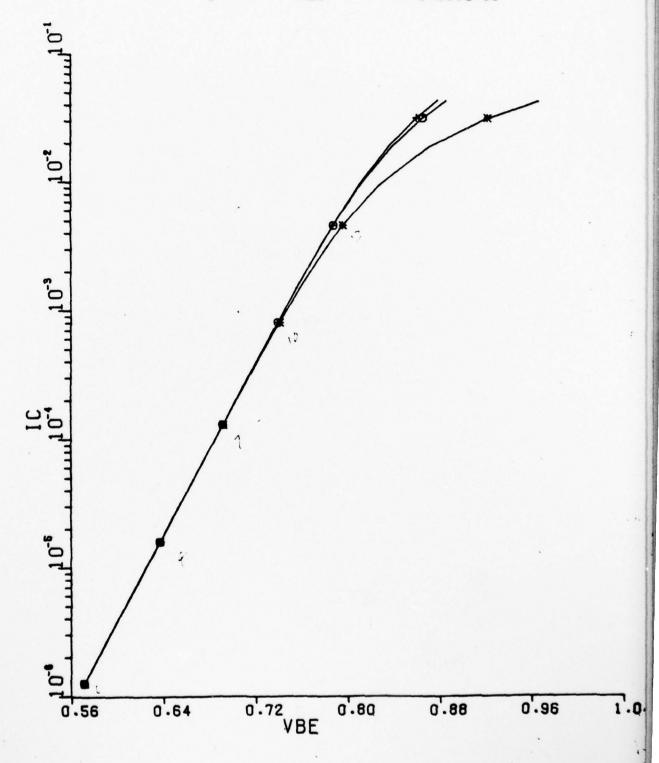
0.56





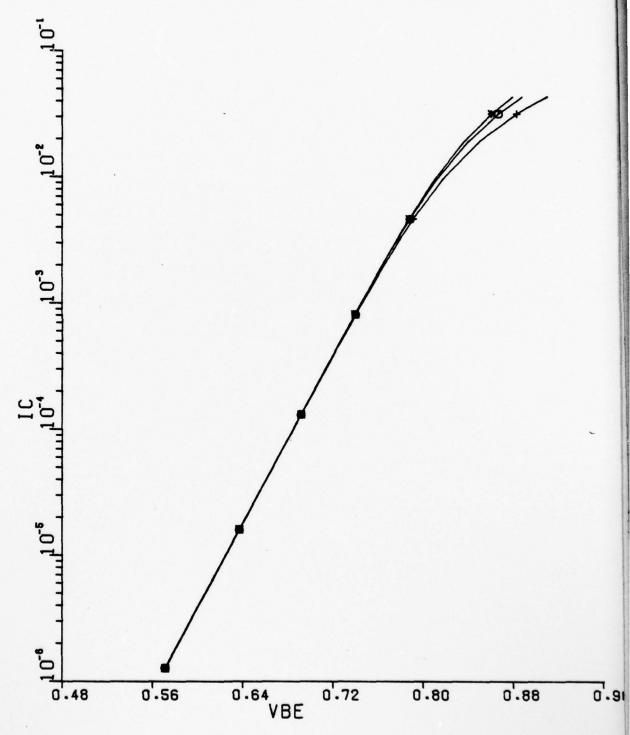
B.2 I<sub>c</sub> vs. V<sub>BE</sub> Curves

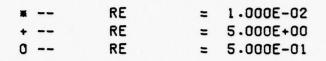
# -- RBO = 9.000E+01 + -- RBO = 9.000E-01 0 -- RBO = 9.000E+00

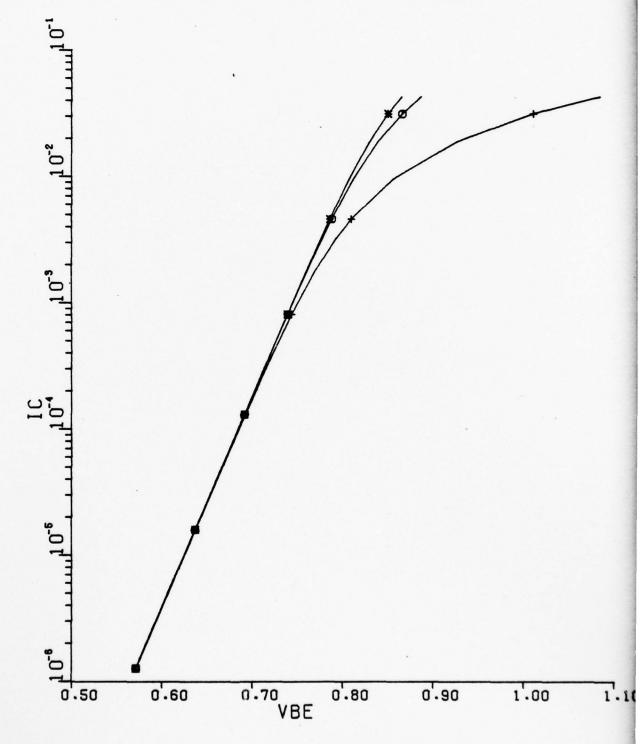


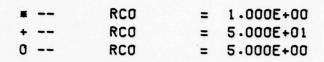
1+.)

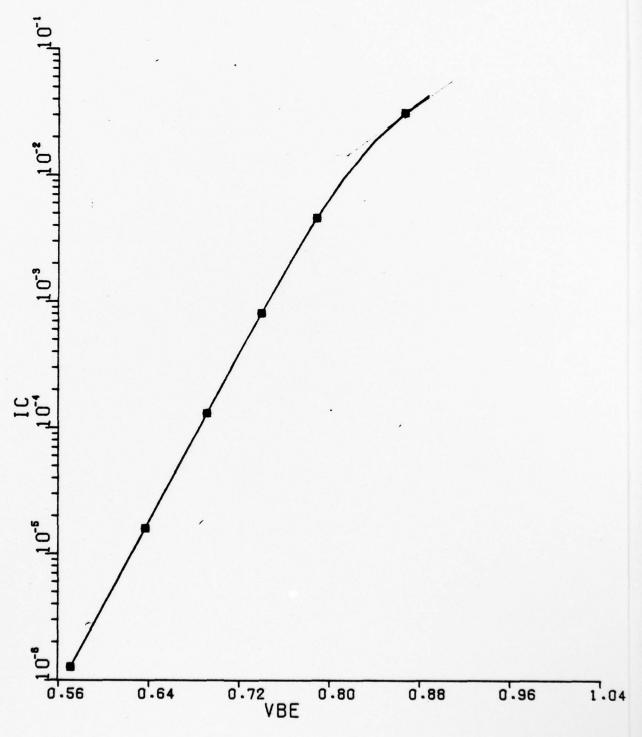




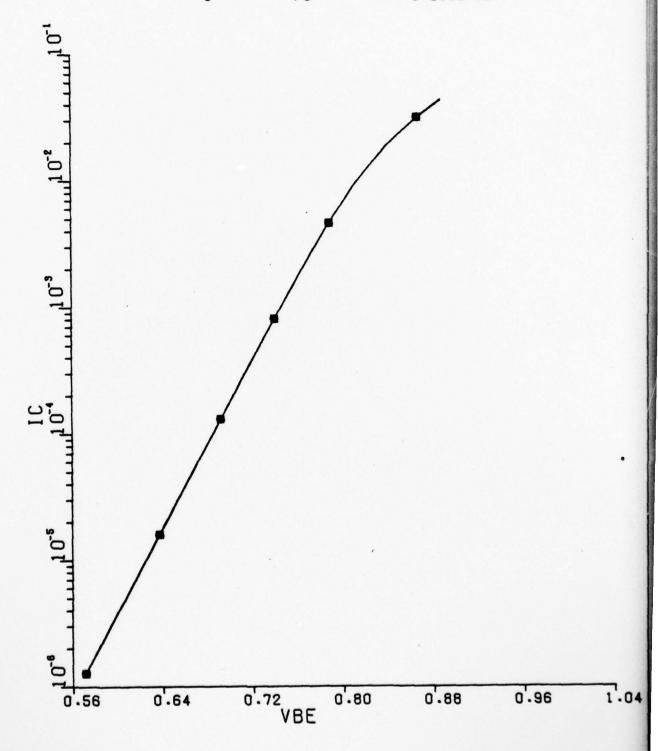




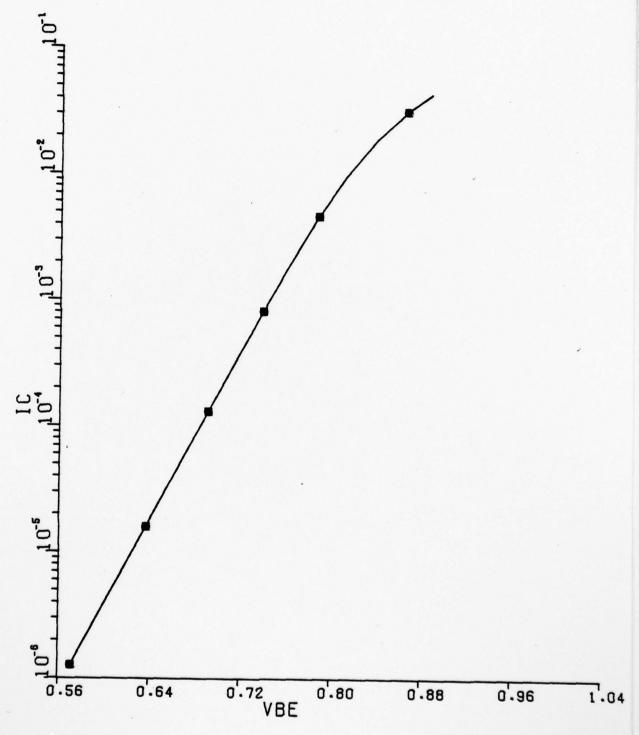




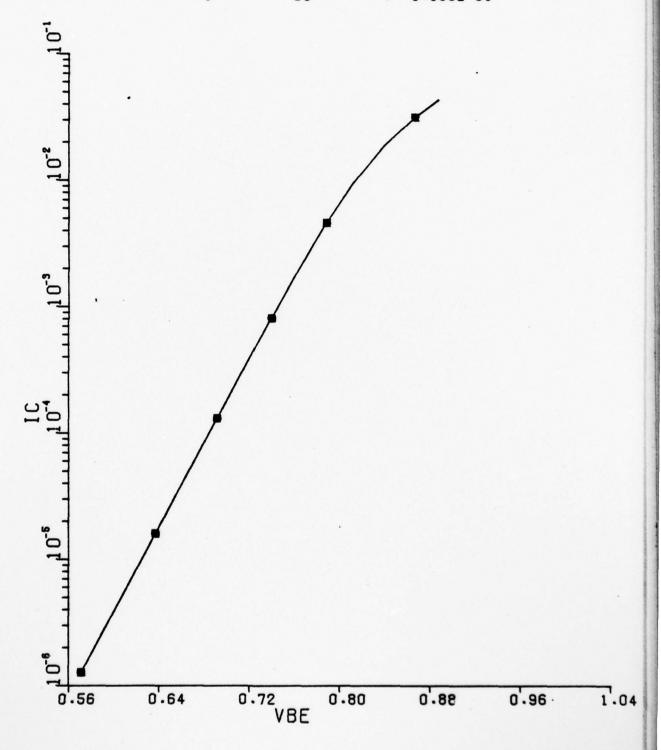
# -- PC = 6.000E-03 + -- PC = 1.000E-03 0 -- PC = 3.250E-03



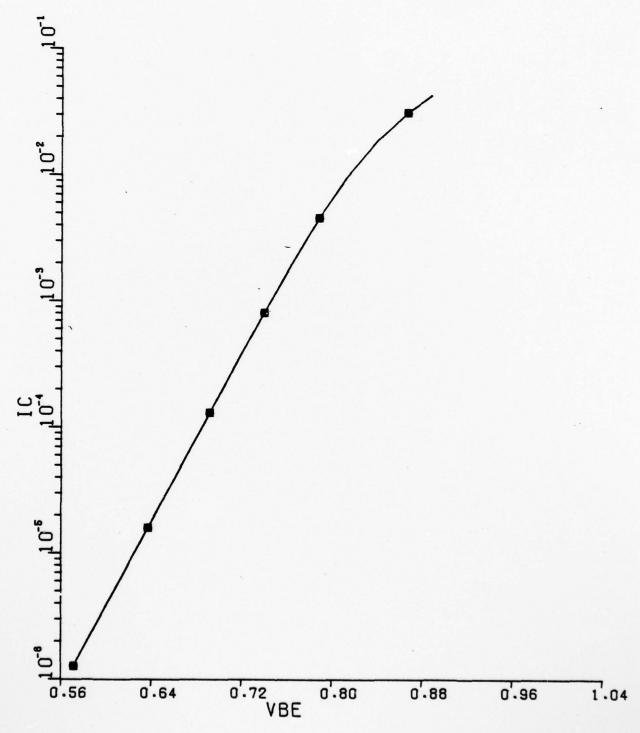




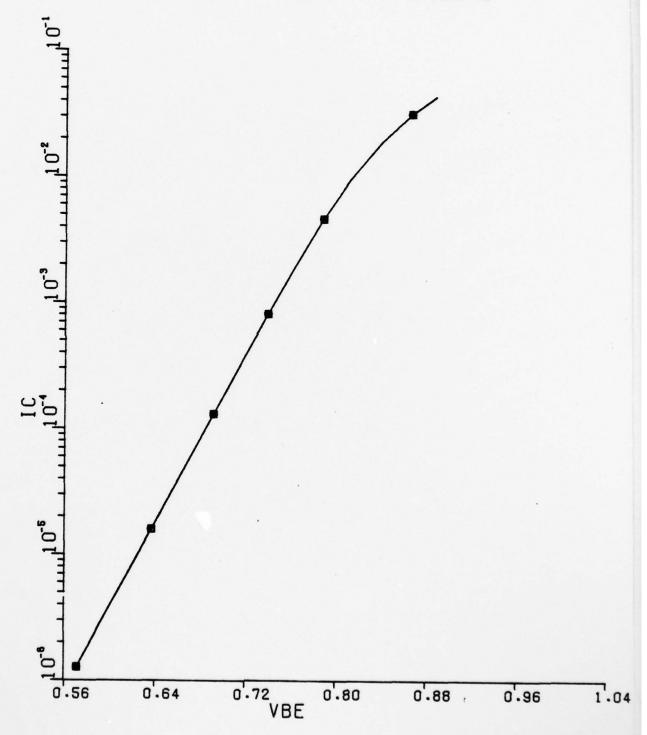
# -- BC = 2.500E-01 + -- BC = 3.333E-01 0 -- BC = 5.000E-01



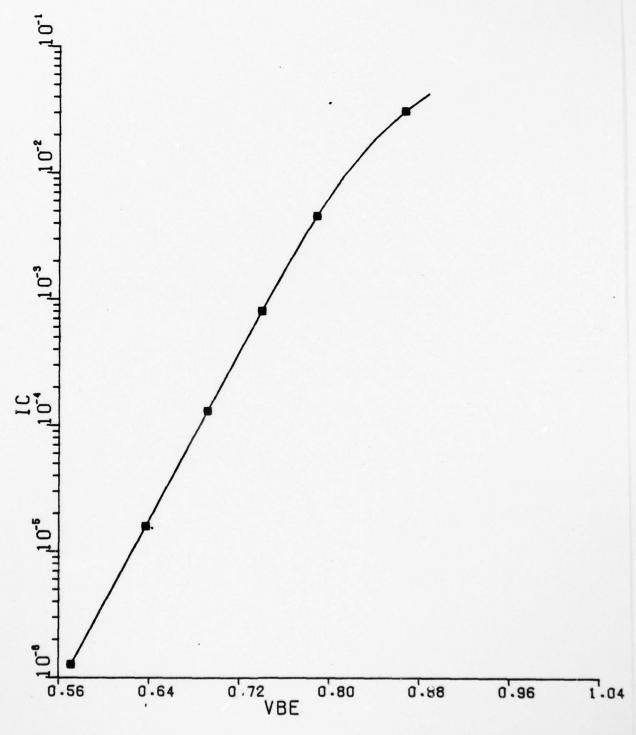


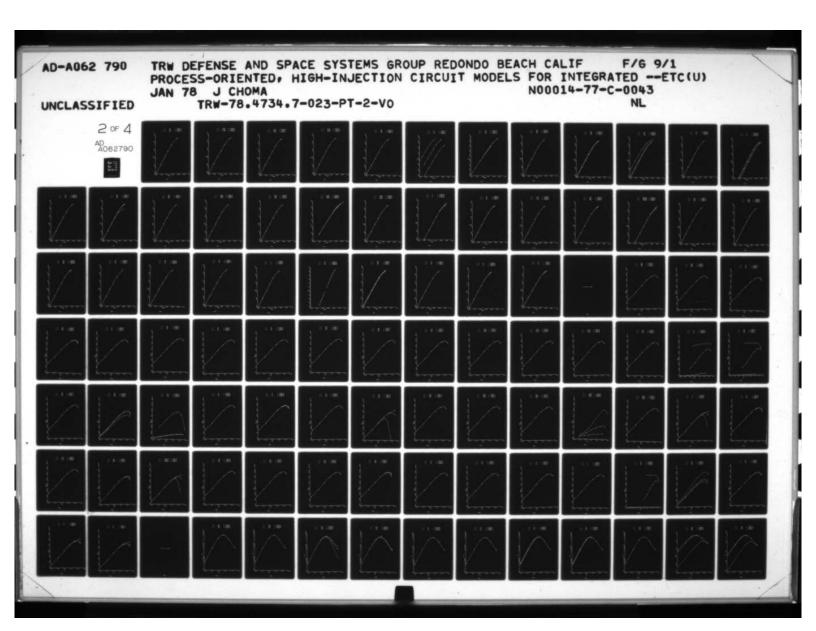


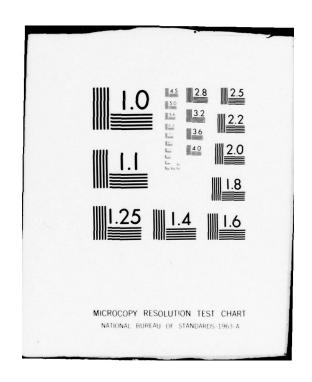
\* -- CSO = 4.000E-13 + -- CSO = 1.000E-11 0 -- CSO = 2.000E-12

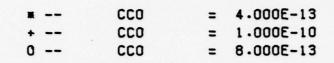


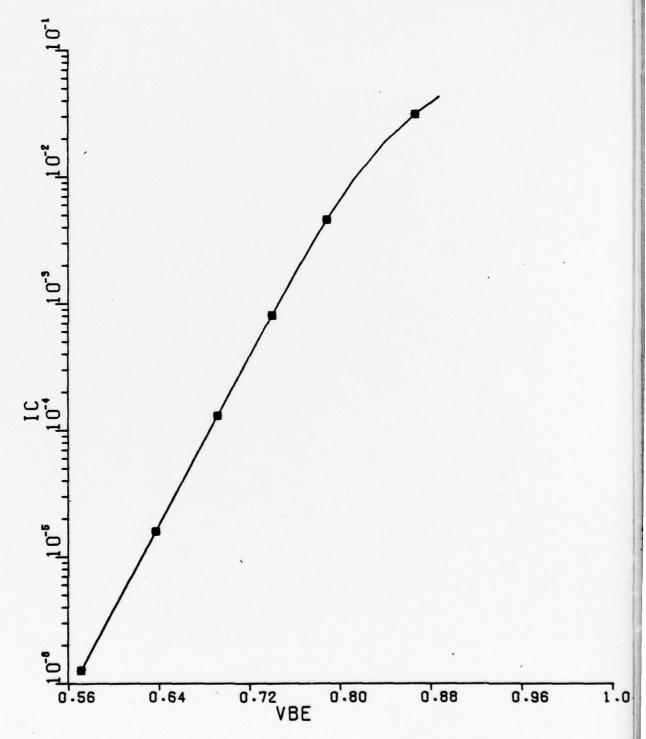




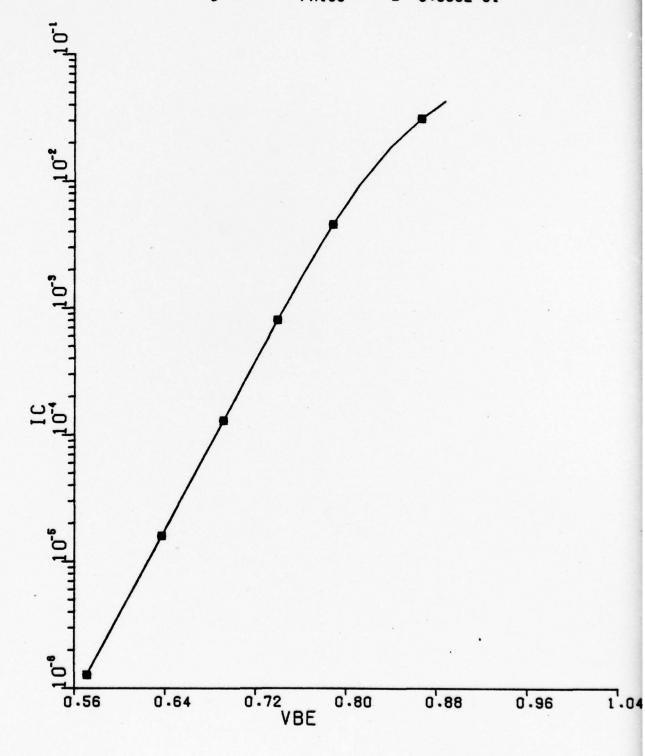




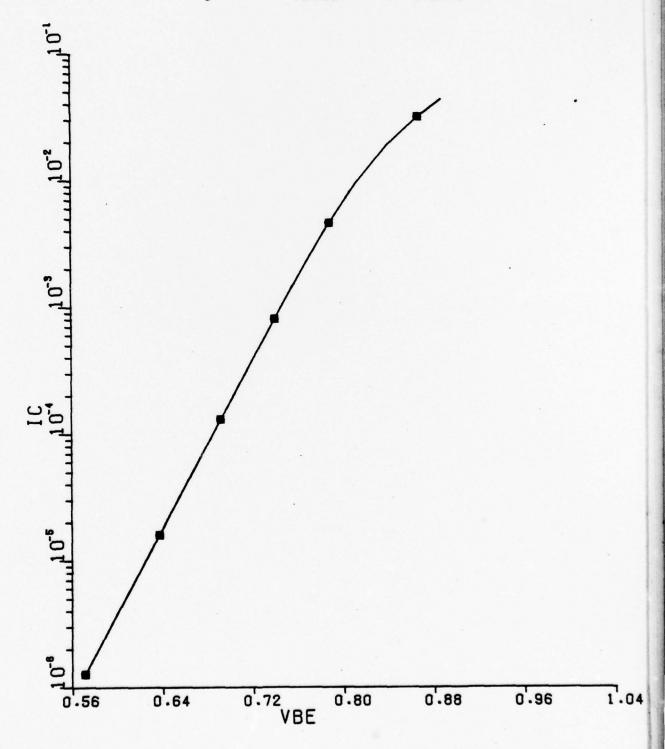




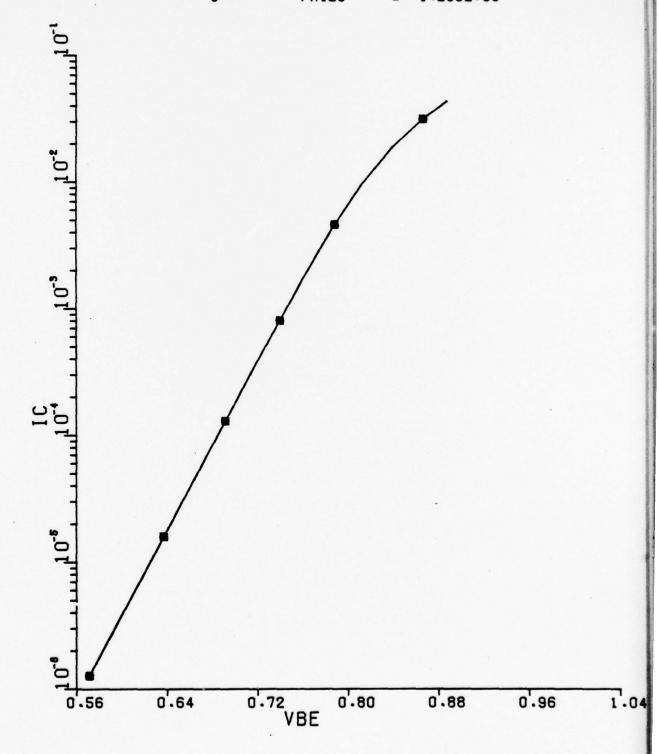
# -- PHISS = 5.000E-01 + -- PHISS = 1.200E+00 0 -- PHISS = 9.000E-01



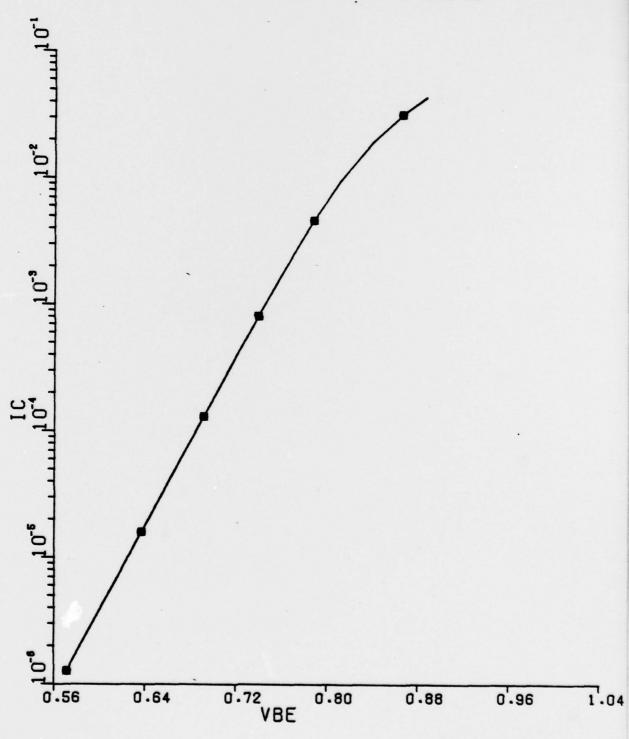
# -- PHICO = 5.000E-01 + -- PHICO = 1.200E+00 0 -- PHICO = 9.000E-01



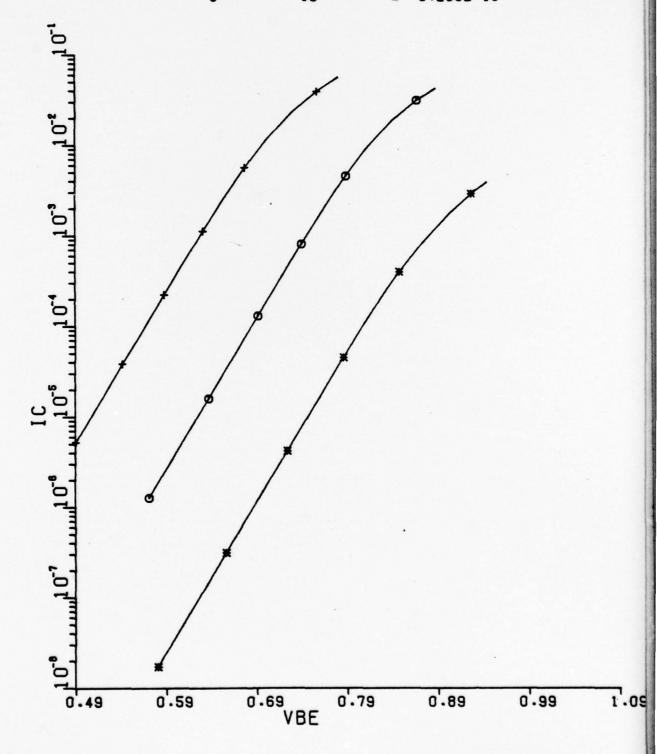
# -- PHIEO = 5.000E-01 + -- PHIEO = 1.500E+00 0 -- PHIEO = 1.200E+00

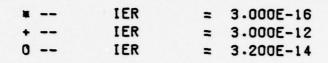


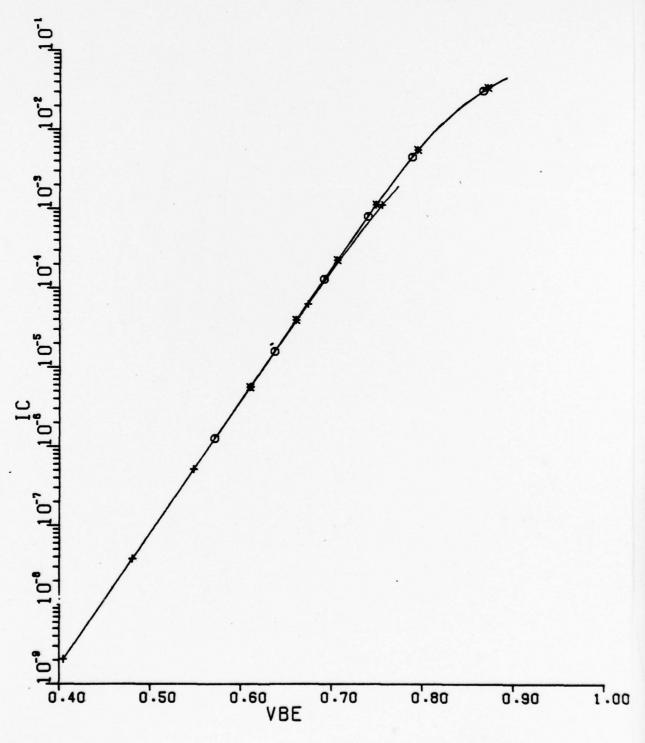




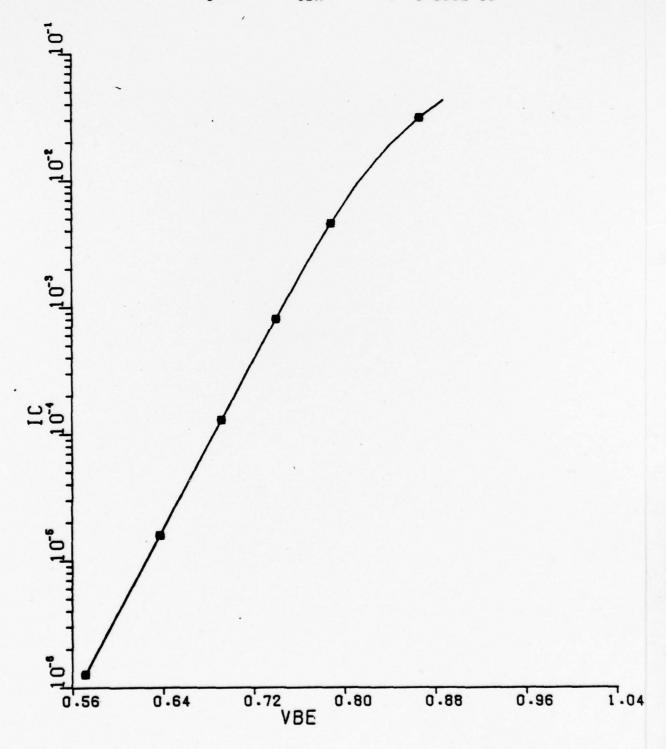
# -- IS = 3.000E-18 + -- IS = 3.000E-14 0 -- IS = 3.200E-16



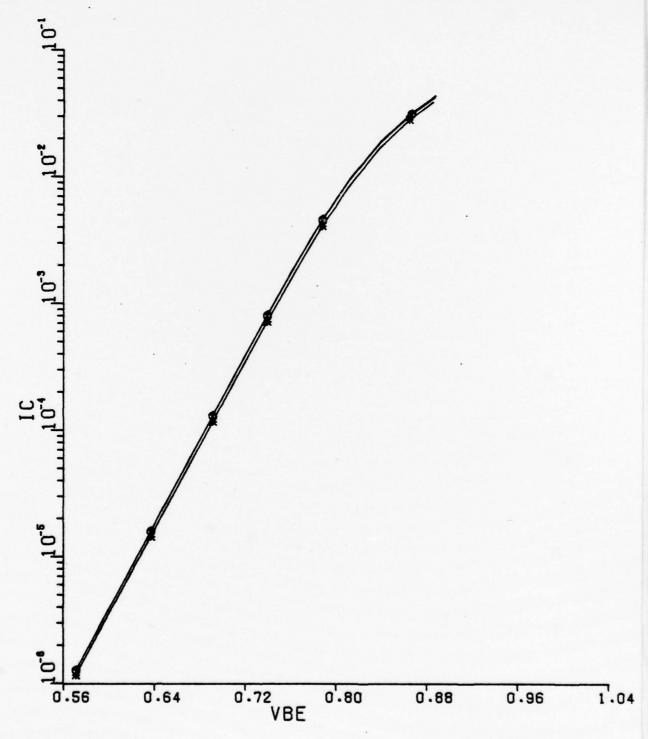


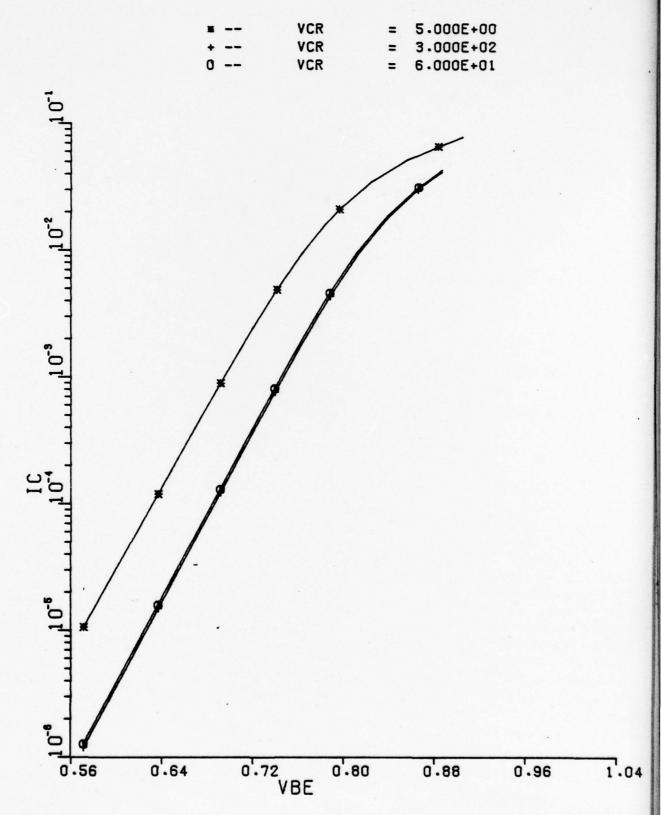


# -- IBR = 3.000E-15 + -- IBR = 3.000E-11 0 -- IBR = 3.000E-13

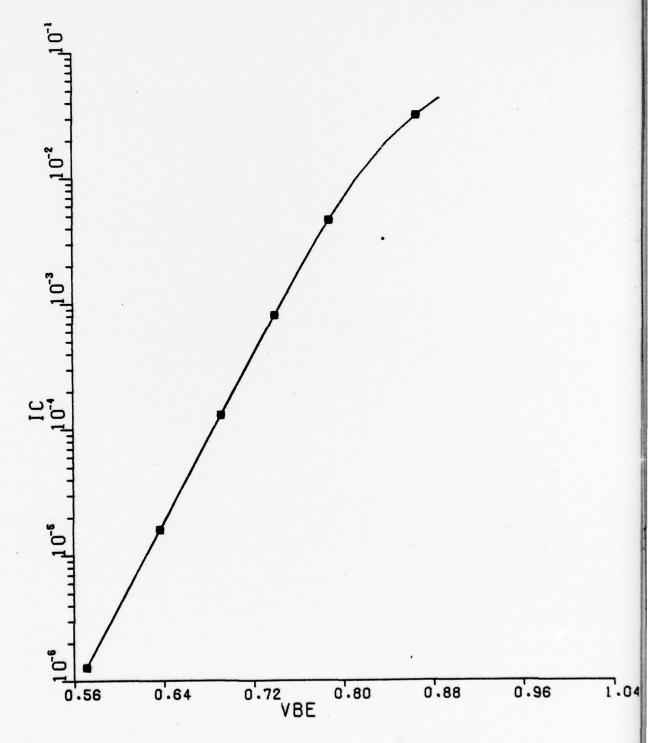




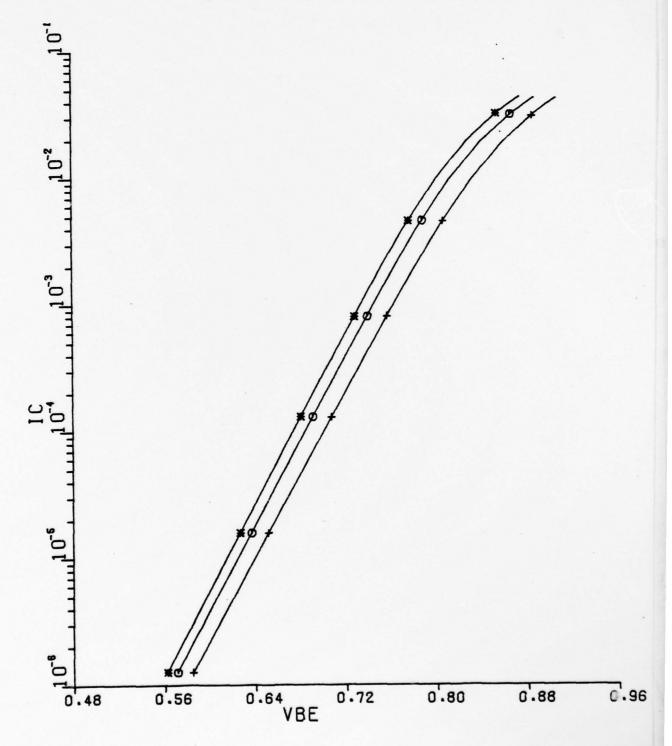




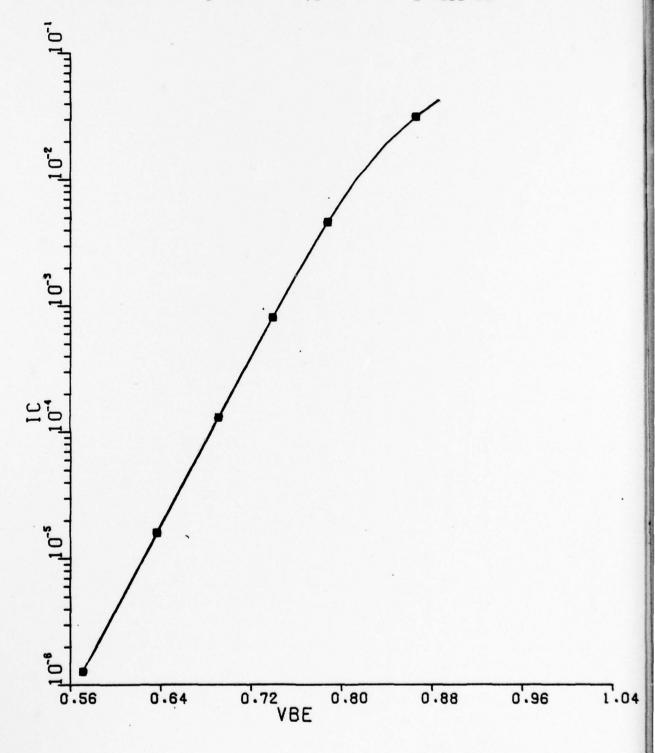
# -- NC = 8.000E-01 + -- NC = 2.300E+00 0 -- NC = 1.800E+00

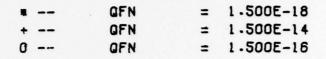


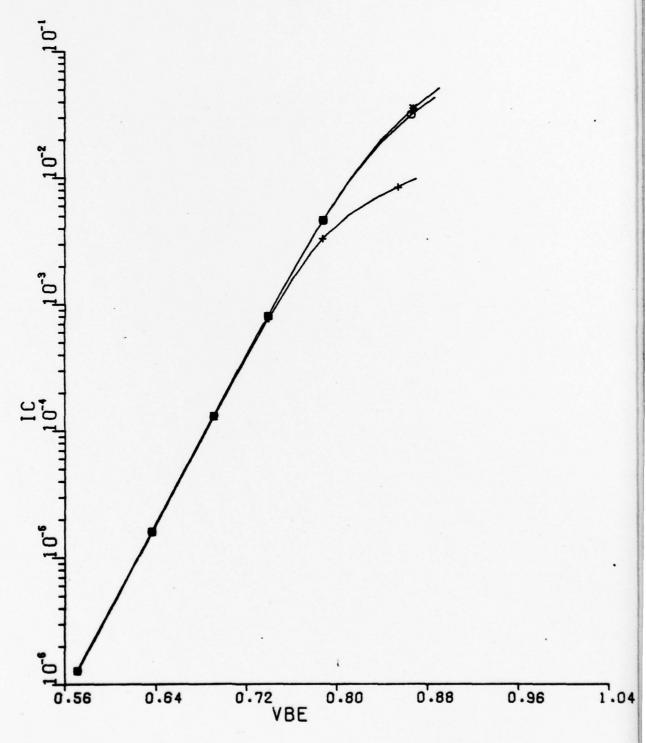
\* -- VT = 2.550E-02 + -- VT = 2.650E-02 0 -- VT = 2.590E-02



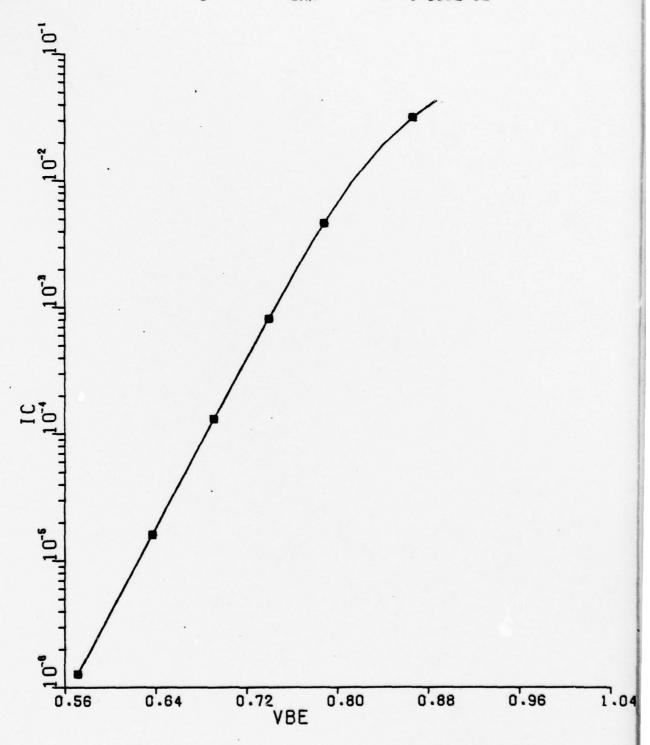
\* -- VTS = 2.600E-02 + -- VTS = 2.500E-02 0 -- VTS = 2.700E-02



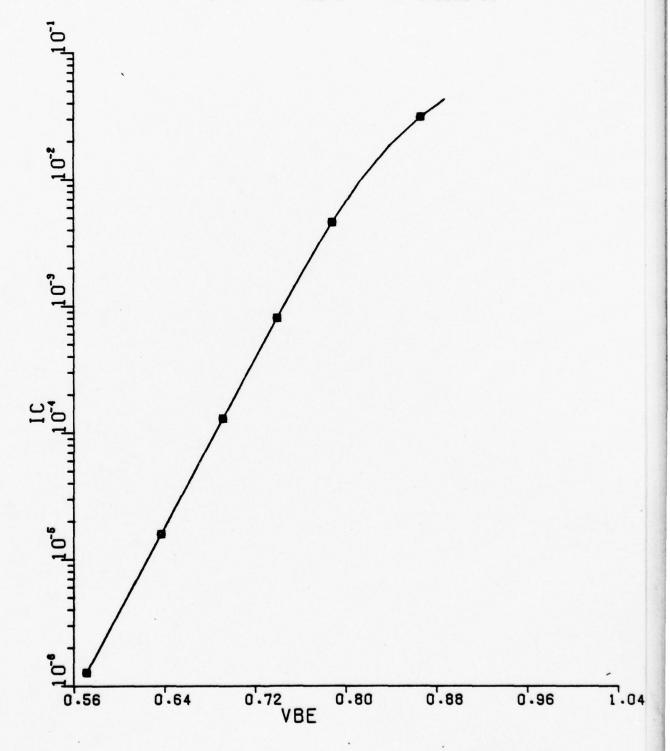




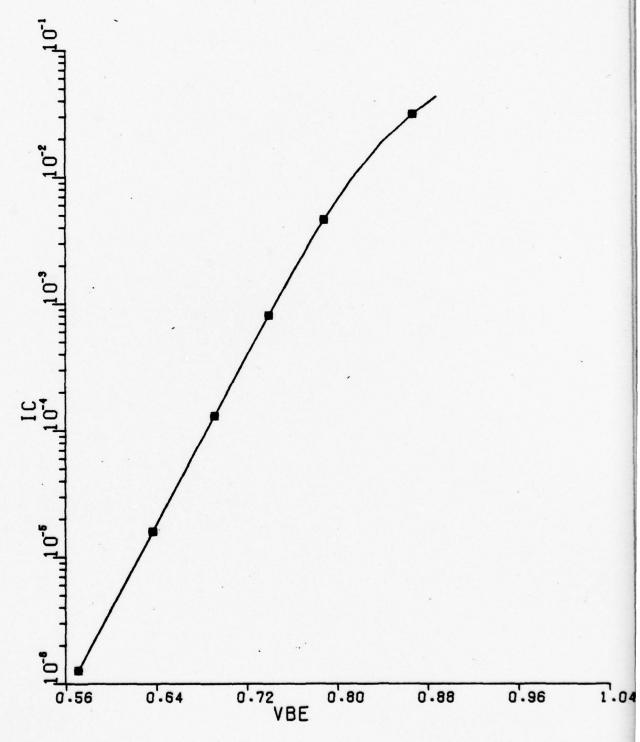
# -- QRN = 1.500E-14 + -- QRN = 1.500E-10 0 -- QRN = 1.500E-12



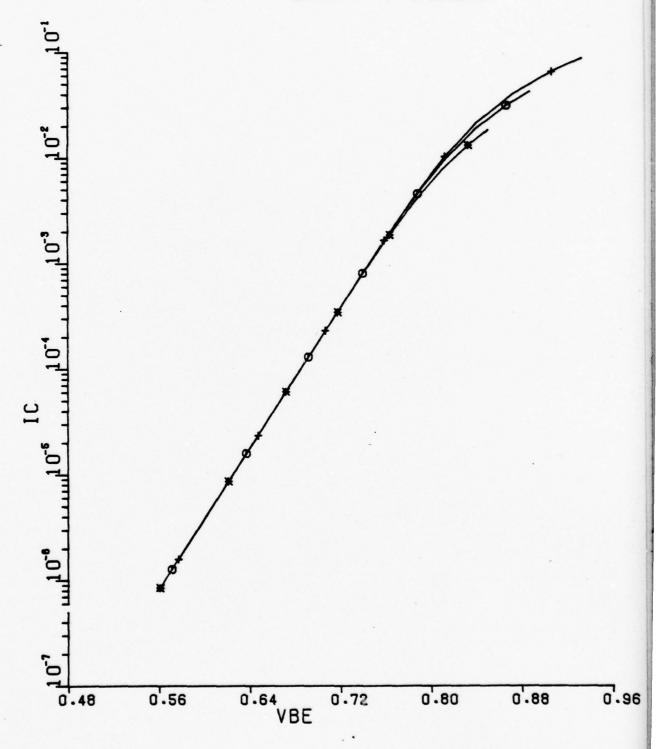
TAUFO = 1.590E-09 +-- TAUFO = 1.590E-08 0-- TAUFO = 1.590E-10



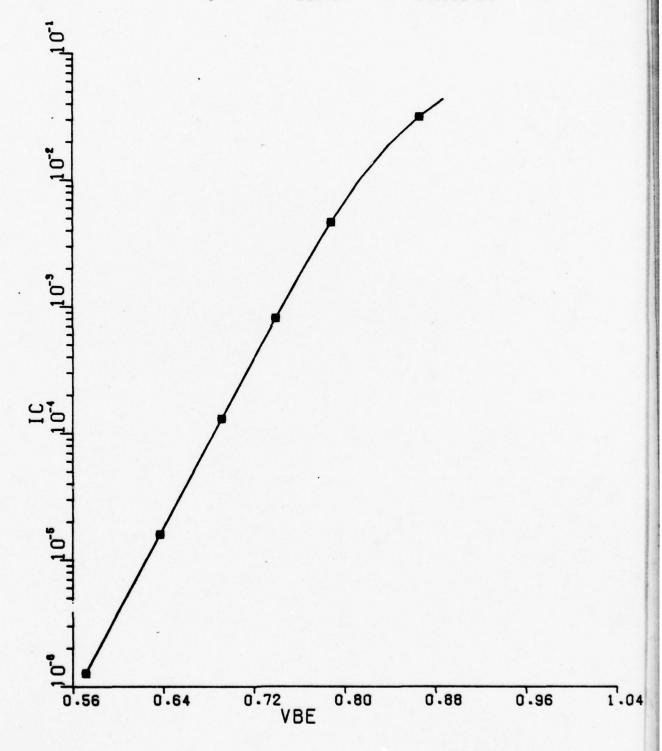




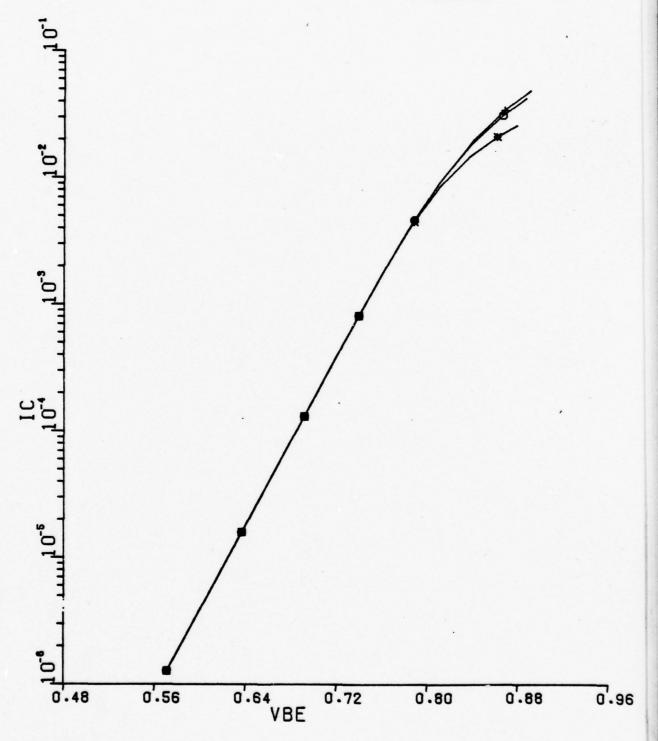
# -- BETAO = 2.000E+01 + -- BETAO = 1.500E+02 0 -- BETAO = 5.500E+01



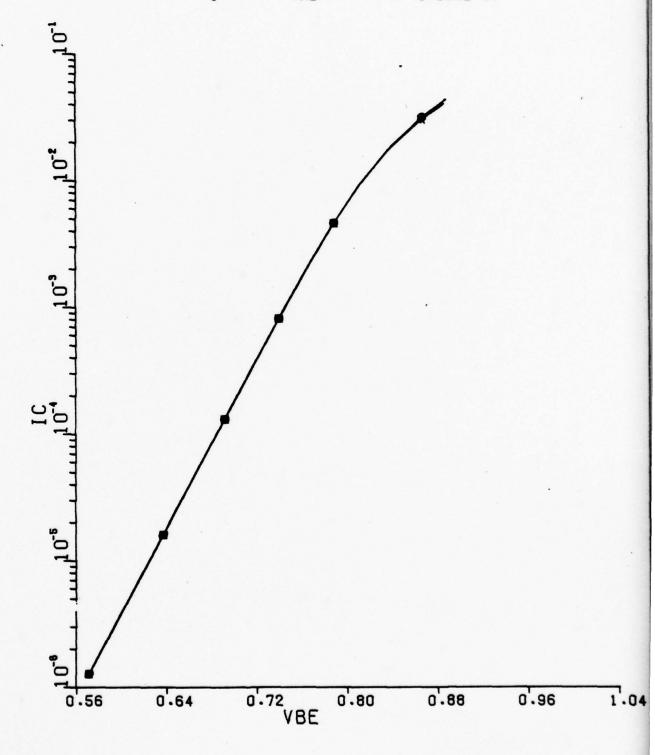
BETAR = 6.000E+00 + -- BETAR = 1.100E+00 0 -- BETAR = 1.500E-01



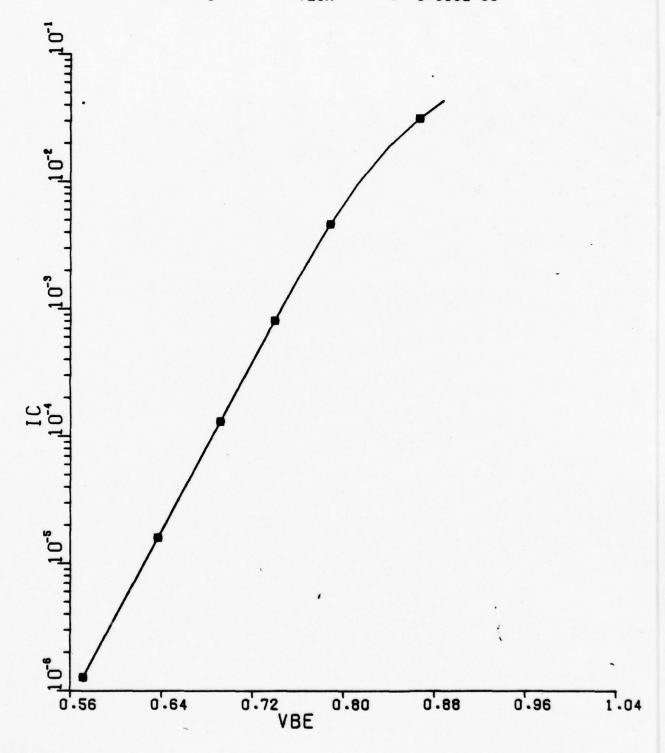
H -- H = 2.000E-05 + -- H = 2.000E-04 0 -- H = 6.000E-05



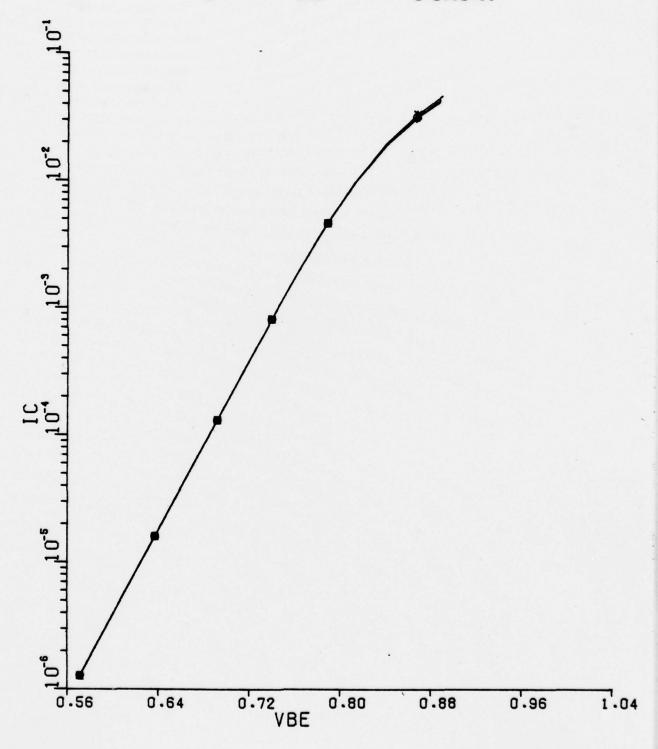
# -- AE = 1.000E-08 + -- AE = 1.000E-04 0 -- AE = 9.800E-07



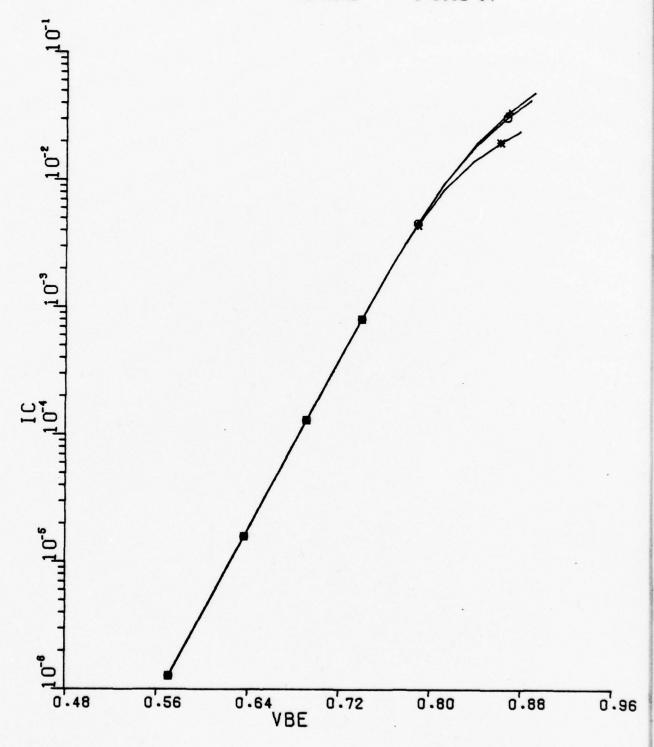
# -- VLIM = 1.000E+06 + -- VLIM = 1.500E+07 0 -- VLIM = 6.000E+06



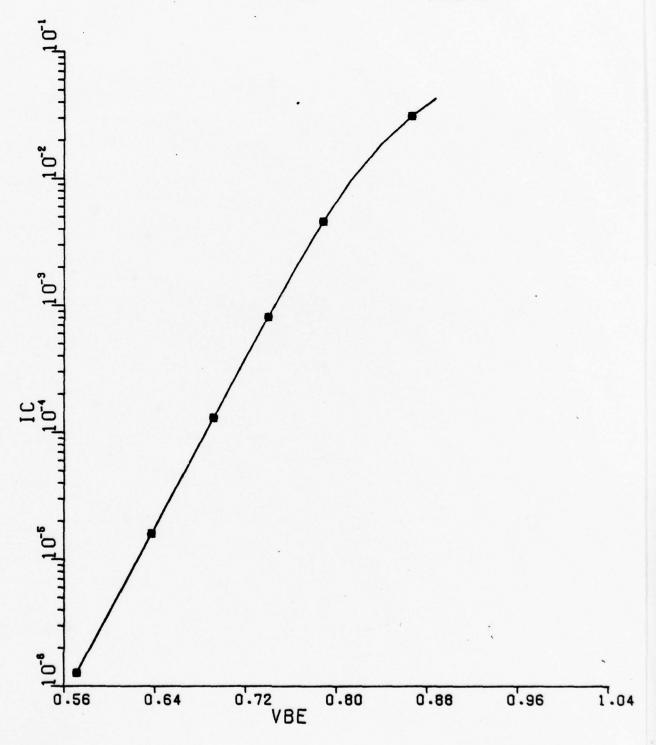
= -- ONB = 1.000E+01 + -- ONB = 3.000E+01 0 -- ONB = 2.260E+01



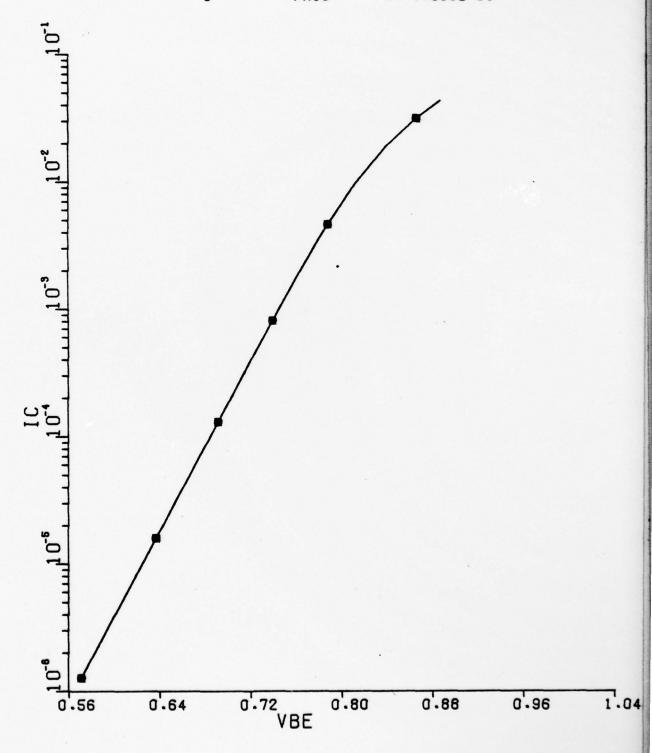
# -- HCPRIME = 1.000E-03 + -- HCPRIME = 1.000E-04 0 -- HCPRIME = 3.000E-04

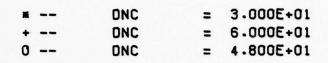


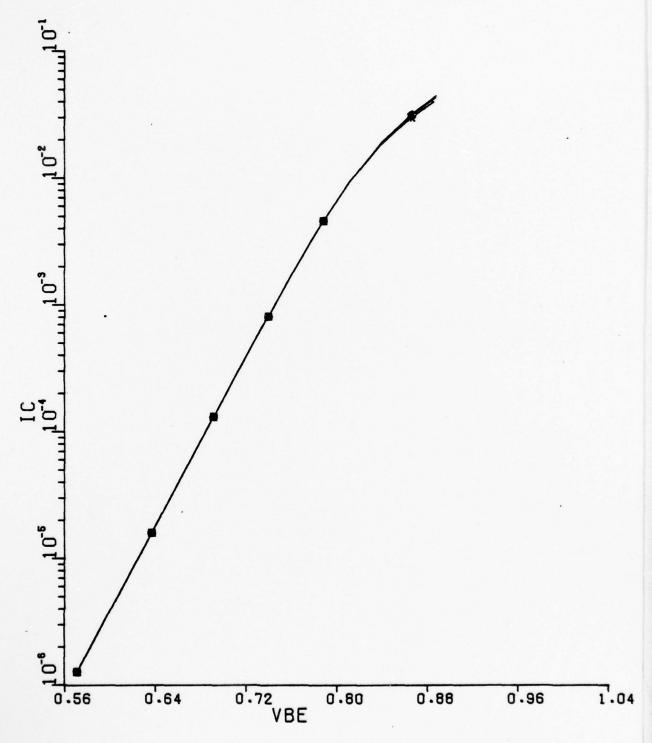
# -- NDC = 1.000E+20 + -- NDC = 1.000E+13 0 -- NDC = 1.000E+16



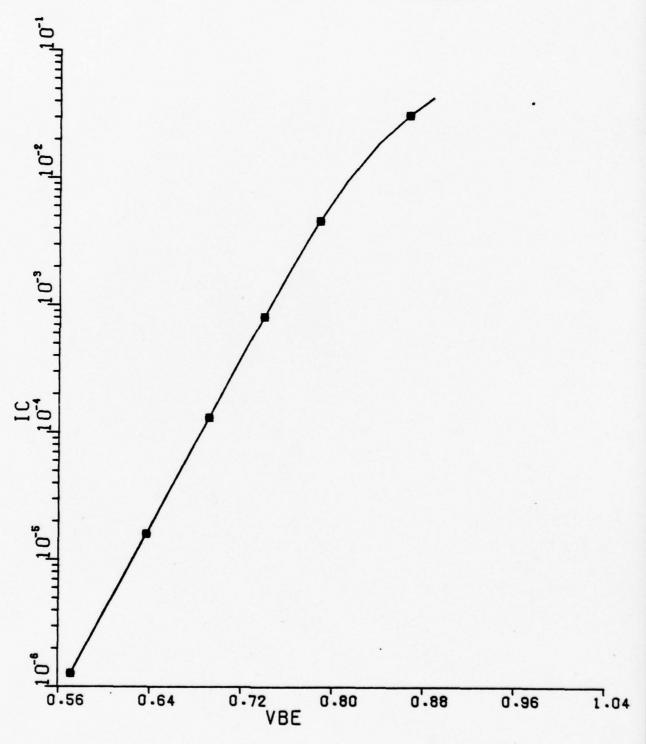
+ -- PHIC = 4.000E-01 + -- PHIC = 1.200E+00 0 -- PHIC = 7.000E-01

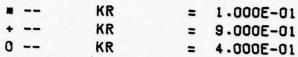


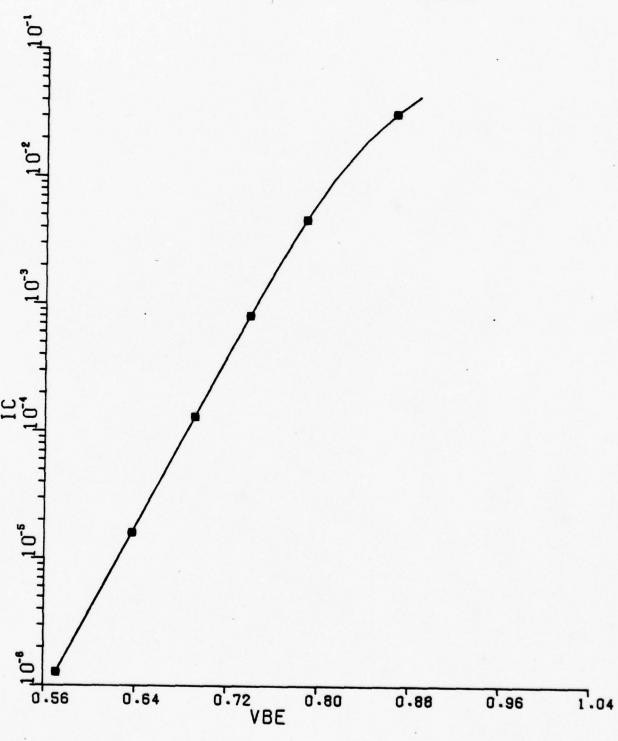


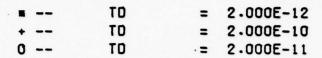


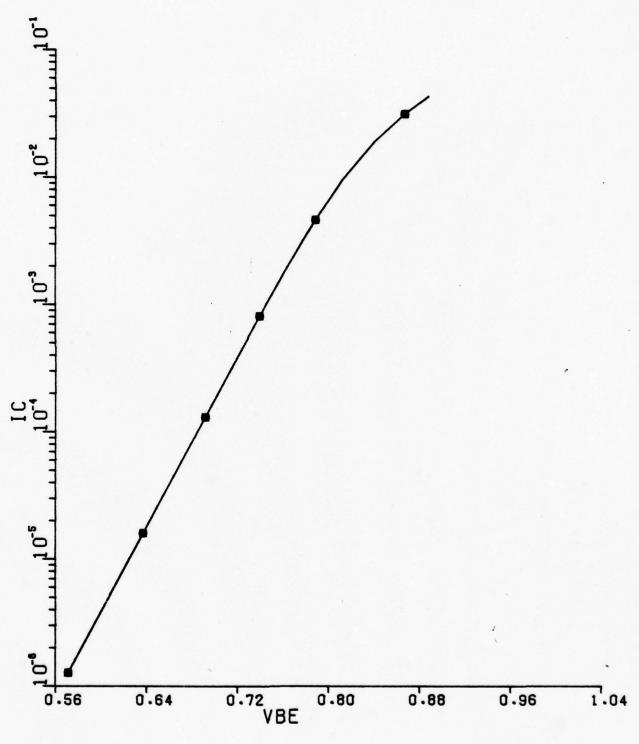




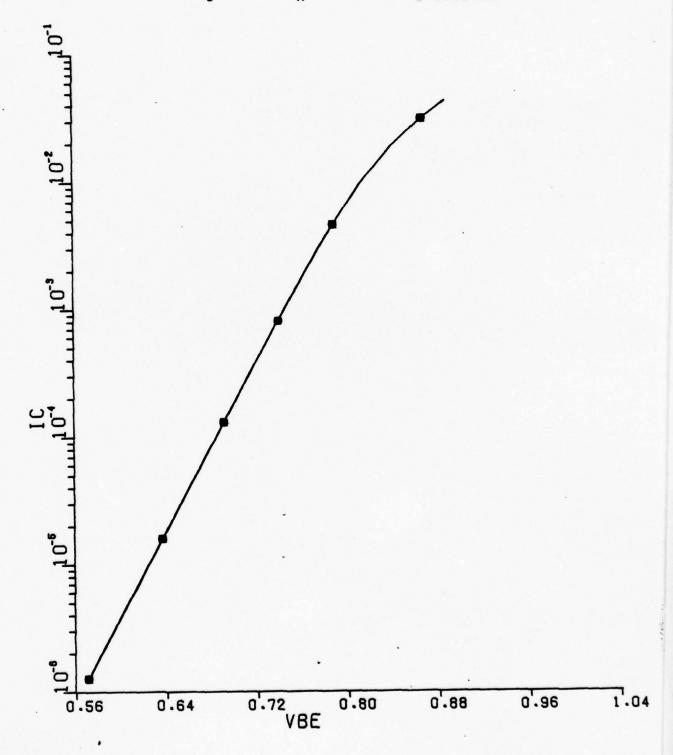


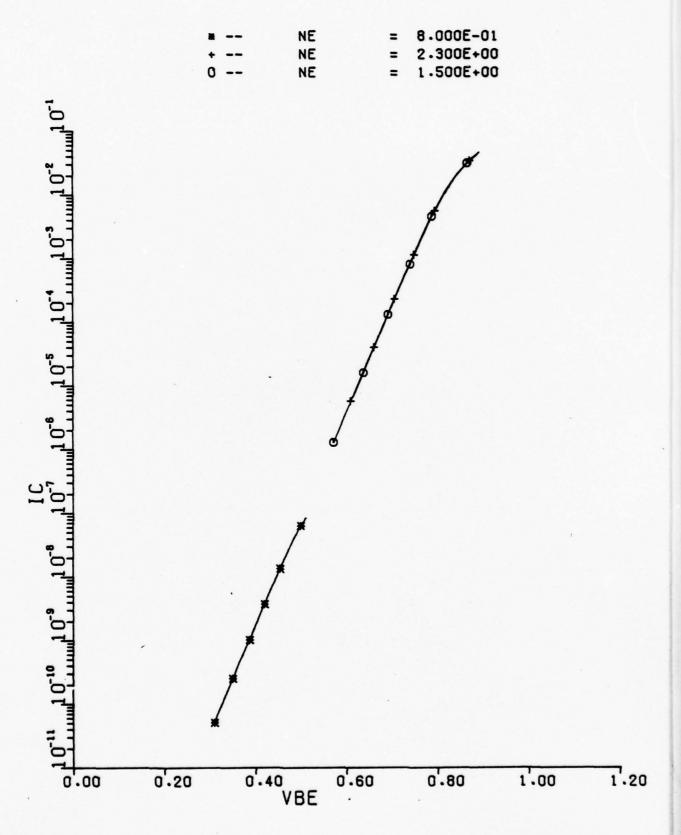




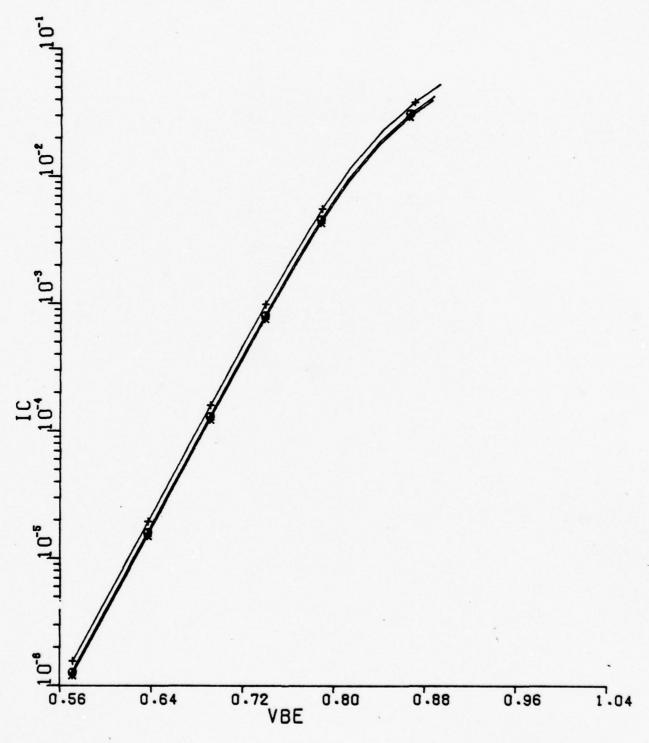


# -- M = 1.000E+00 + -- M = 1.000E+01 0 -- M = 2.000E+00

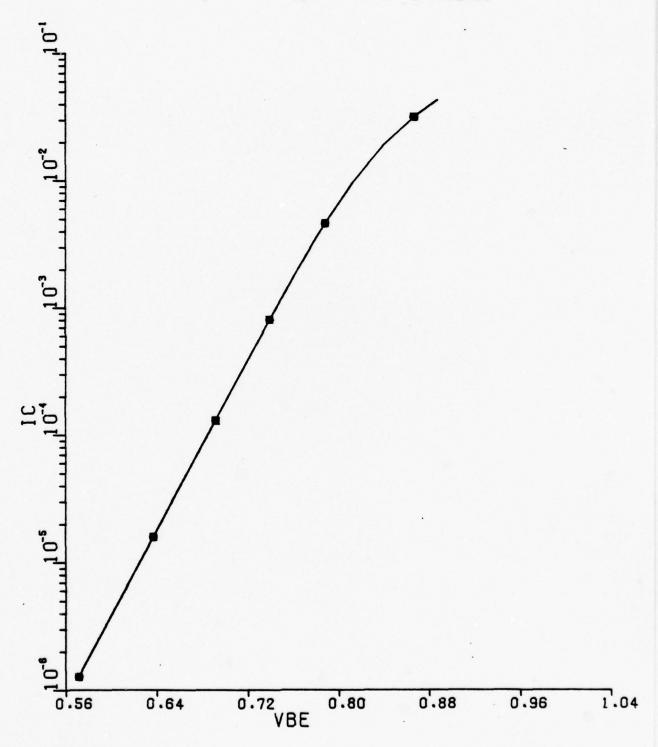


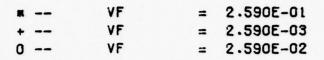


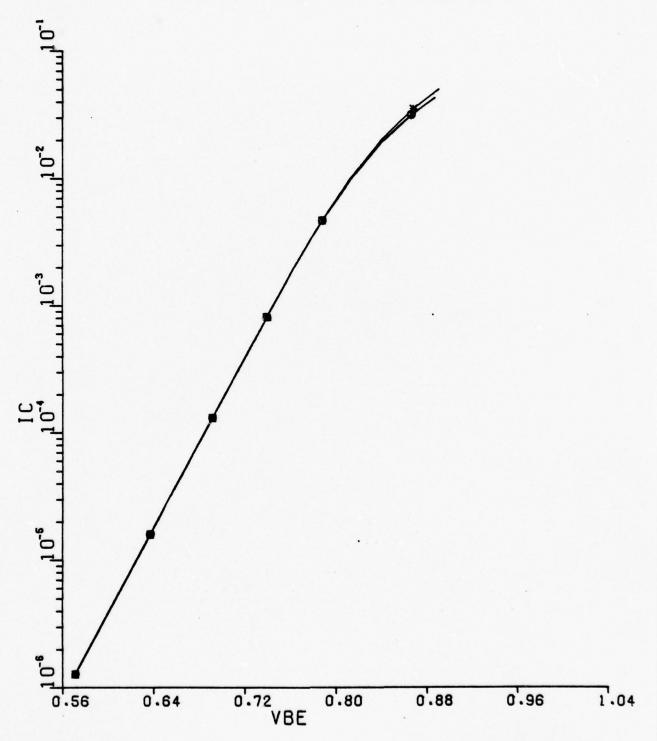




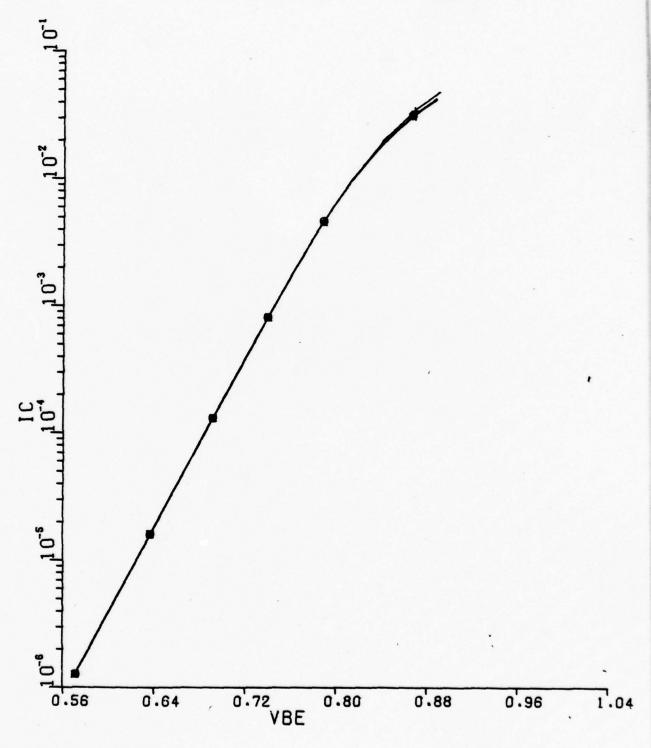
# -- VB = 1.500E+01 + -- VB = 2.000E+01 0 -- VB = 1.779E+01



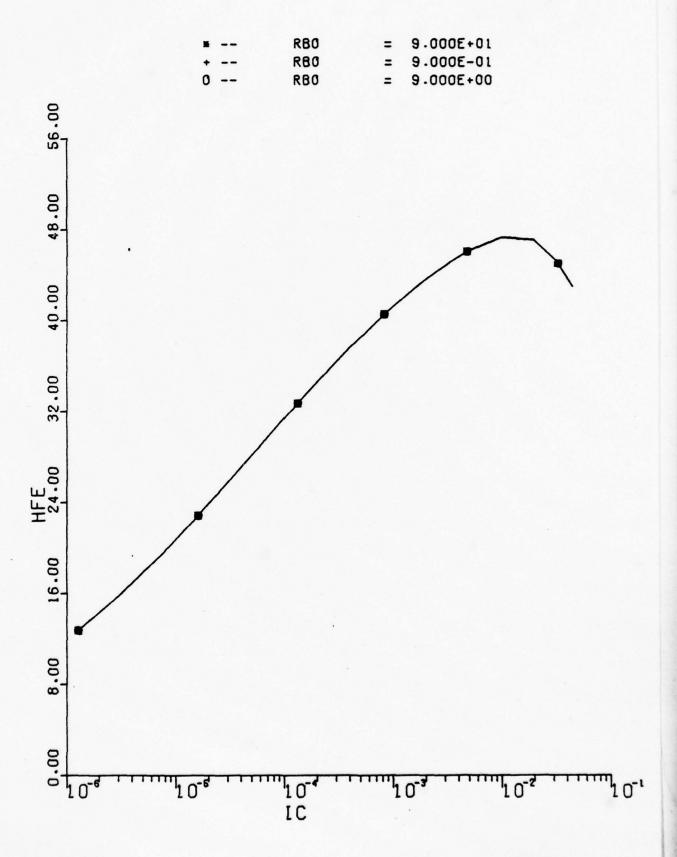


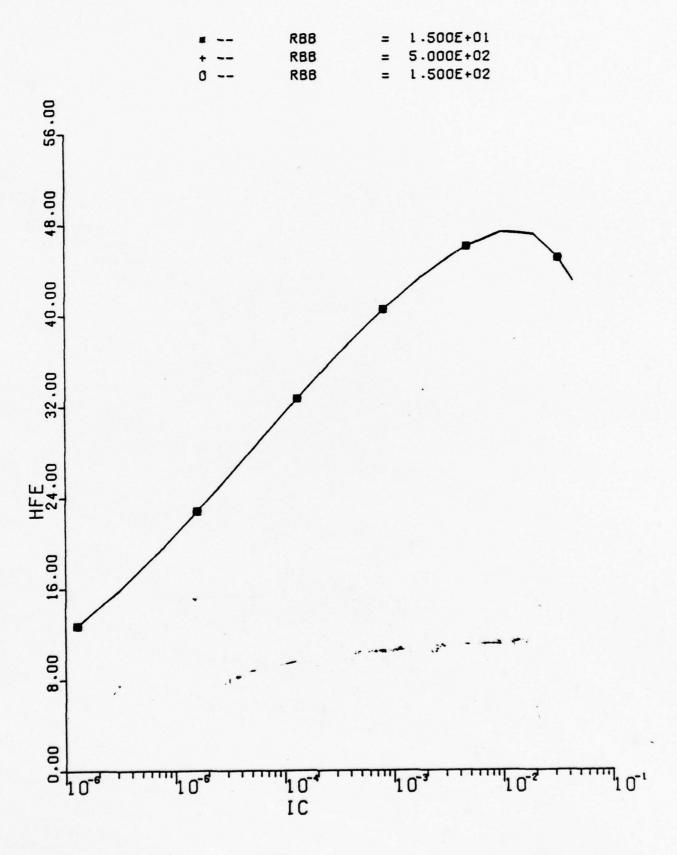


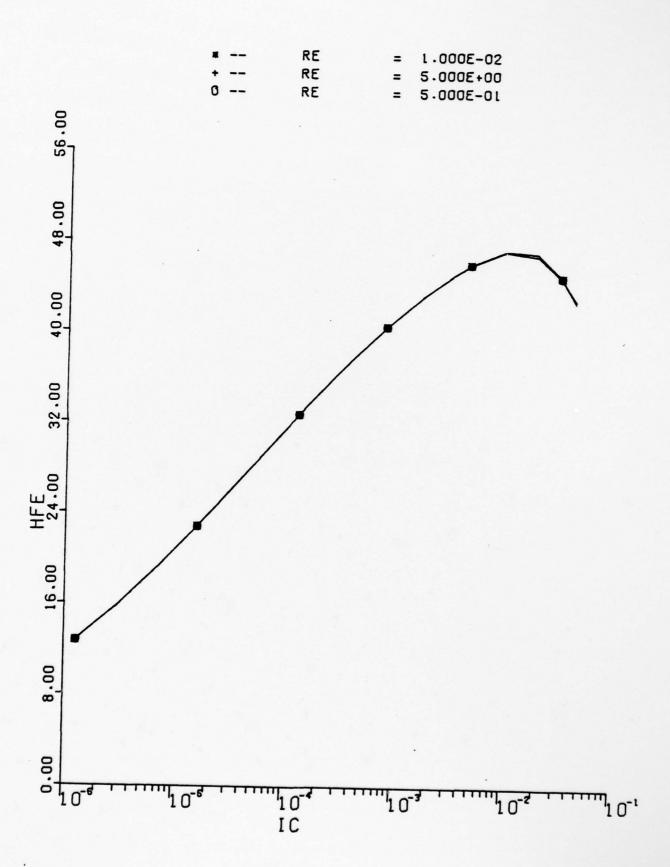
# -- FF = 1.336E-13 + -- FF = 1.336E-15 0 -- FF = 1.336E-14

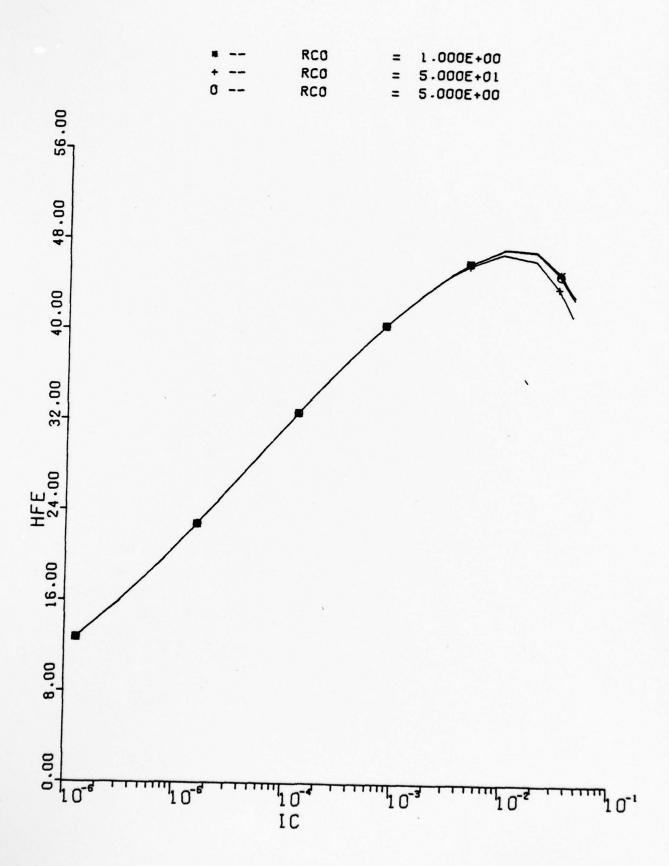


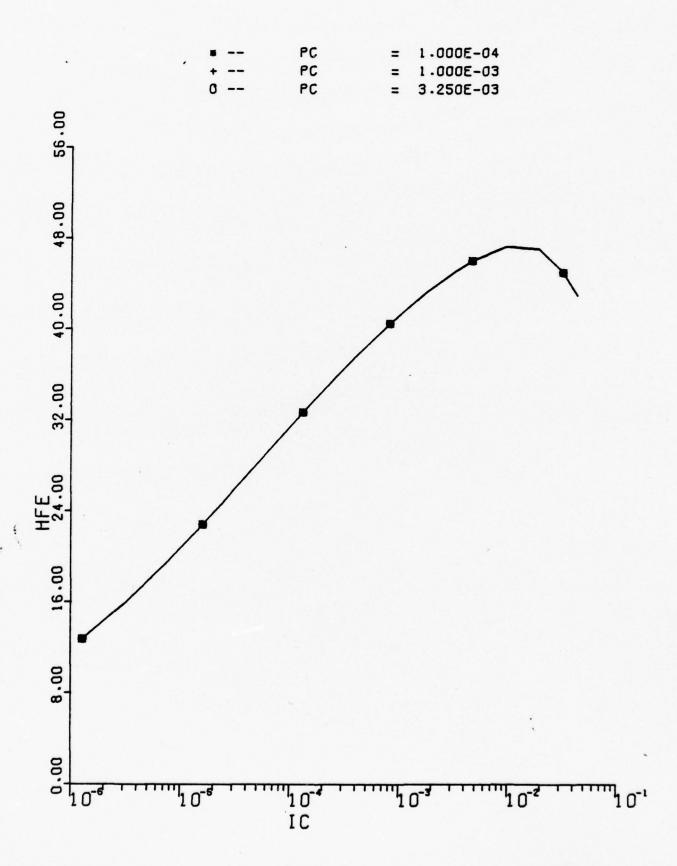
B.3 HFE vs. I Curves

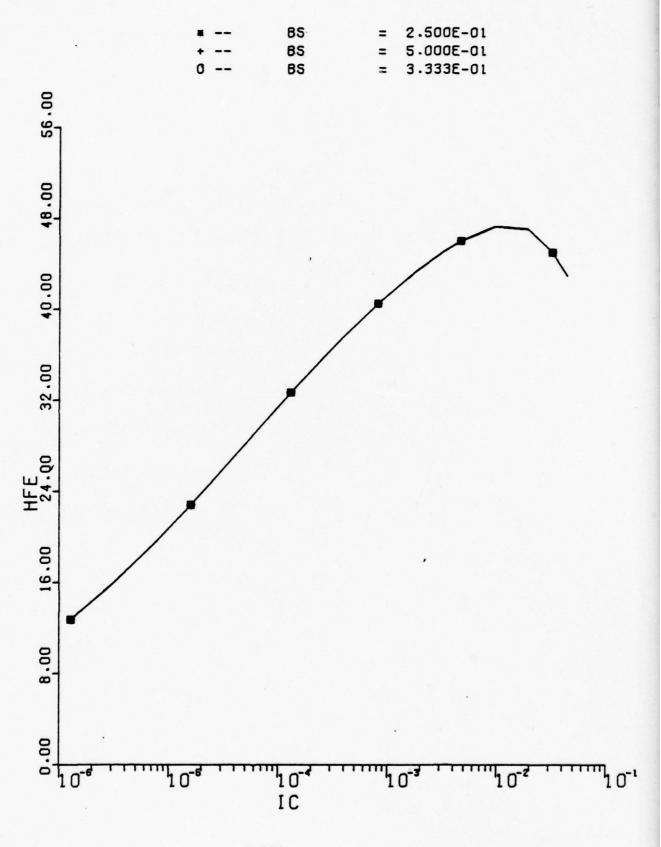










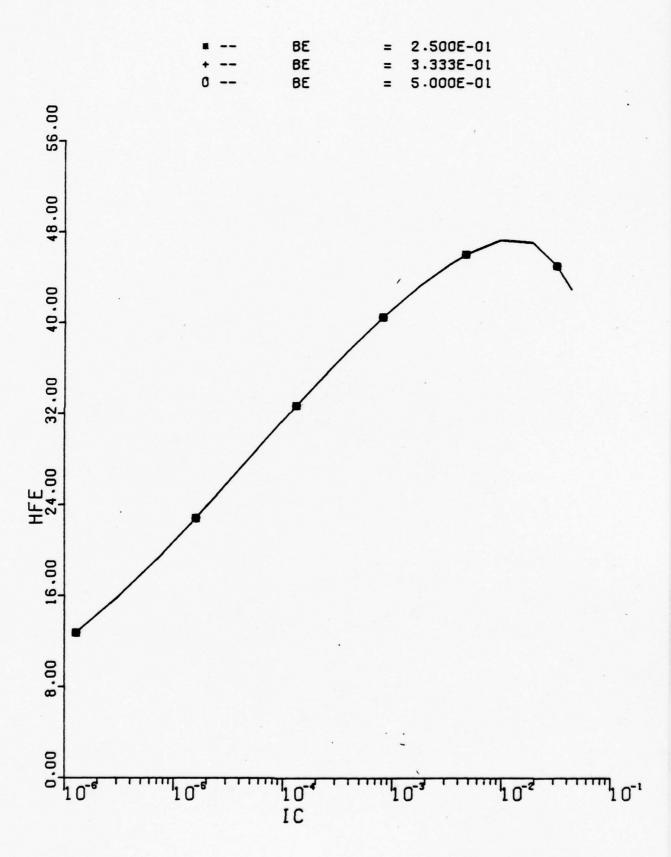


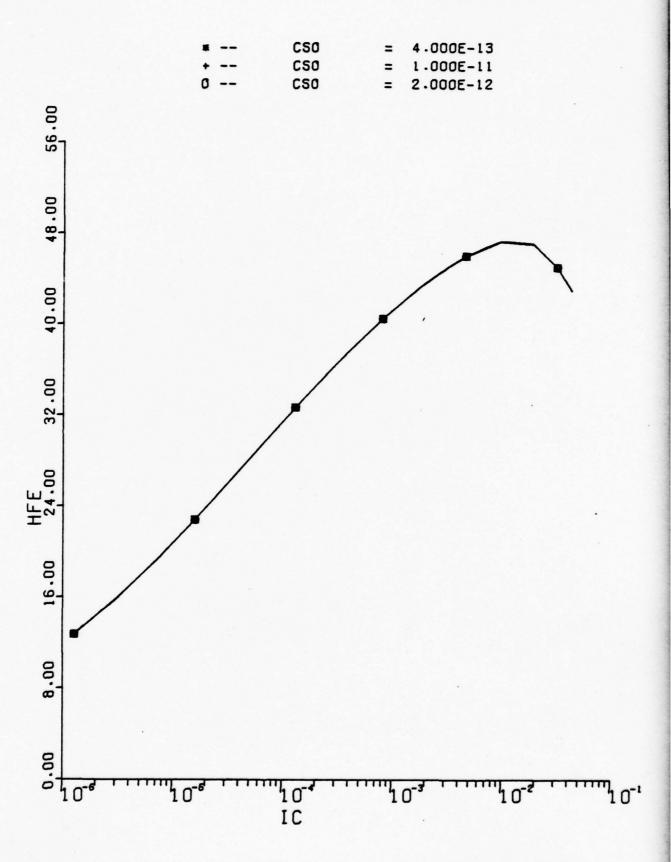
BC BC BC 5.000E-01 3.333E-01 56.00 48.00 40.00 32.00 HFE 24.00 16.00 8.00

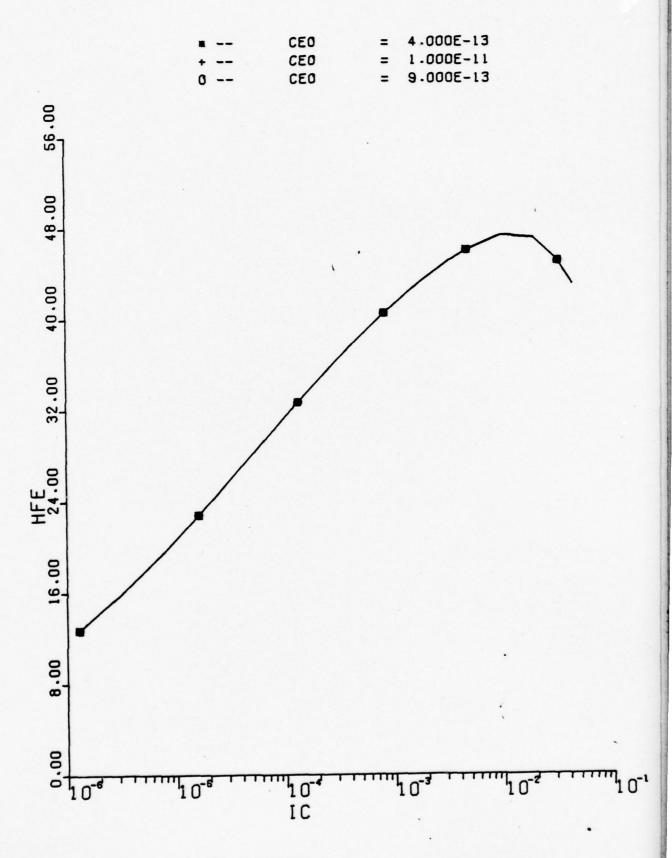
2.500E-01

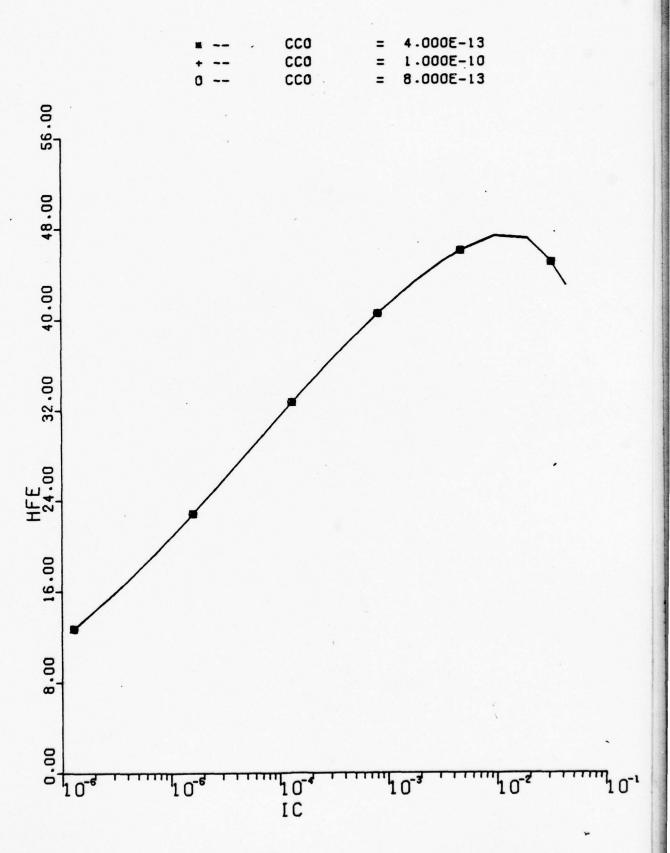
10 1C

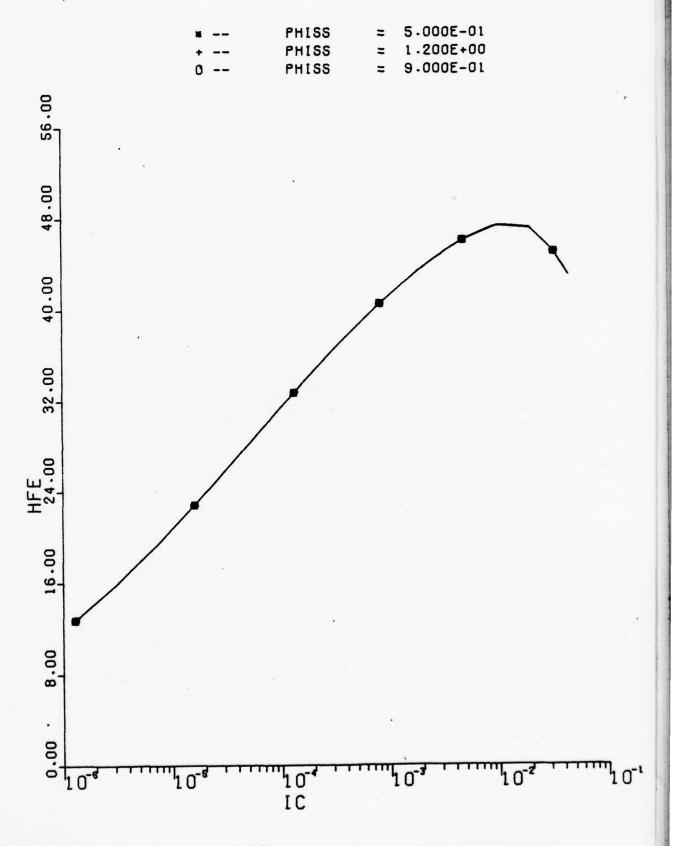
8 10-e

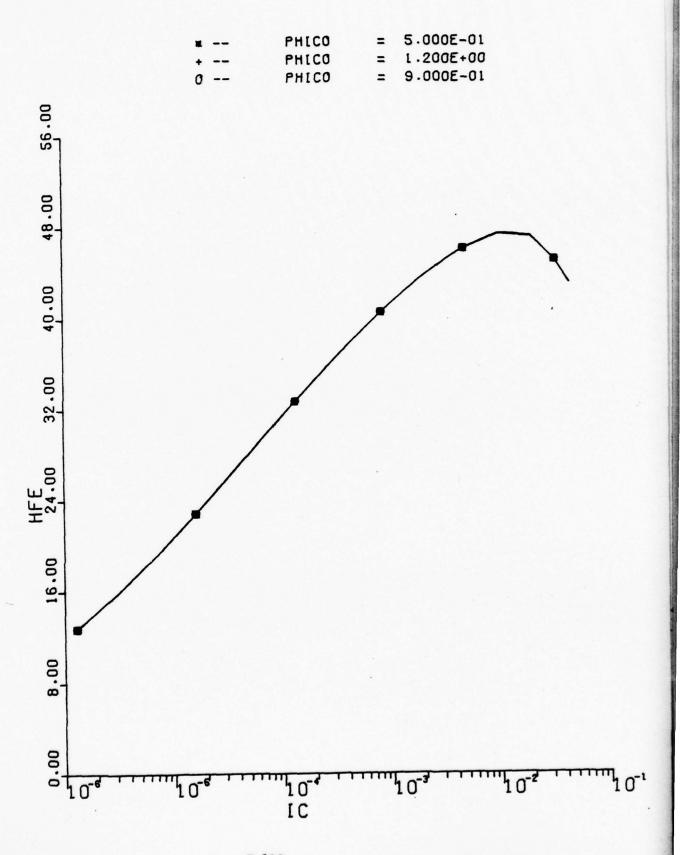


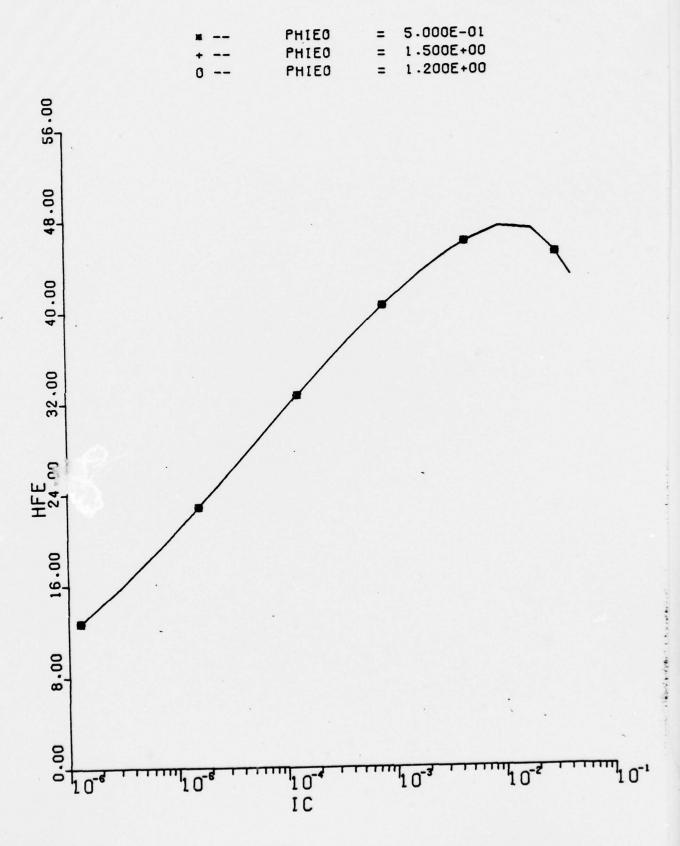


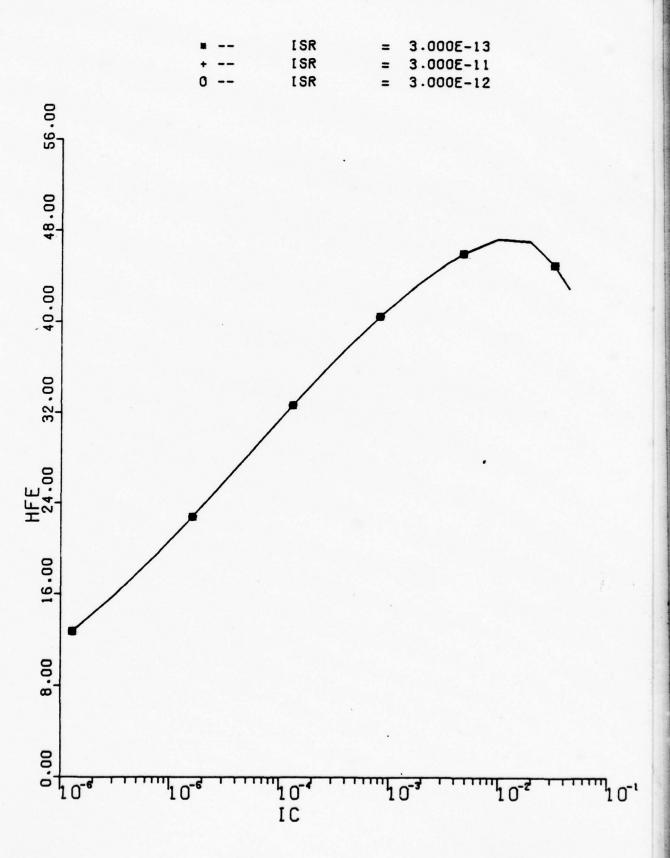


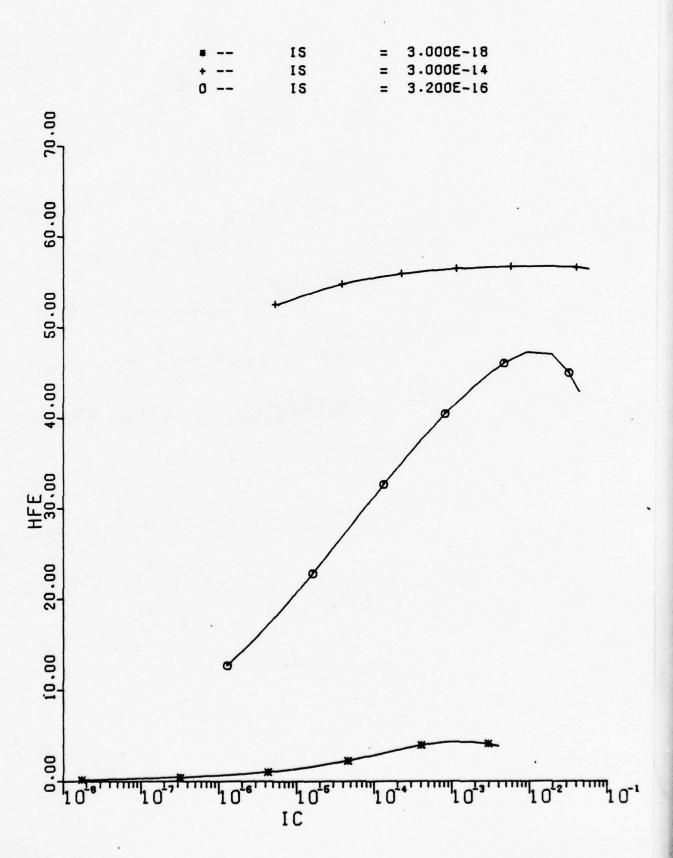


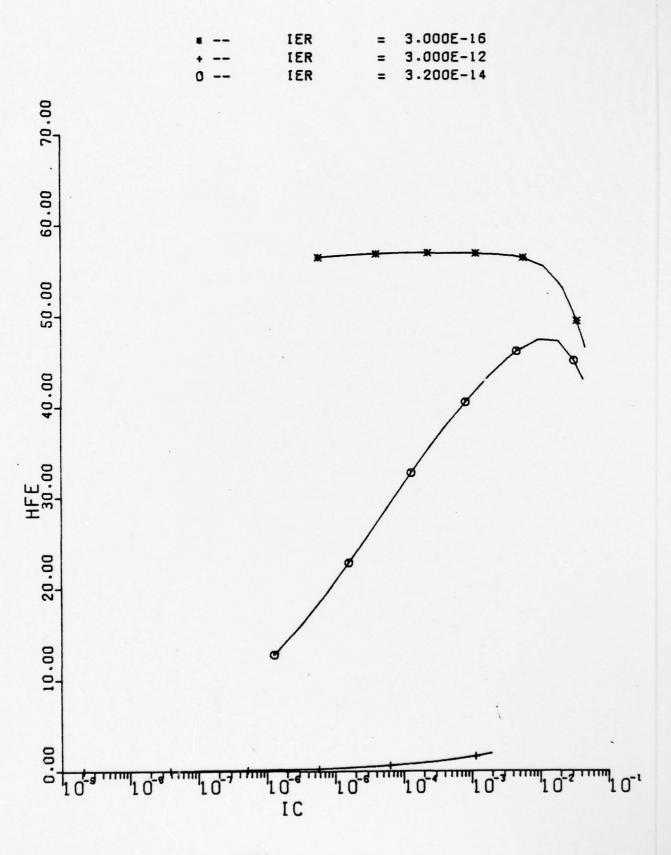


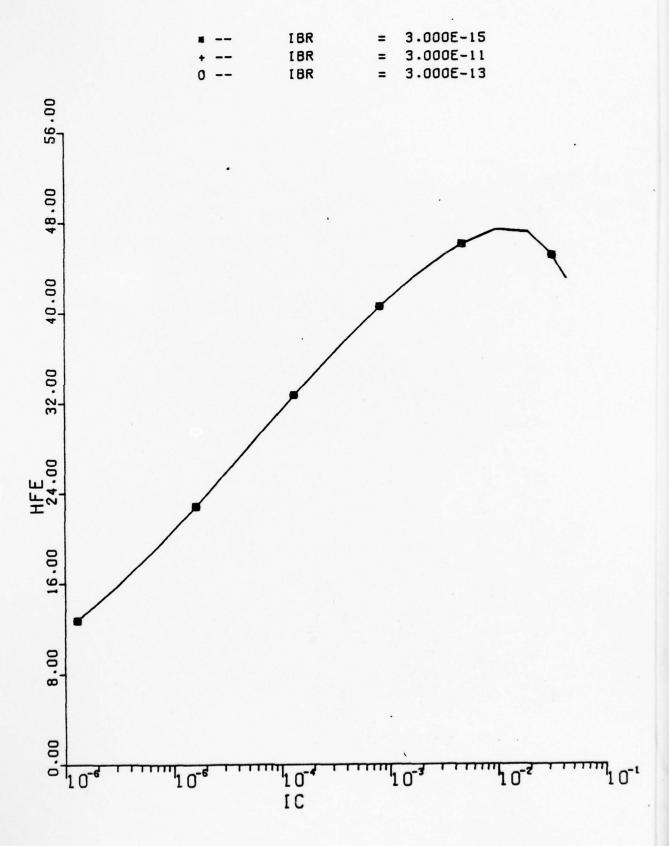


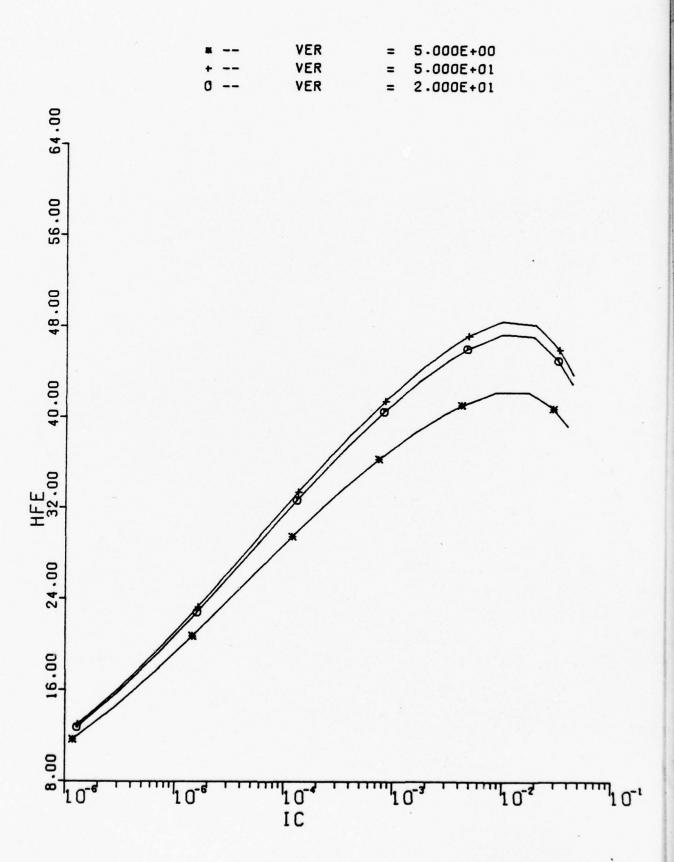


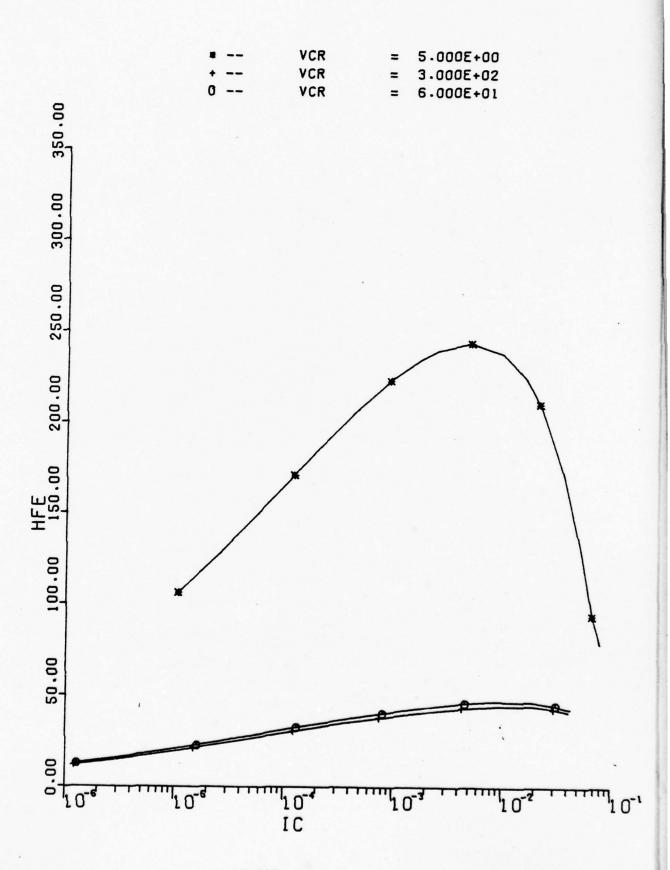


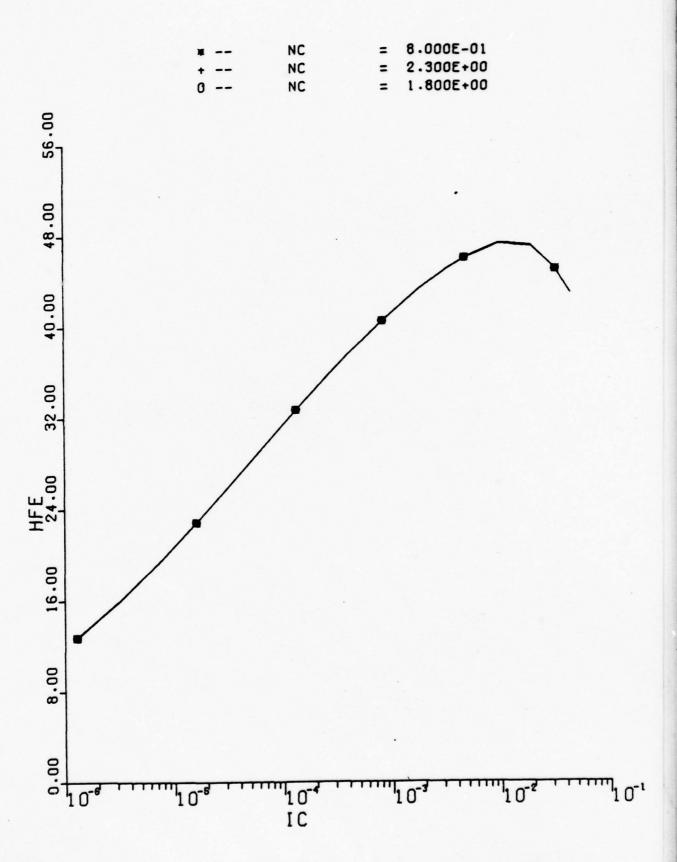


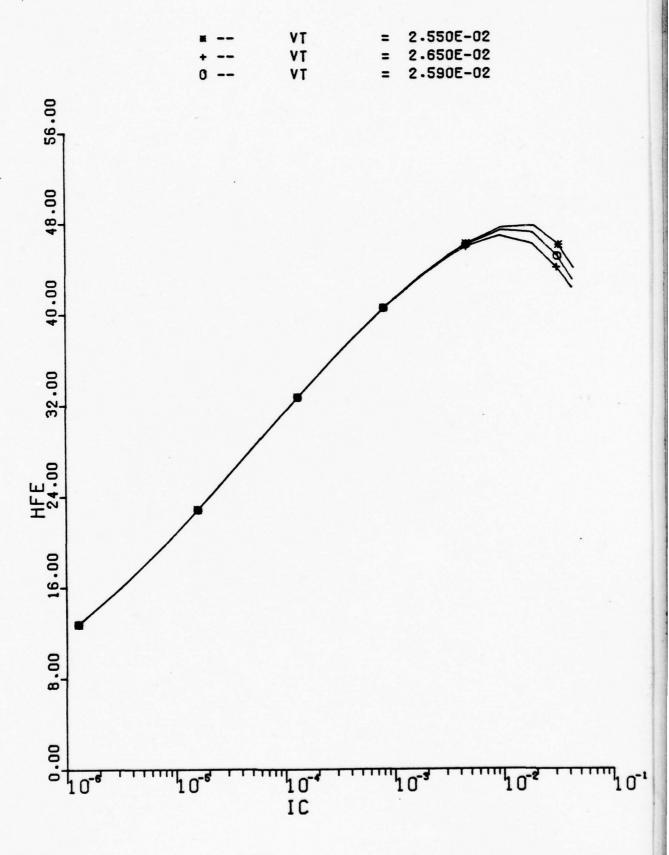


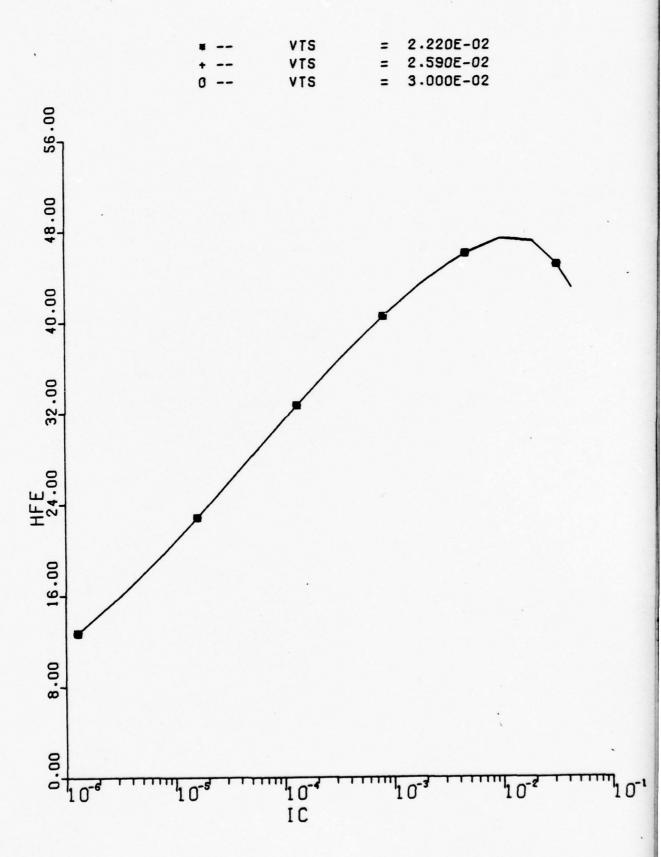


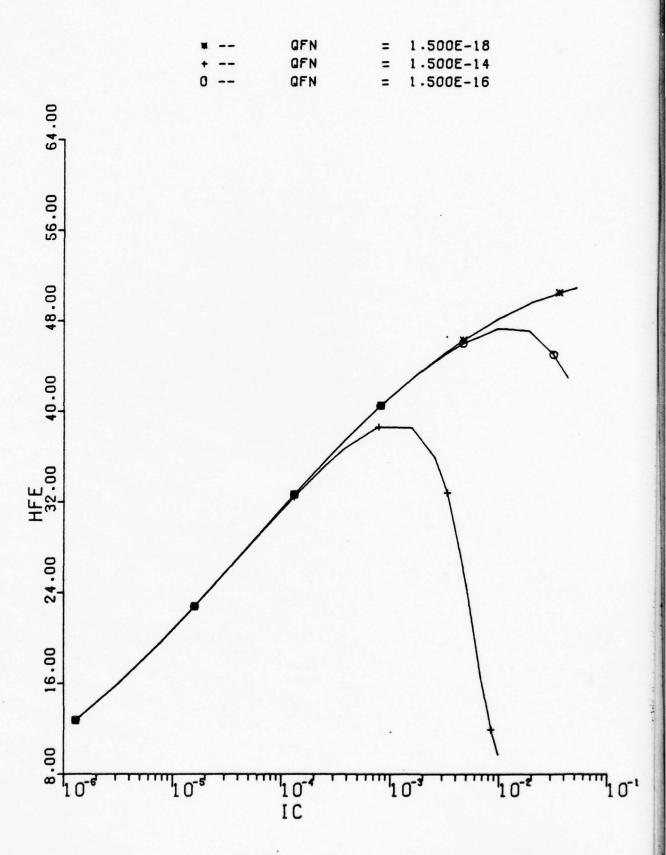


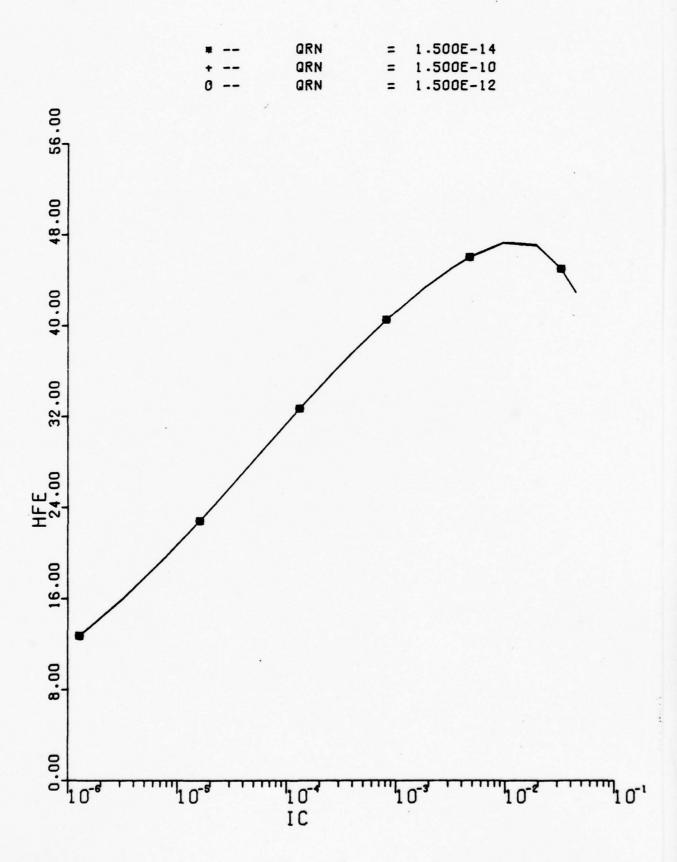


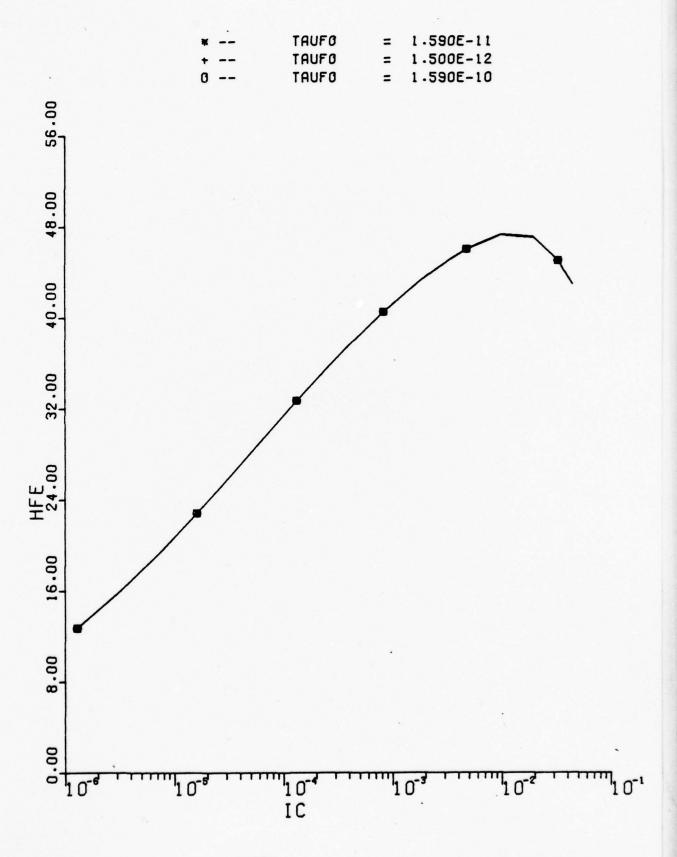












TAUR = 1.000E-06 +-- TAUR = 1.270E-11 0-- TAUR = 1.270E-09



10

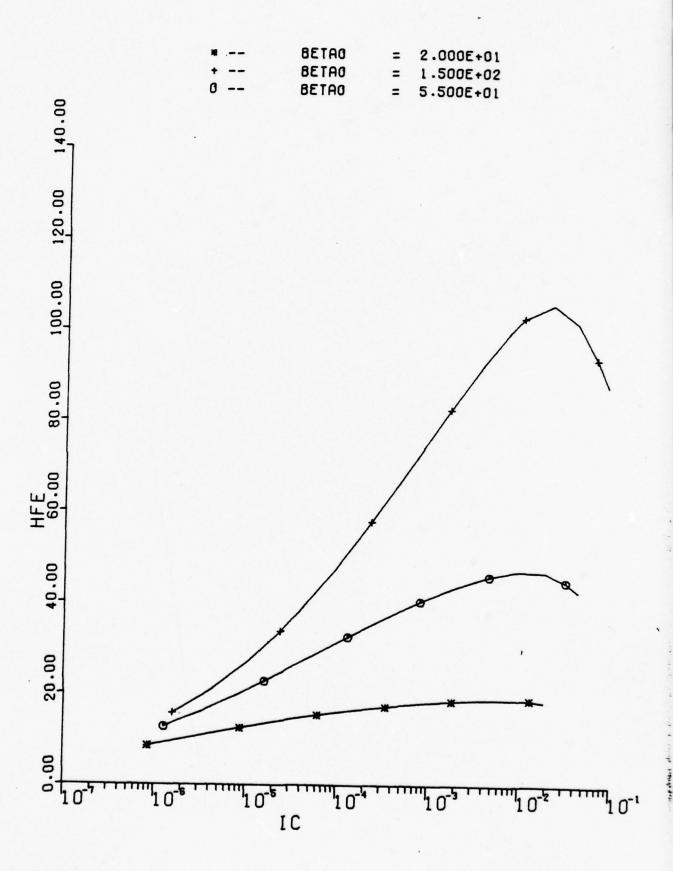
32.00

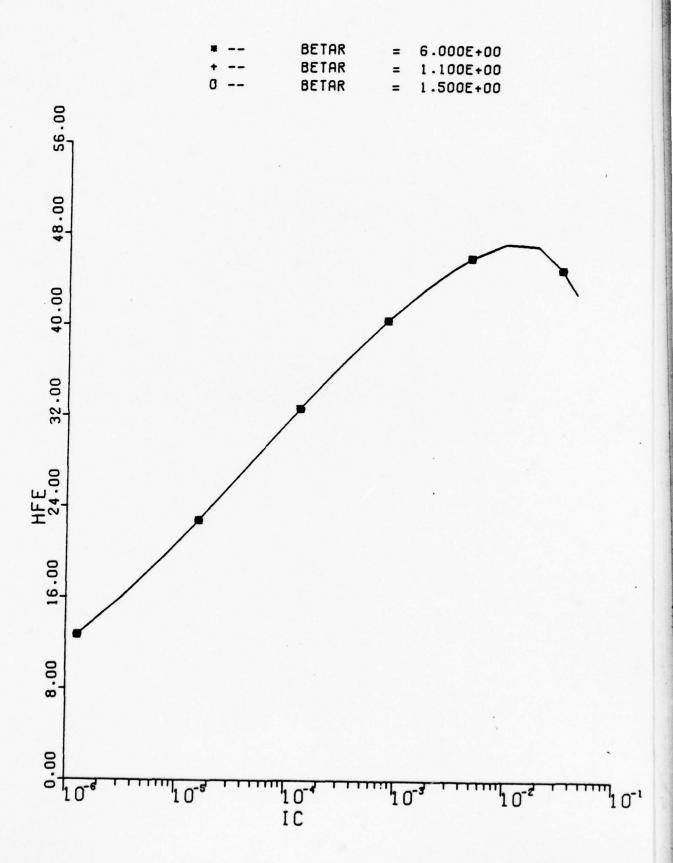
HFE 24.00

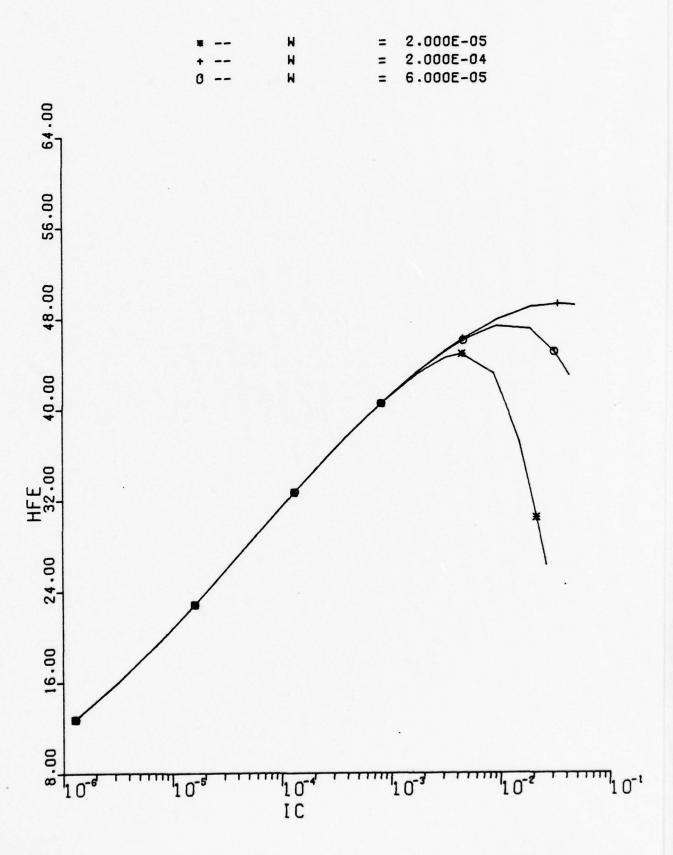
16.00

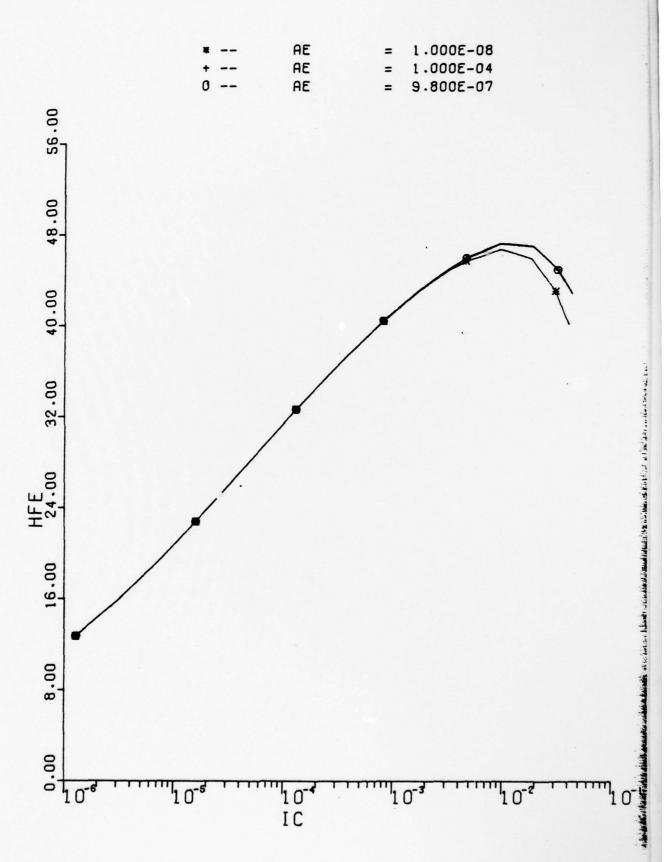
8.00

0.00

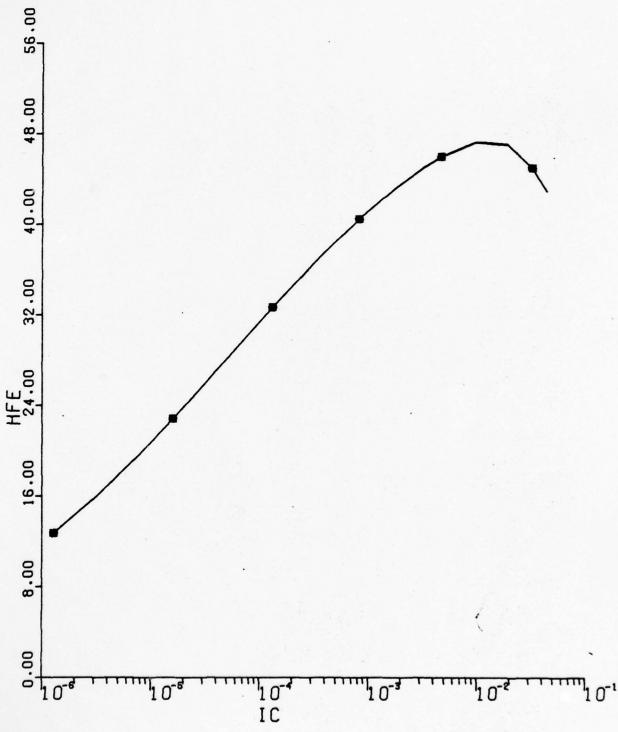


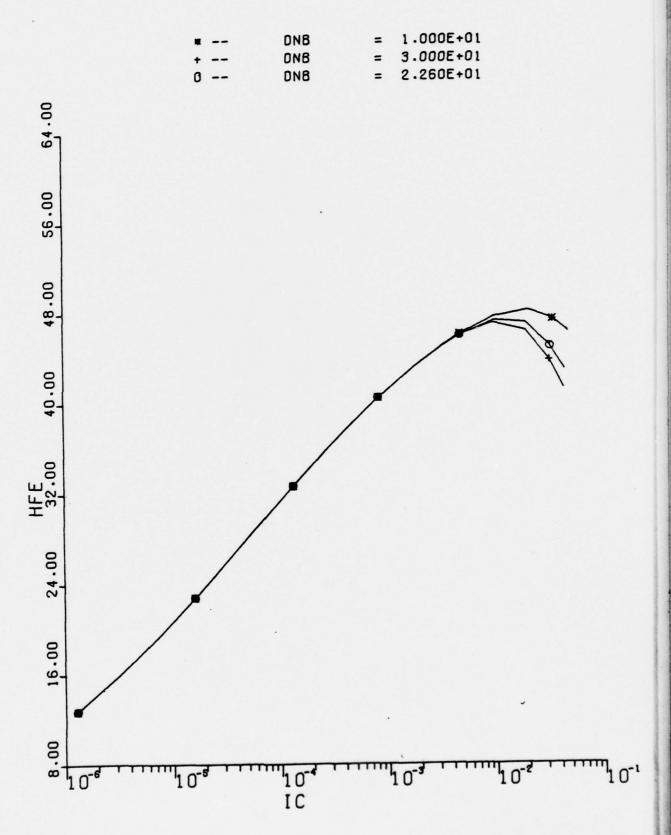


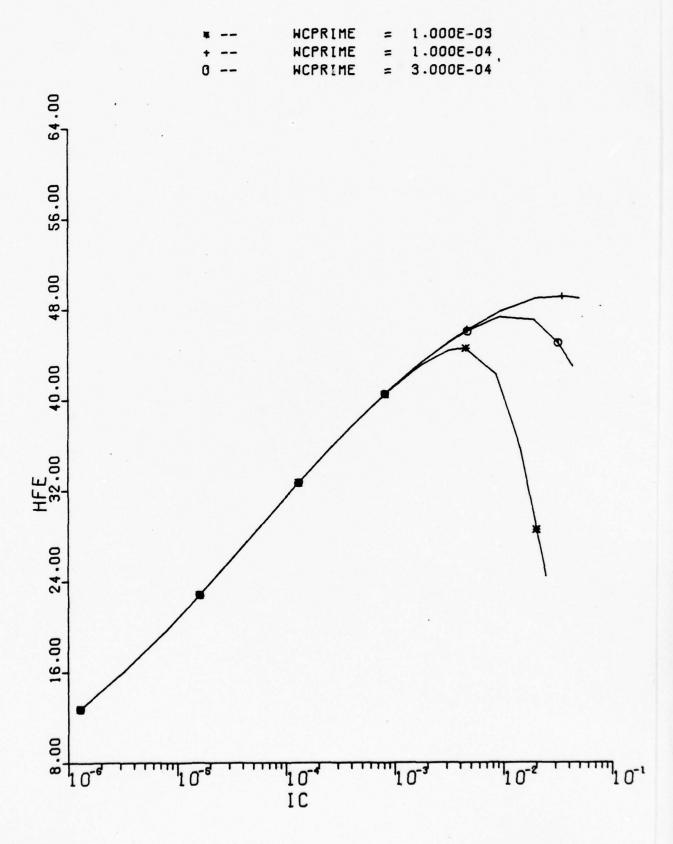


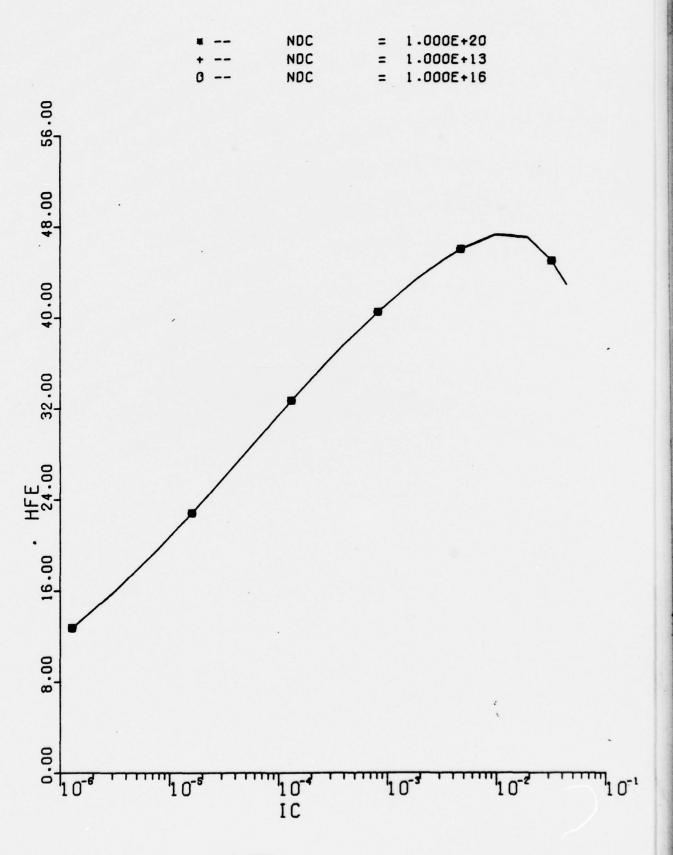


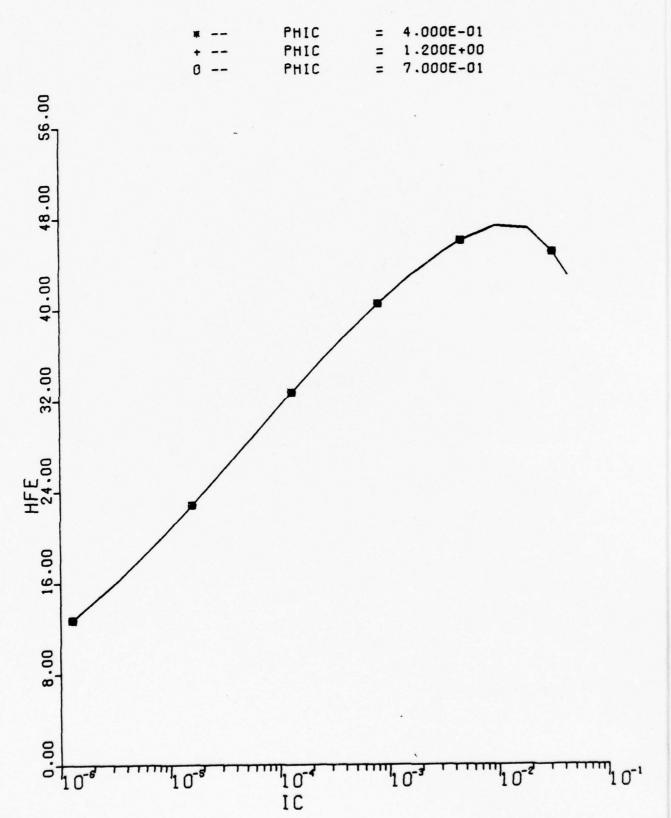


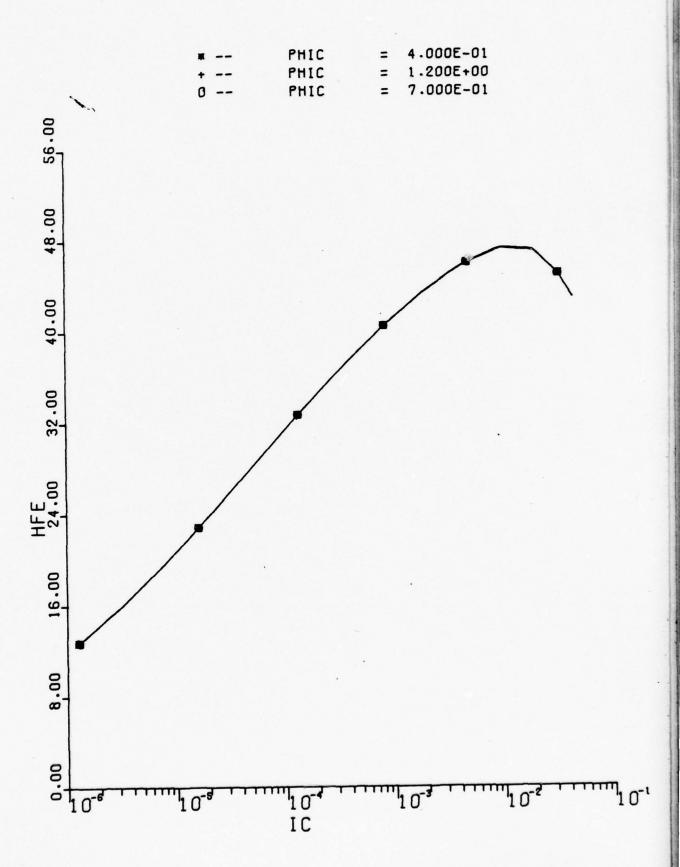


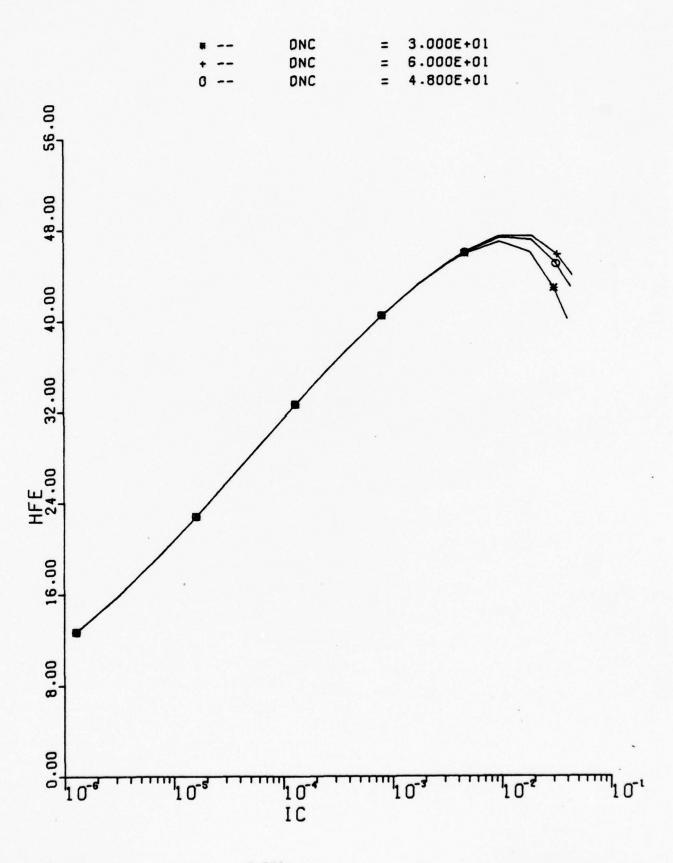


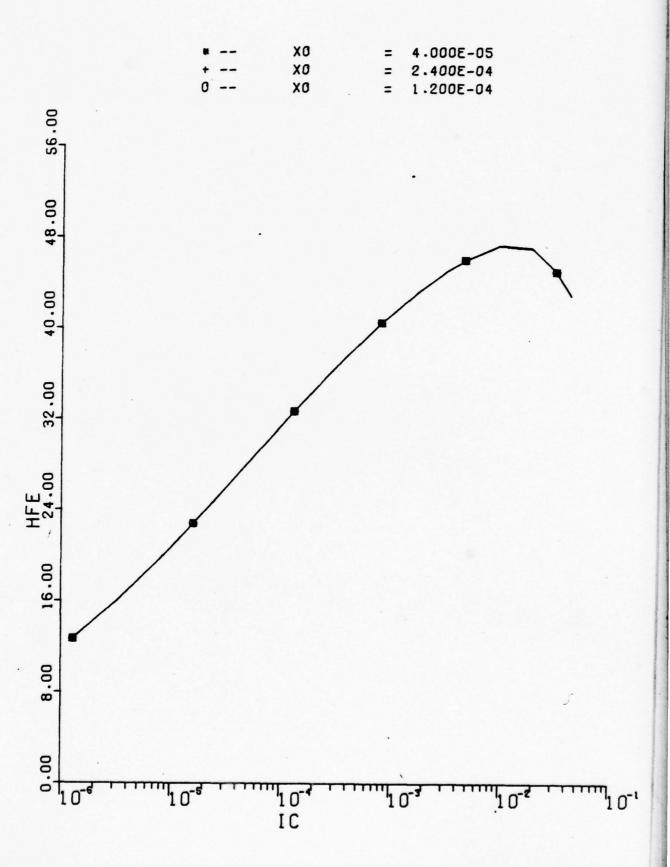


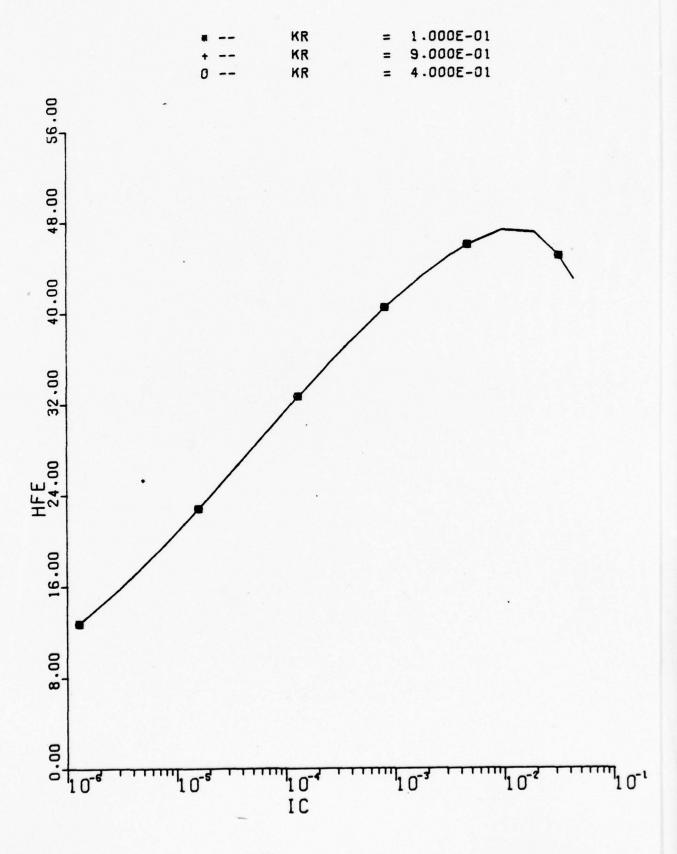




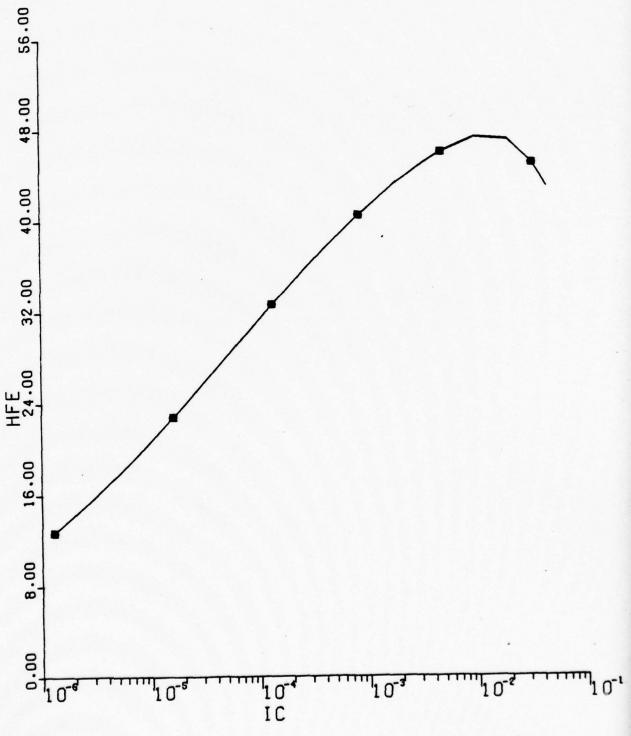


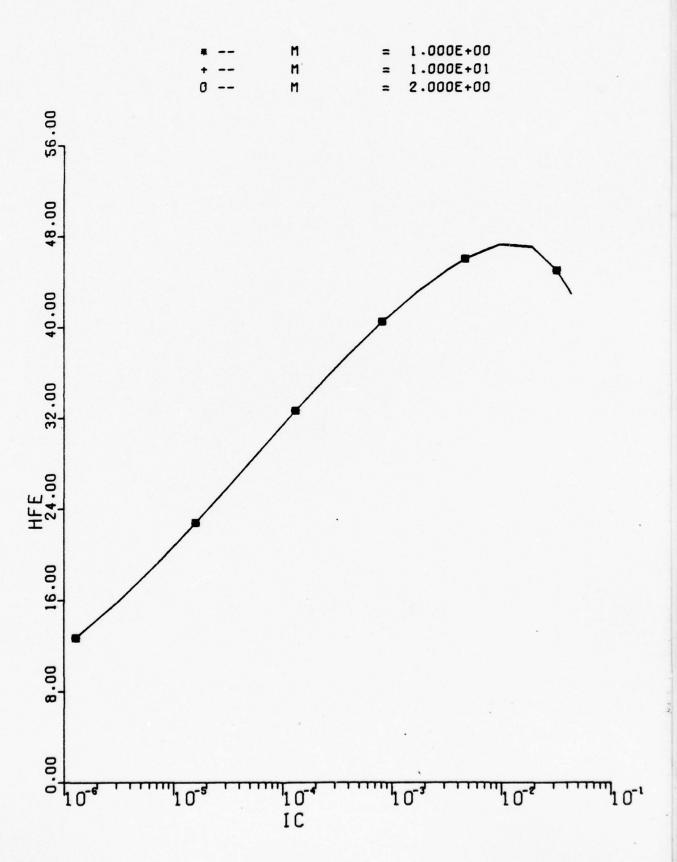


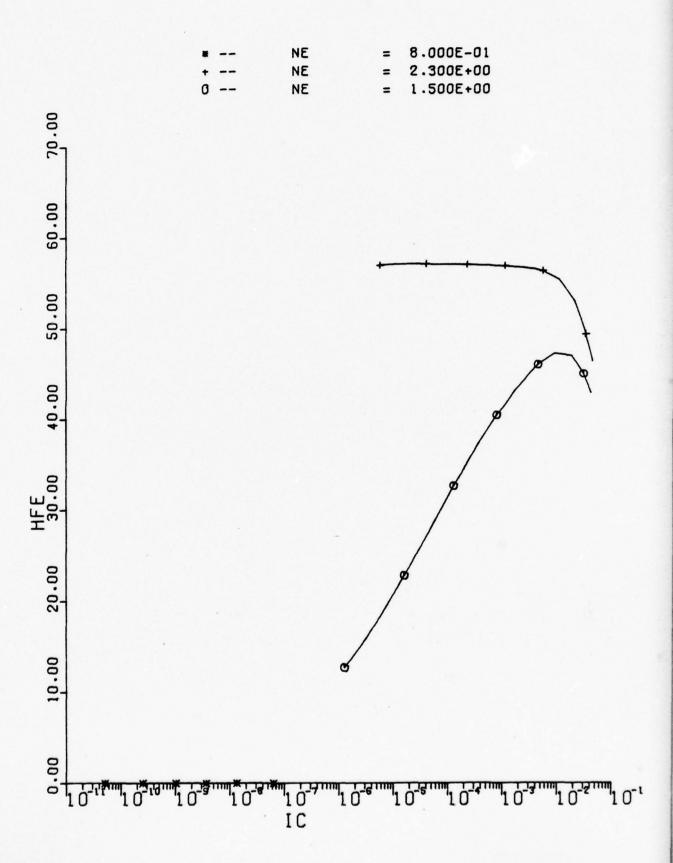


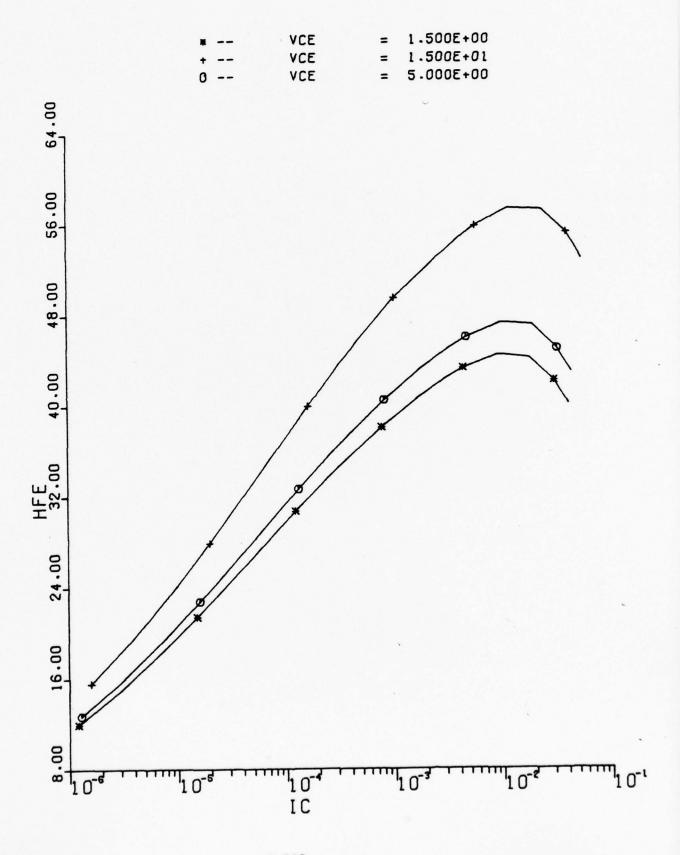


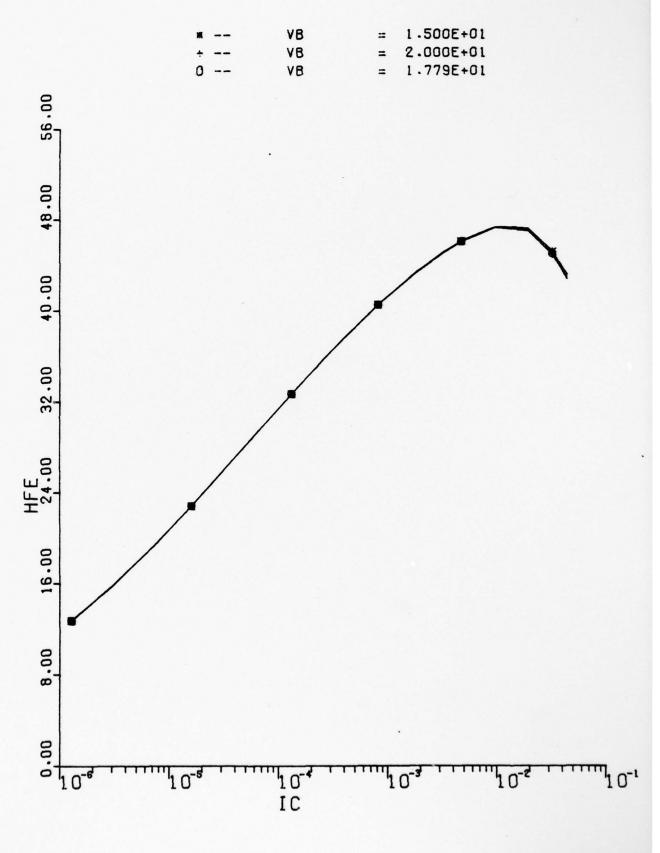


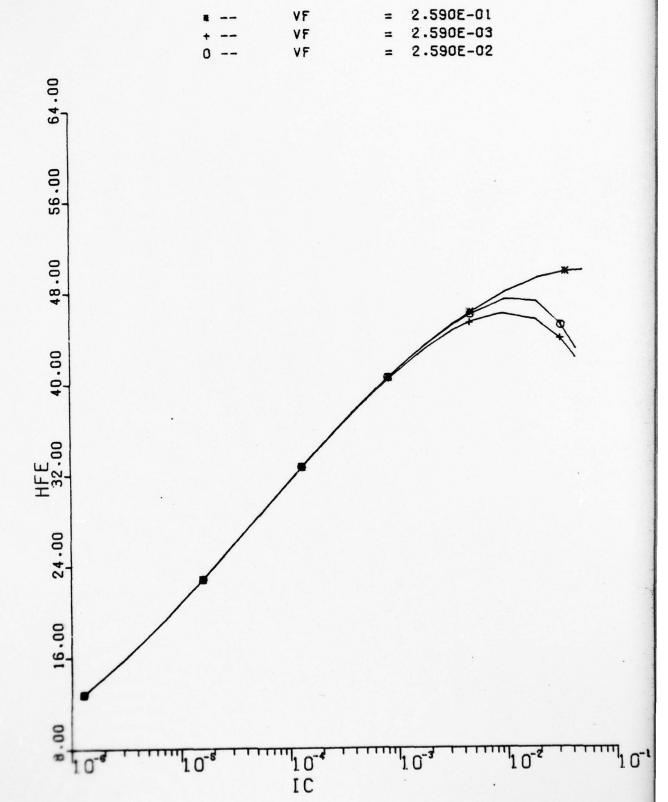




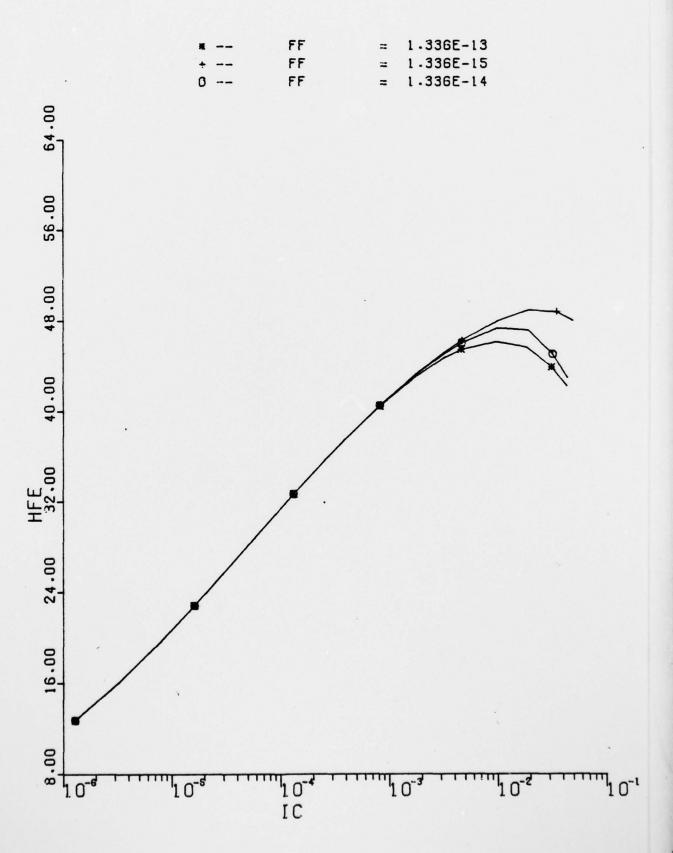






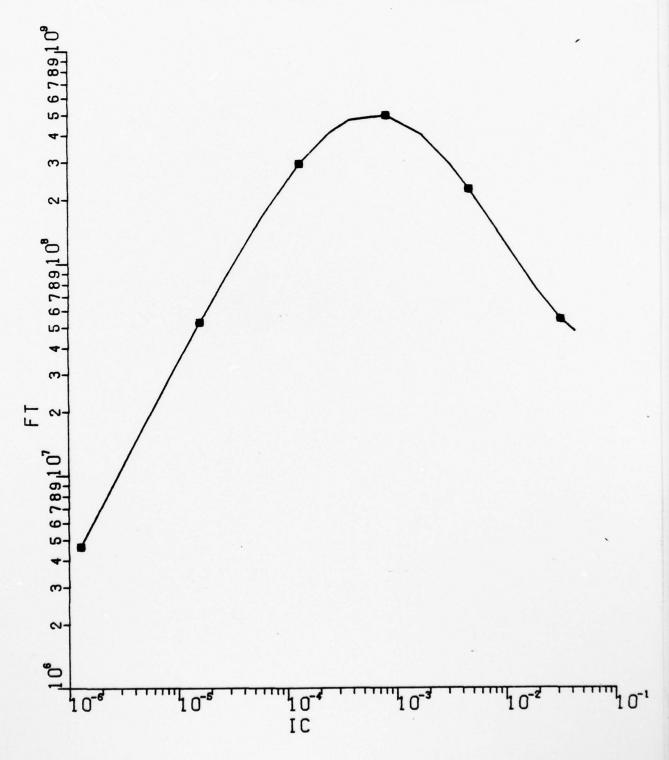


10 1C

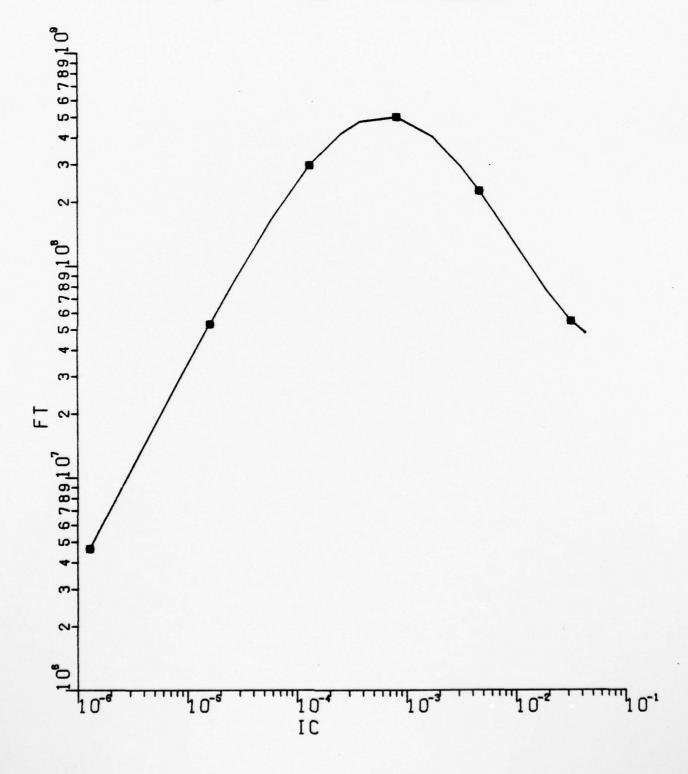


B.4 F<sub>T</sub> vs. I<sub>c</sub> Curves

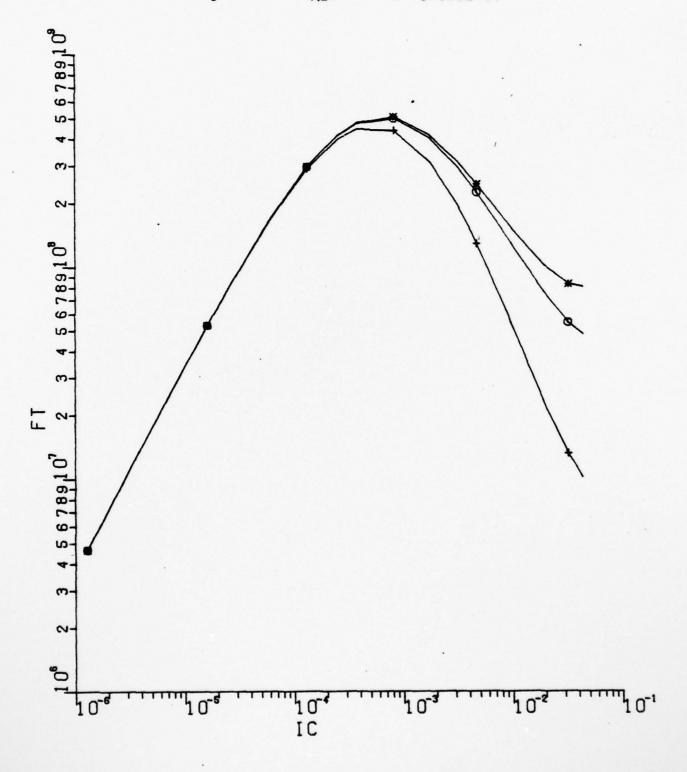




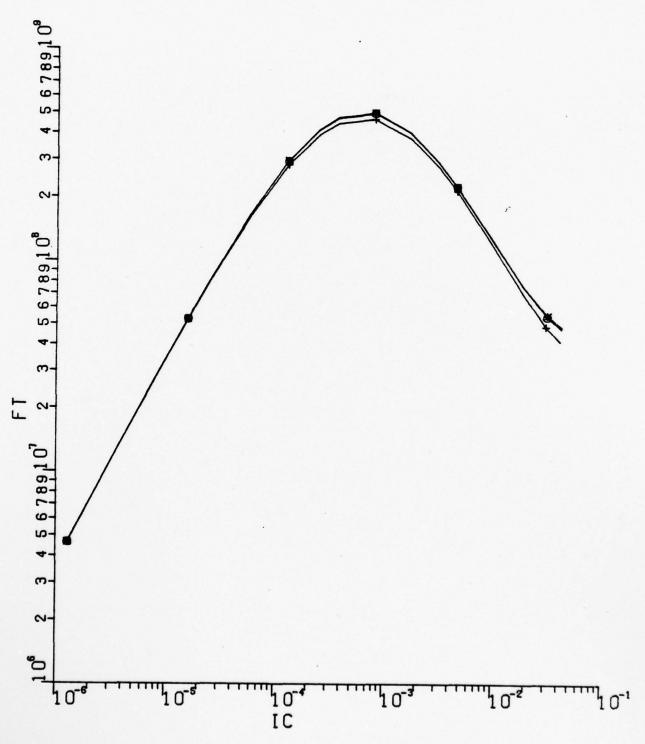
\* -- RBB = 1.500E+01 + -- RBB = 5.000E+02 0 -- RBB = 1.500E+02



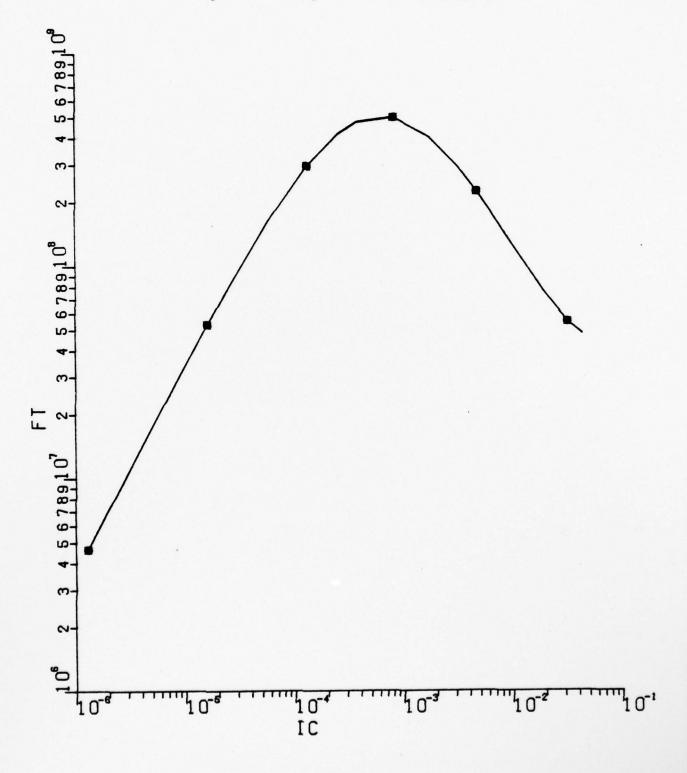
\* -- RE\* = 1.000E-02 + -- RE\* = 5.000E+00 0 -- RE\* = 5.000E-01



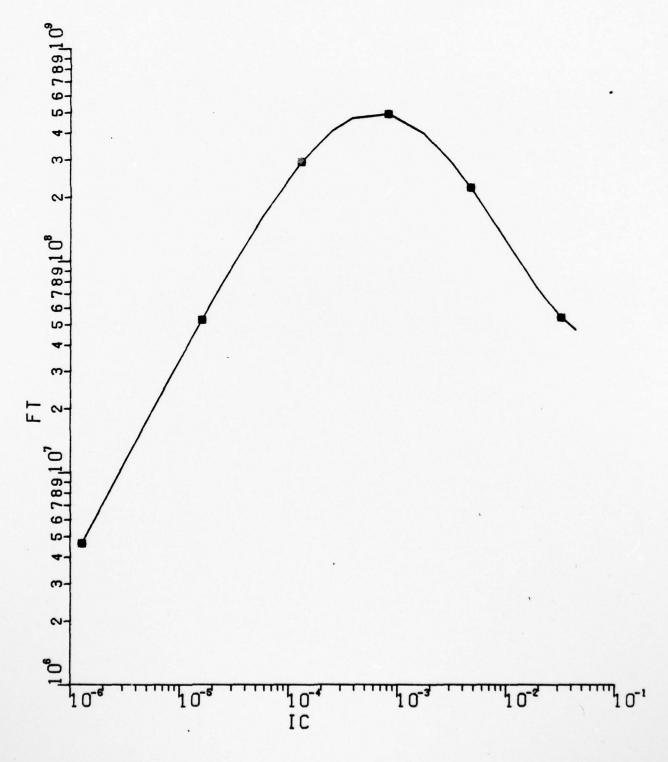




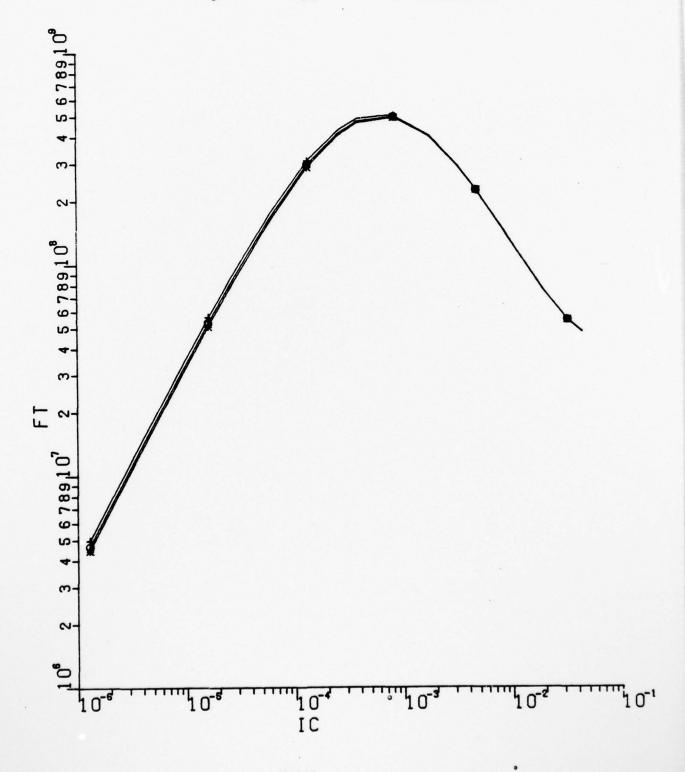
# -- PC = 6.000E-03
+ -- PC = 1.000E-03
0 -- PC = 3.250E-03



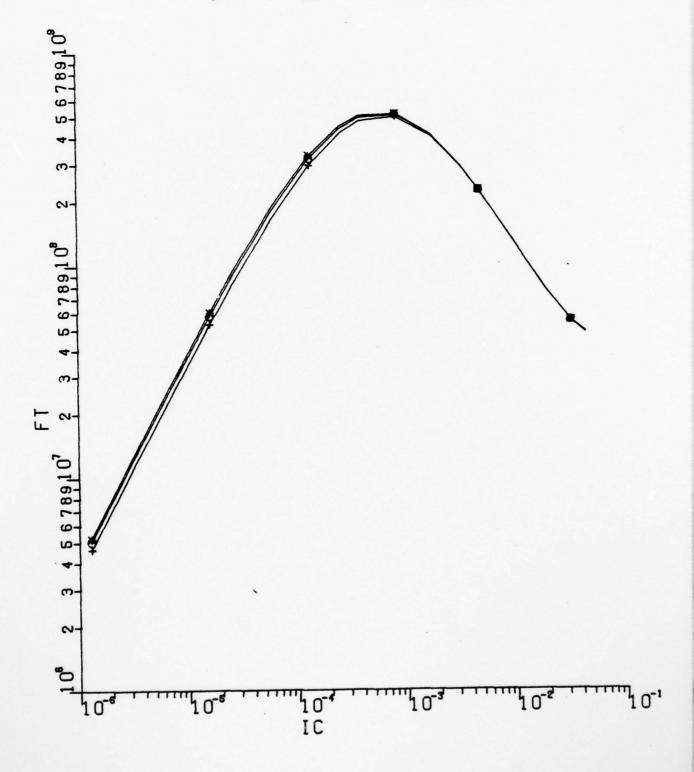




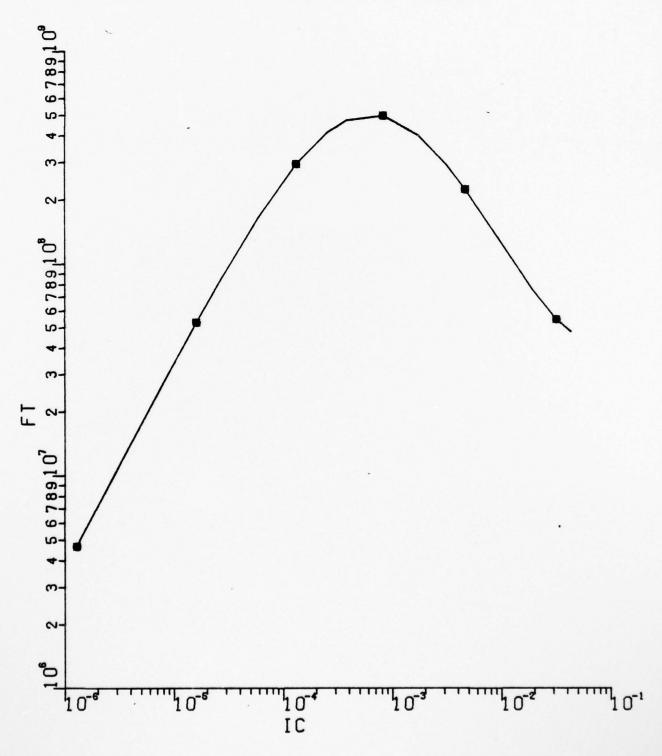
# -- BC = 2.500E-01 + -- BC = 5.000E-01 0 -- BC = 3.333E-01



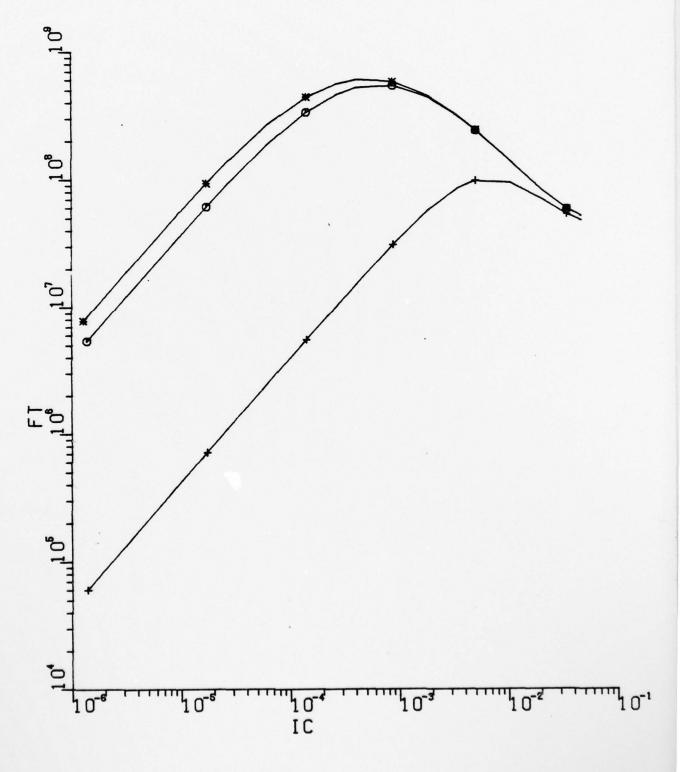
# -- BE = 2.500E-01 + -- BE = 5.000E-01 · 0 -- BE = 3.333E-01



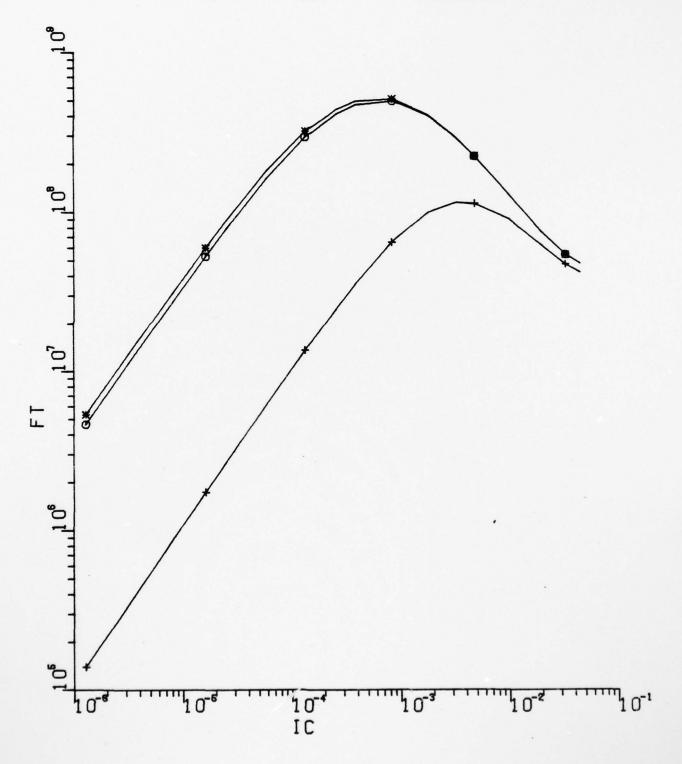


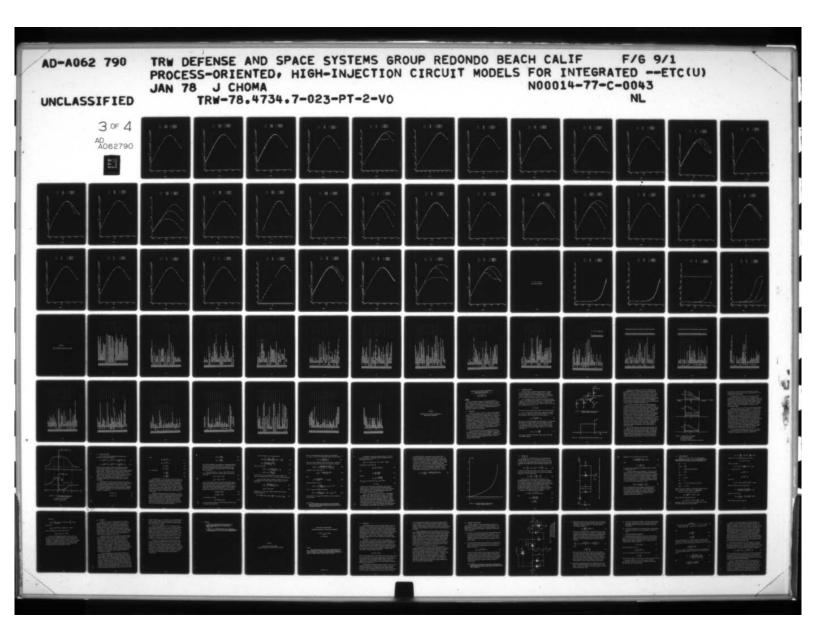


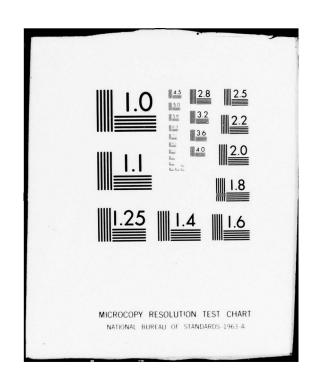
\* -- CEO = 4.000E-13 + -- CEO = 1.000E-10 0 -- CEO = 8.000E-13



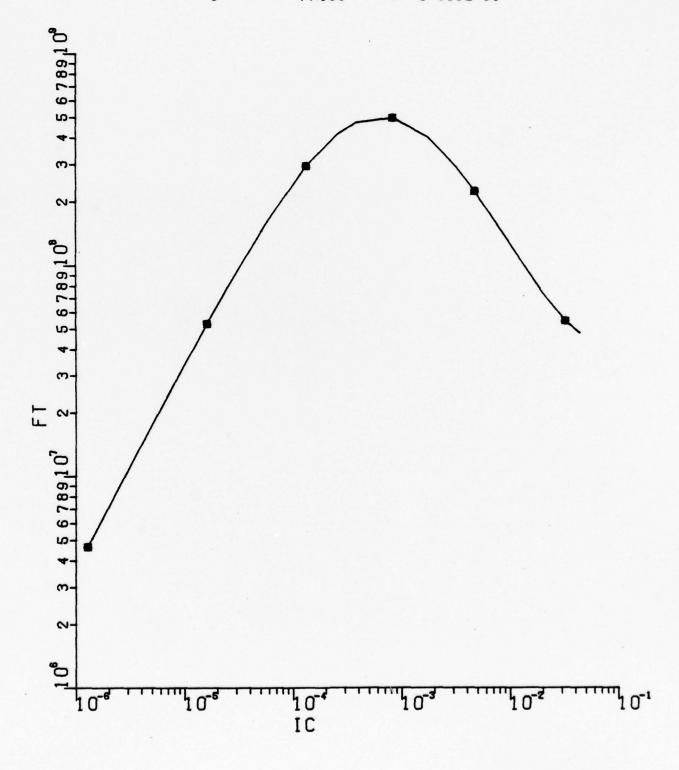
# -- CCO = 4.000E-13 + -- CCO = 1.000E-10 0 -- CCO = 8.000E-13



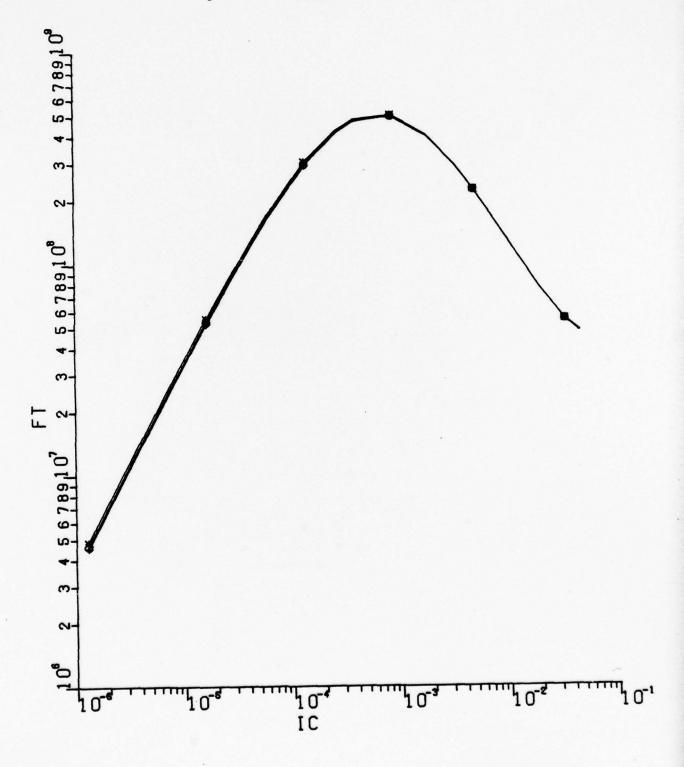




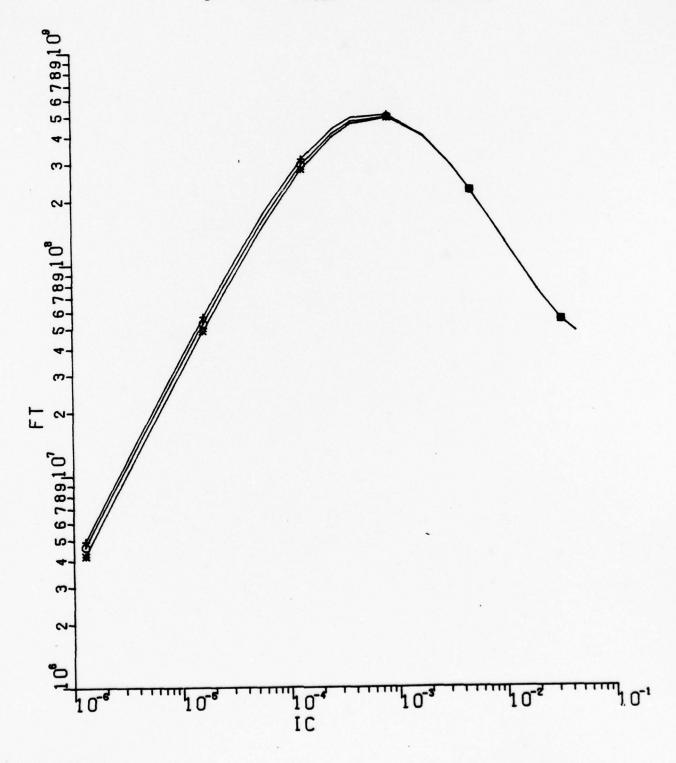
# -- PHISS = 5.000E-01 + -- PHISS = 1.200E+00 0 -- PHISS = 9.000E-01



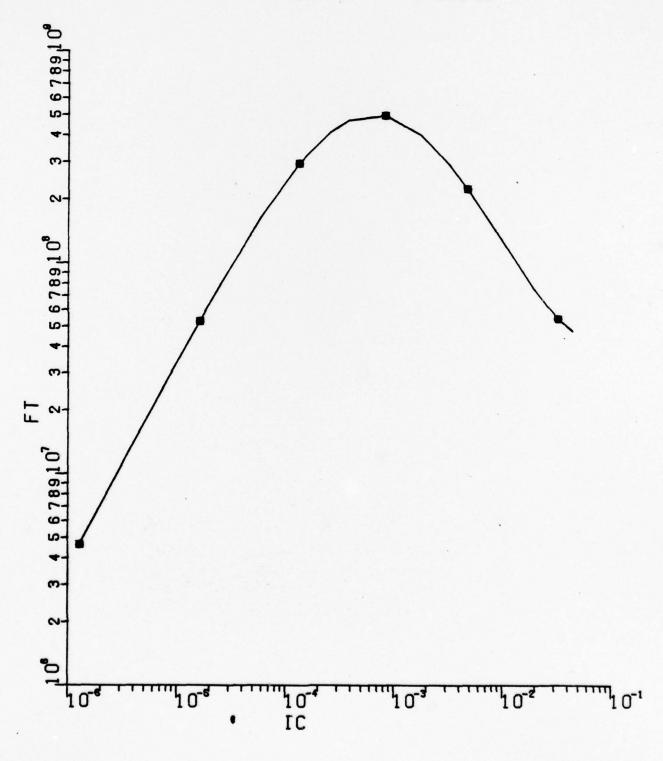
# -- PHICO = 5.000E-01 + -- PHICO = 1.200E+00 0 -- PHICO = 9.000E-01



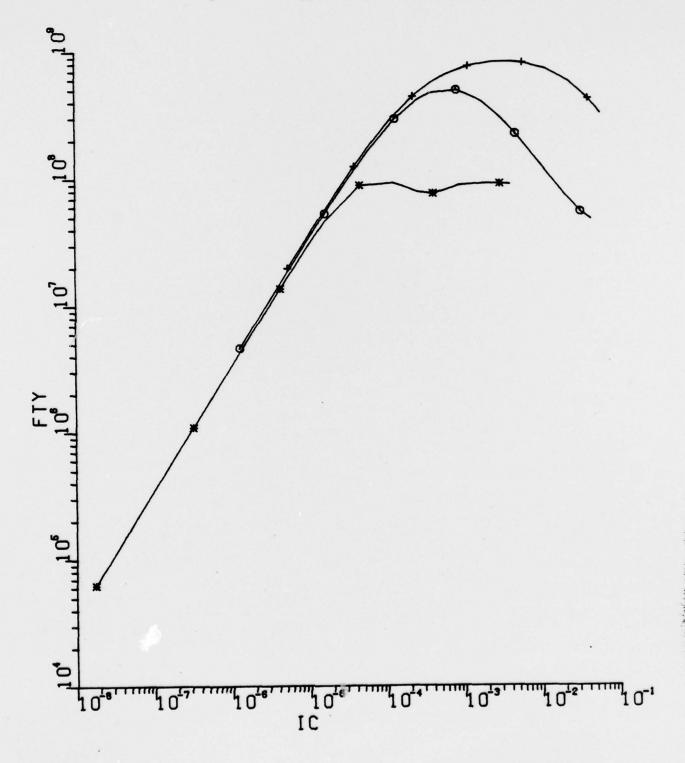
# -- PHIEO = 5.000E-01 + -- PHIEO = 1.500E+00 0 -- PHIEO = 1.200E+00



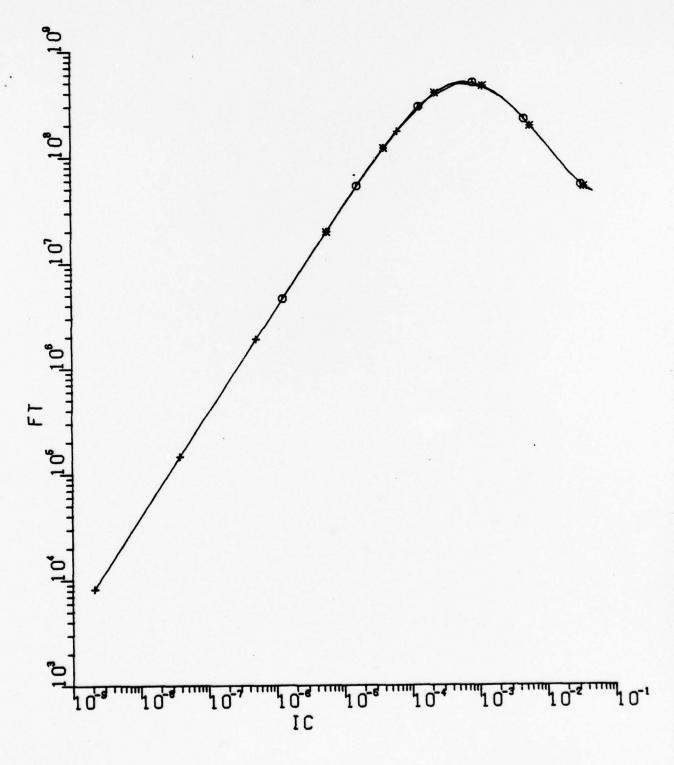
# -- ISR = 3.000E-13 + -- ISR = 3.000E-11 0 -- ISR = 3.000E-12



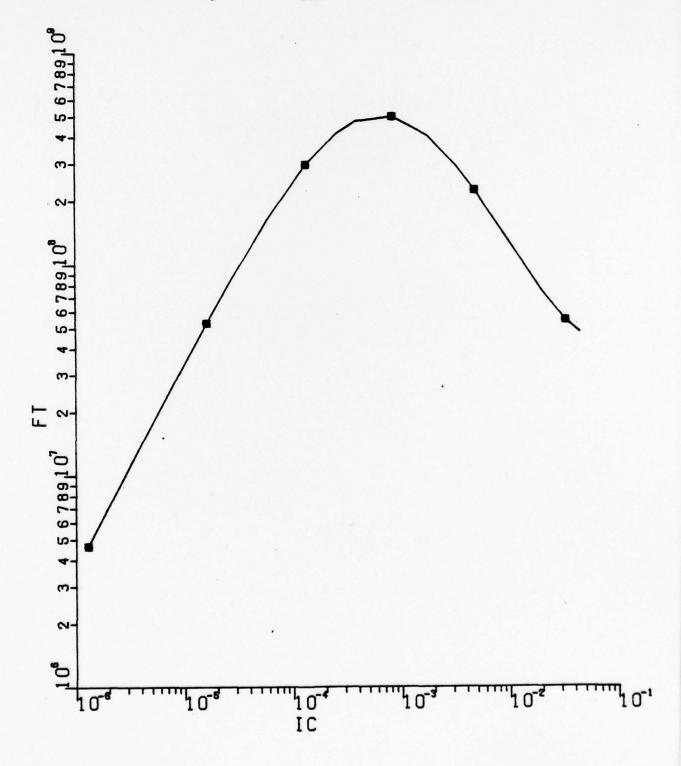
# -- IS = 3.000E-18 + -- IS = 3.000E-14 0 -- IS = 3.200E-16



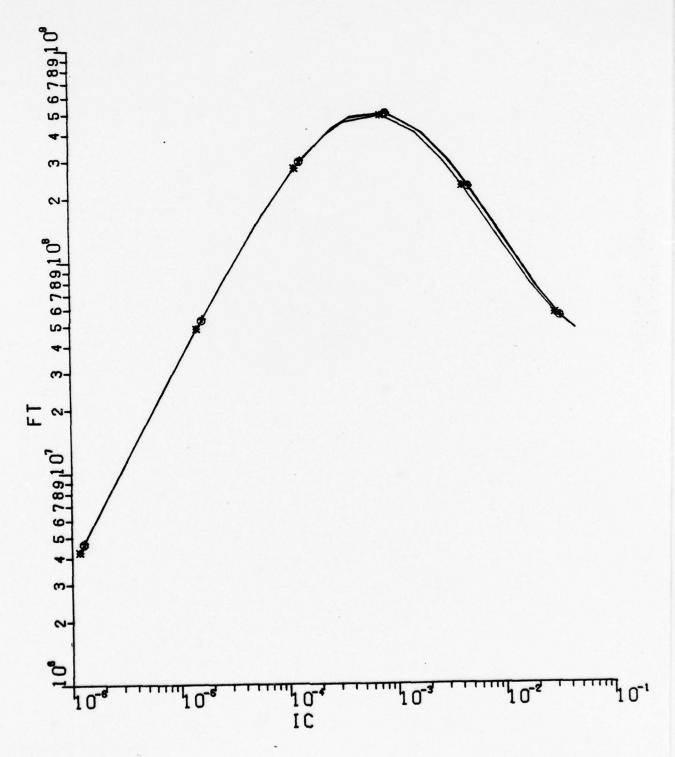


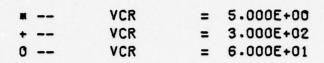


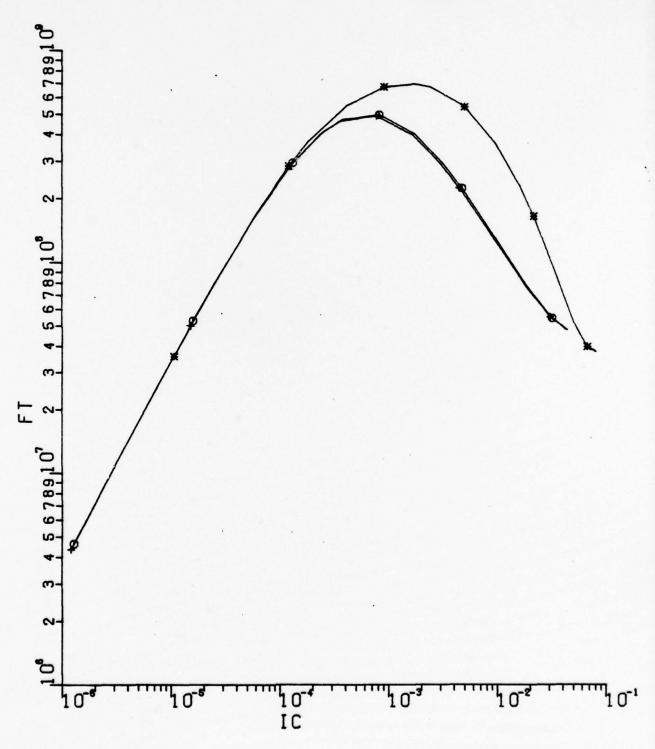
# -- IBR = 3.000E-15 + -- IBR = 3.000E-11 0 -- IBR = 3.000E-13



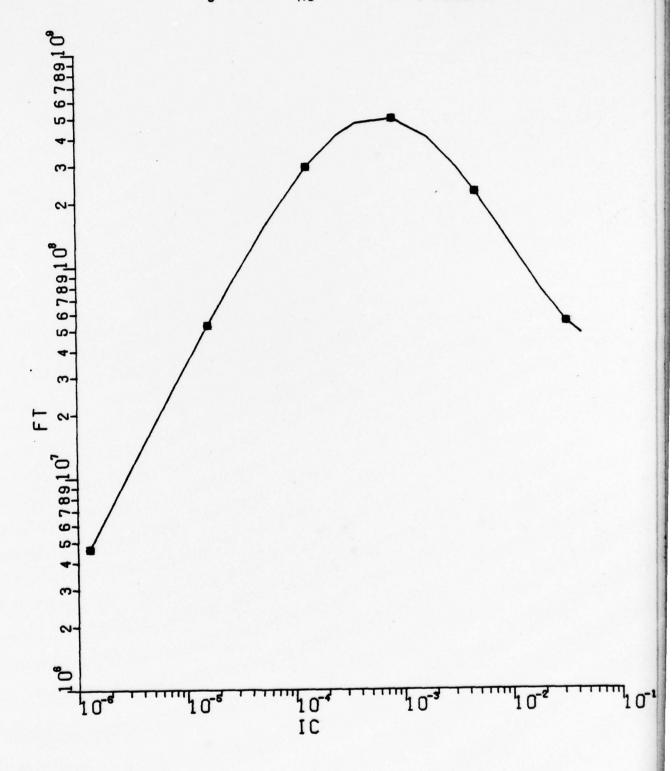


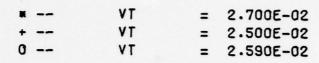


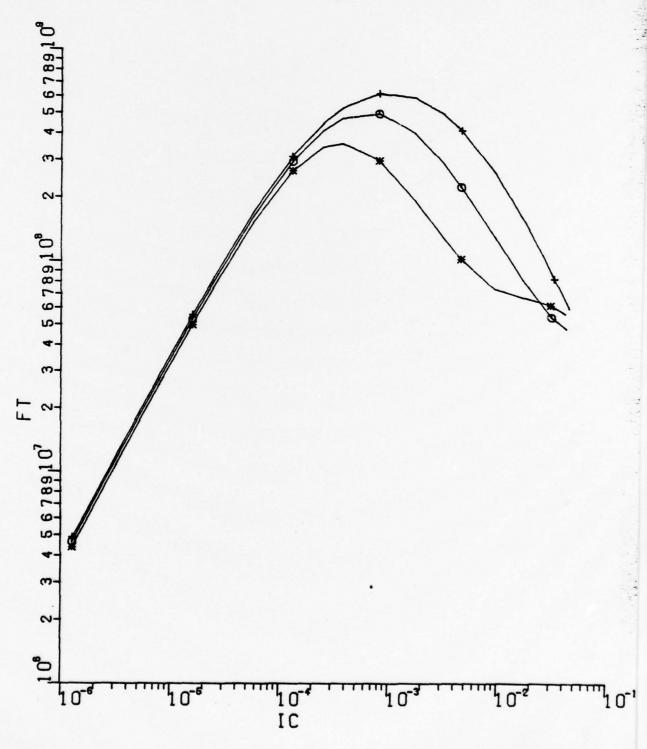




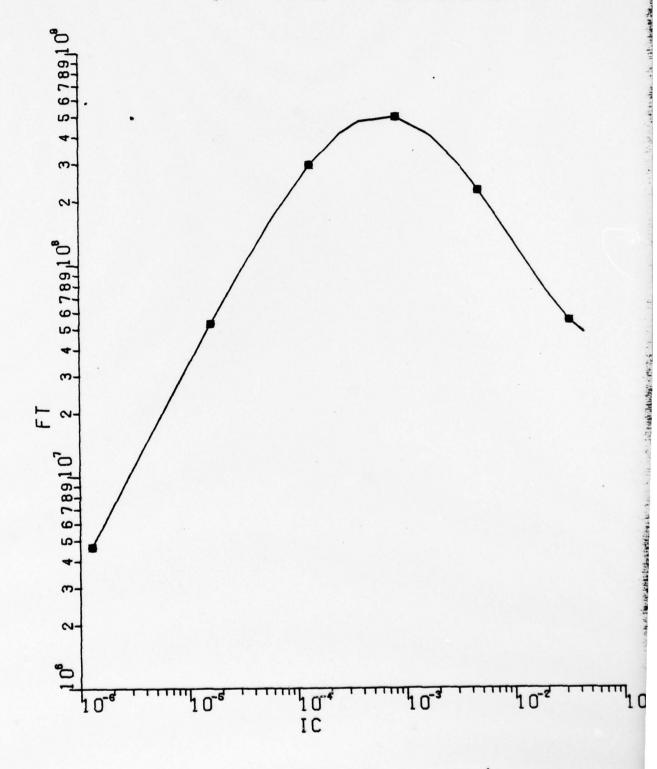
\* -- NC = 8.000E-01 + -- NC = 2.300E+00 0 -- NC = 1.800E+00



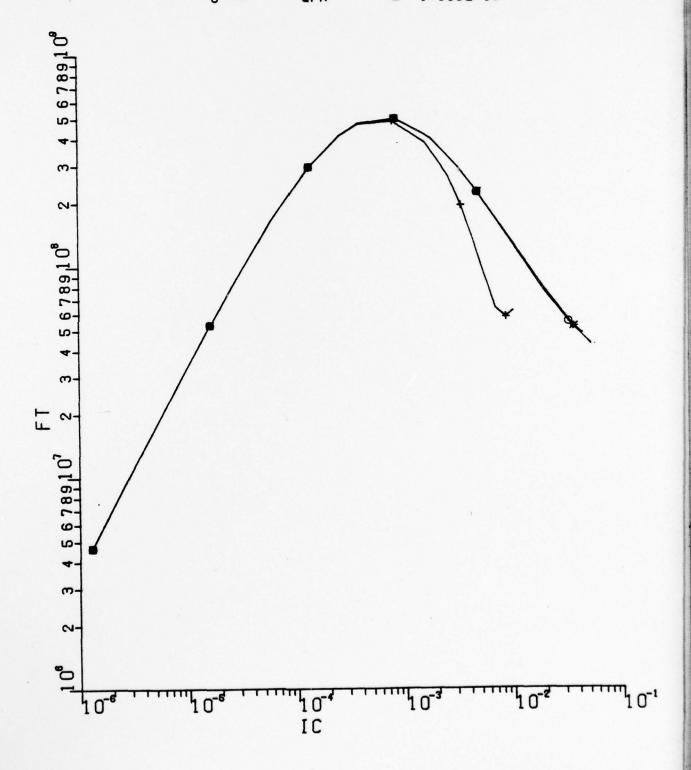




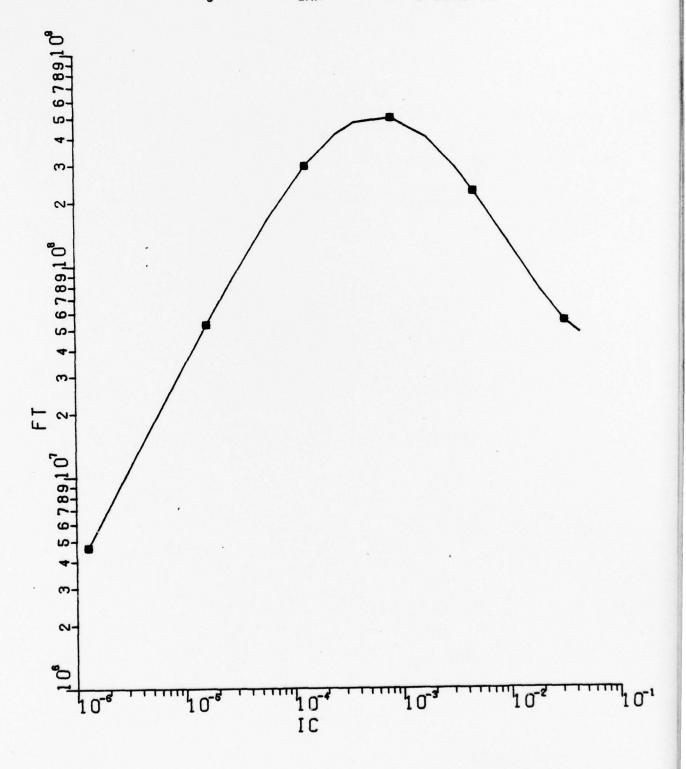
\* -- VTS = 2.500E-02 + -- VTS = 2.700E-02 0 -- VTS = 2.590E-02



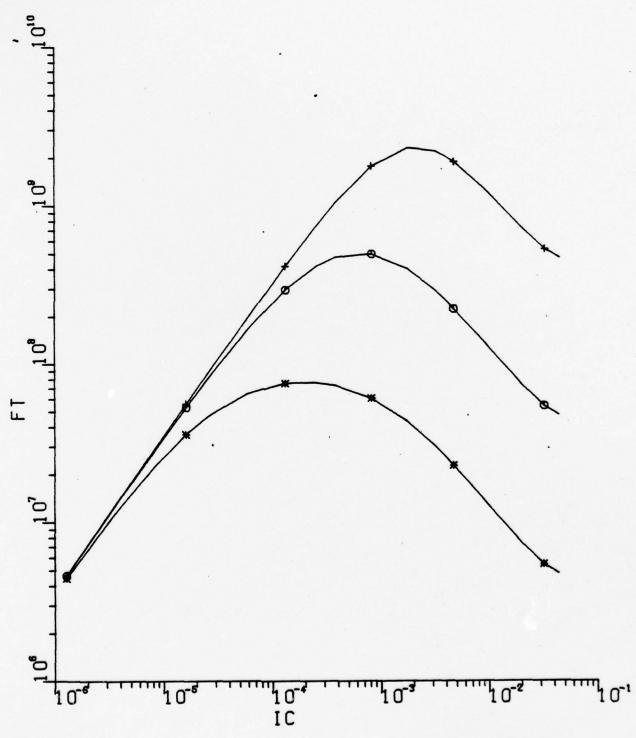
# -- QFN = 1.500E-18 + -- QFN = 1.500E-14 0 -- QFN = 1.500E-16



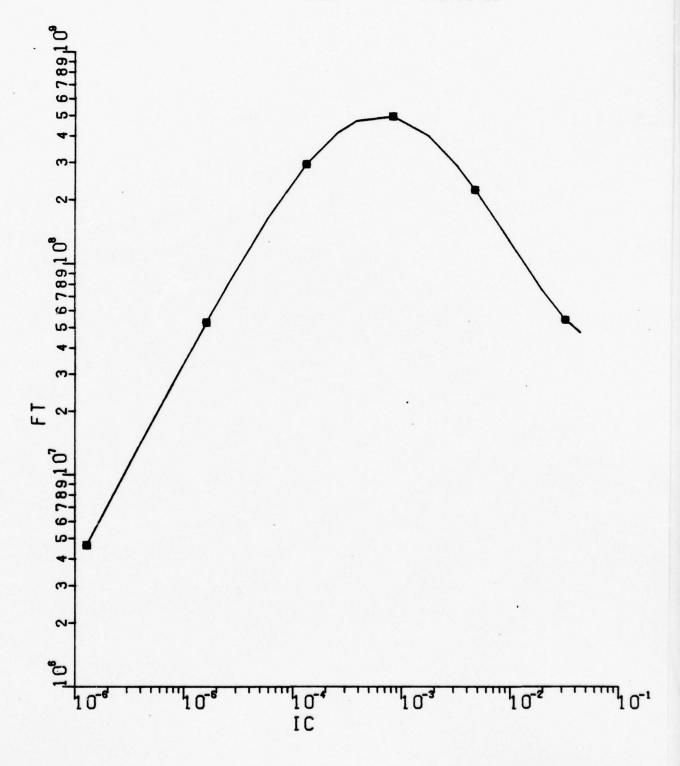
# -- QRN = 1.500E-14 + -- QRN = 1.500E-10 0 -- QRN = 1.500E-12



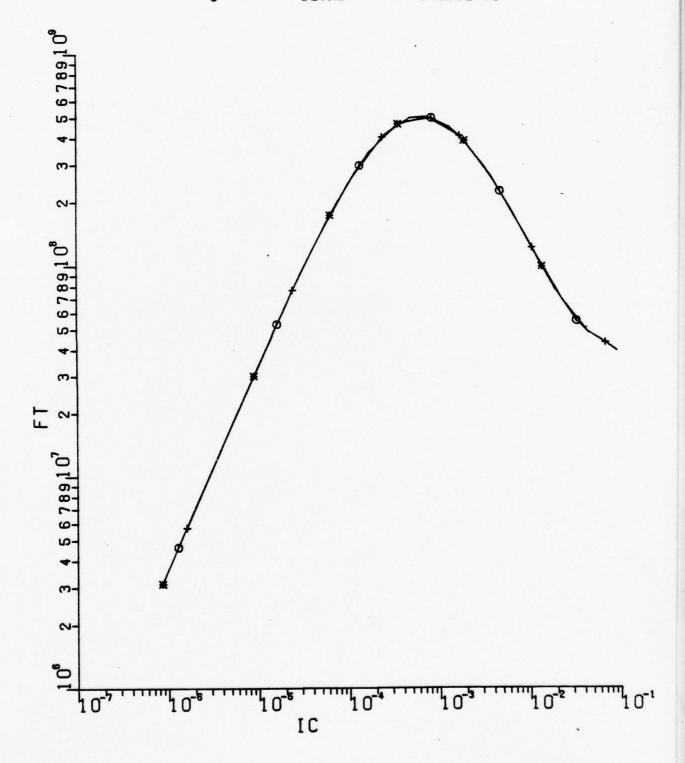




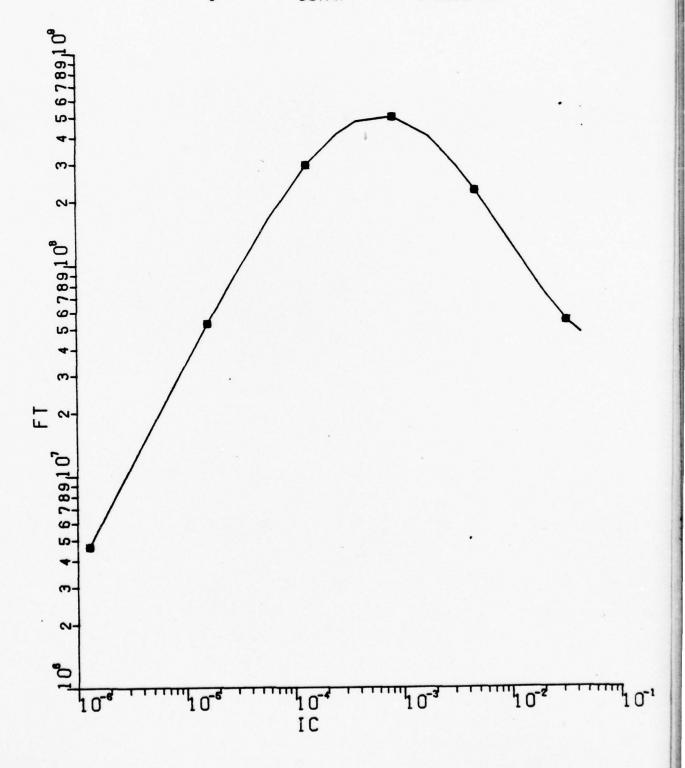
TAUR = 1.000E-06 +-- TAUR = 1.270E-11 0 -- TAUR = 1.270E-09

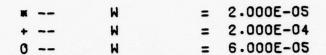


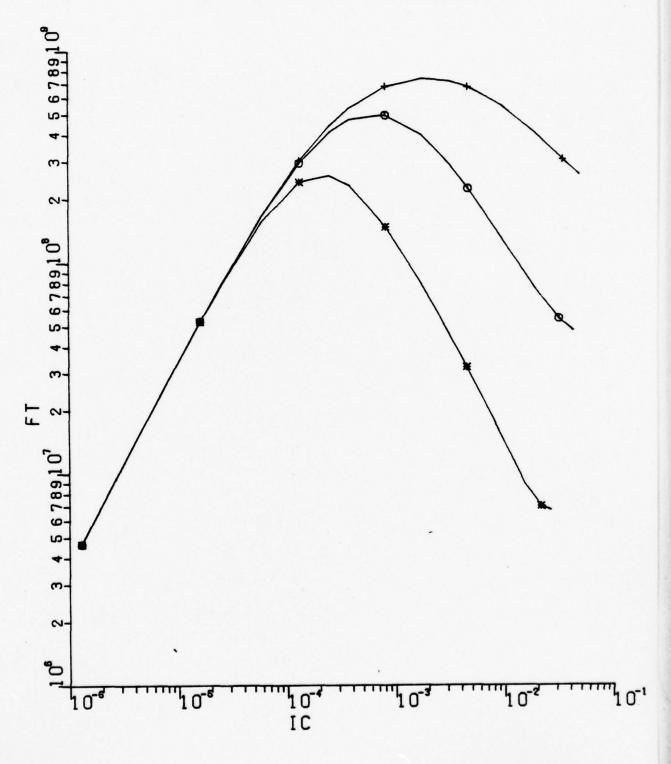
# -- BETAO = 2.000E+01
+ -- BETAO = 1.500E+02
0 -- BETAO = 5.500E+01



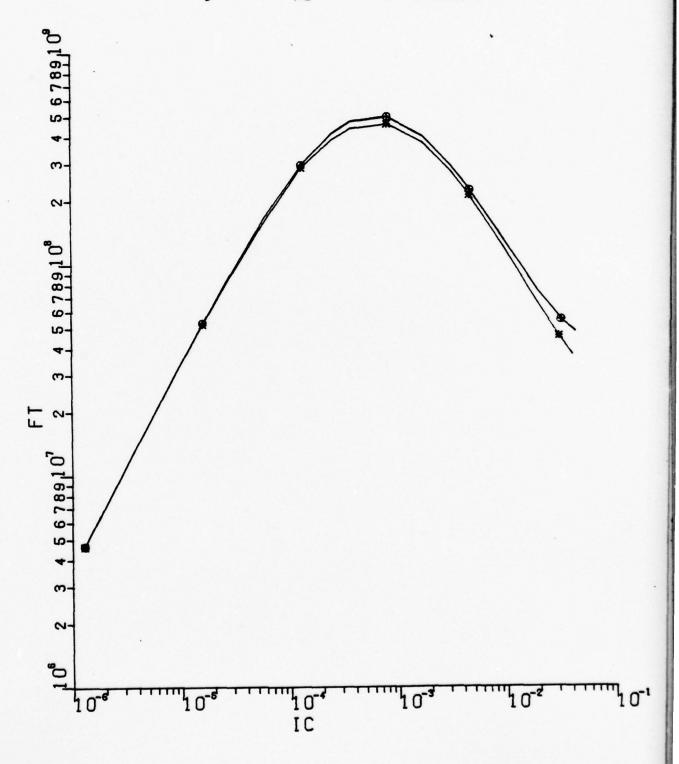
# -- BETAR = 6.000E+00 + -- BETAR = 1.100E+00 0 -- BETAR = 1.500E+00



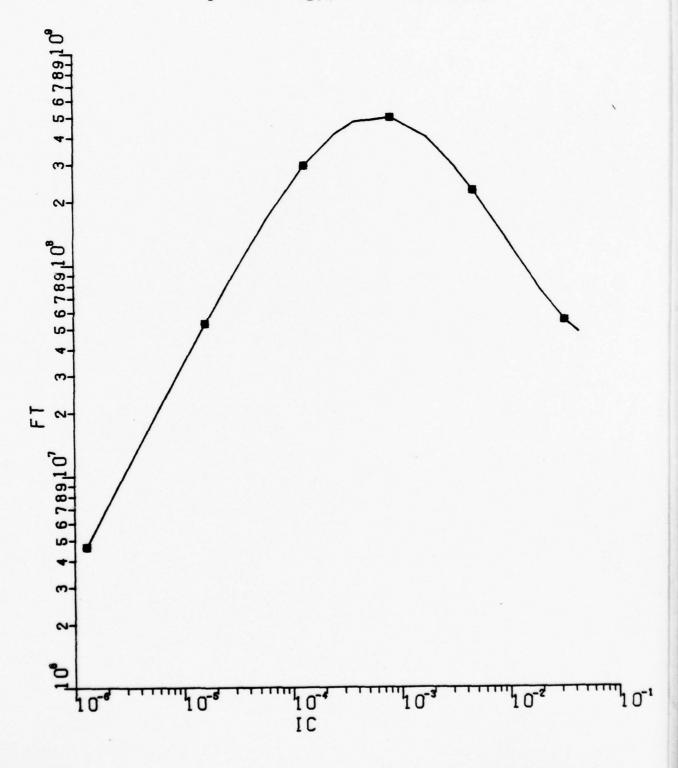




# -- AE = 1.000E-08 + -- AE = 1.000E-04 0 -- AE = 9.800E-07



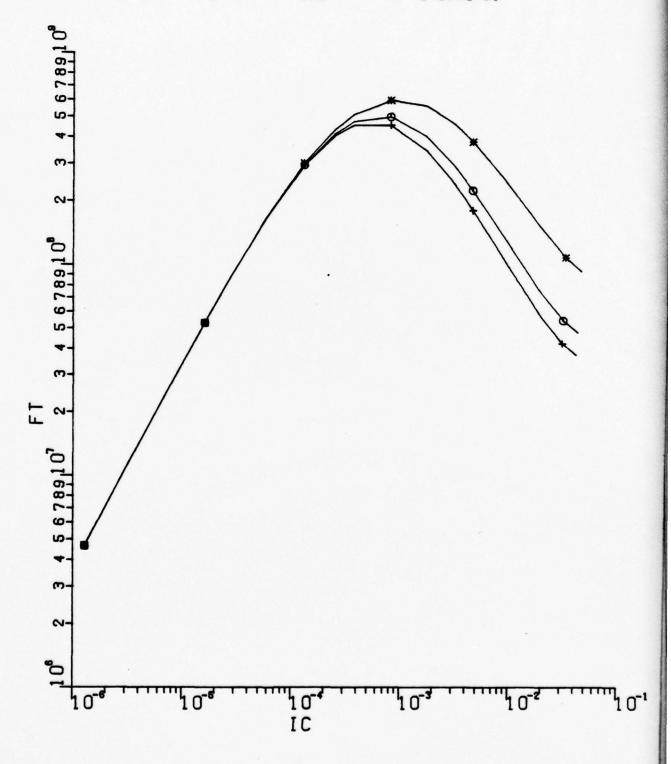
\* -- VLIM = 1.000E+06 + -- VLIM = 1.500E+07 0 -- VLIM = 6.000E+06



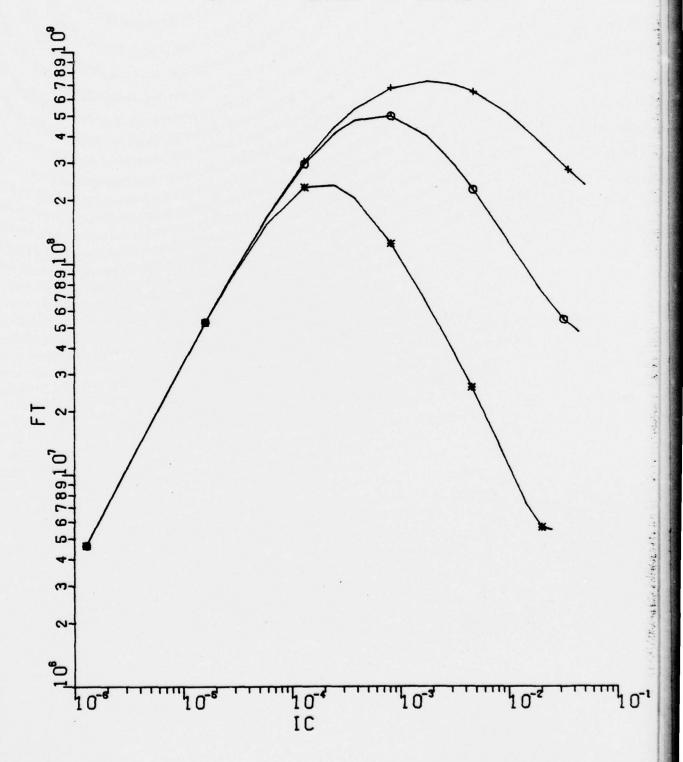
 \* - DNB
 = 1.000E+01

 + - DNB
 = 3.000E+01

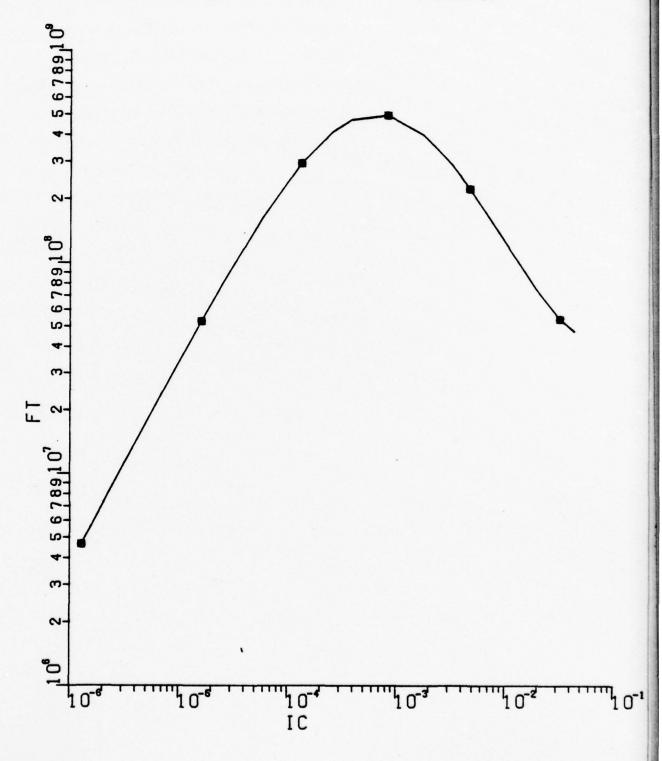
 0 - DNB
 = 2.260E+01



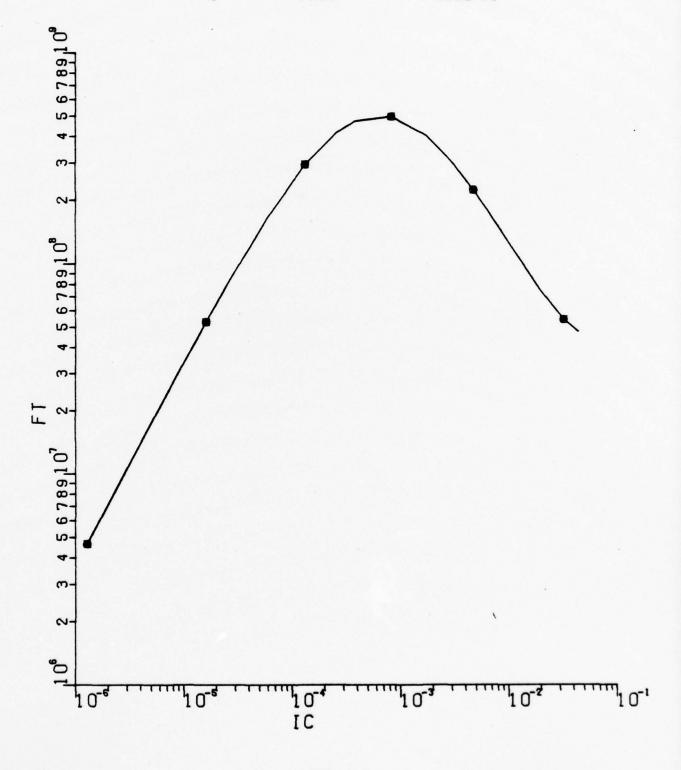
# -- WCPRIME = 1.000E-03 + -- WCPRIME = 1.000E-04 0 -- WCPRIME = 3.000E-04



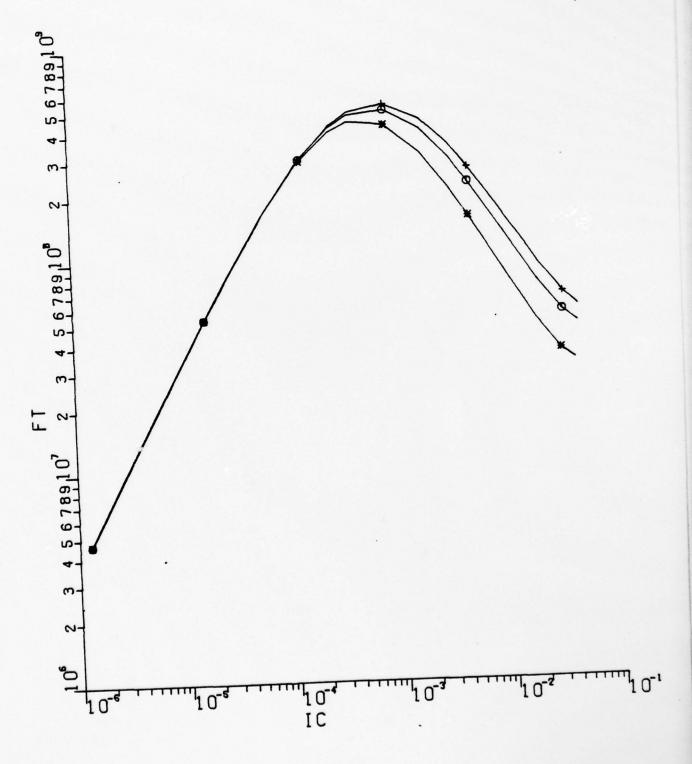




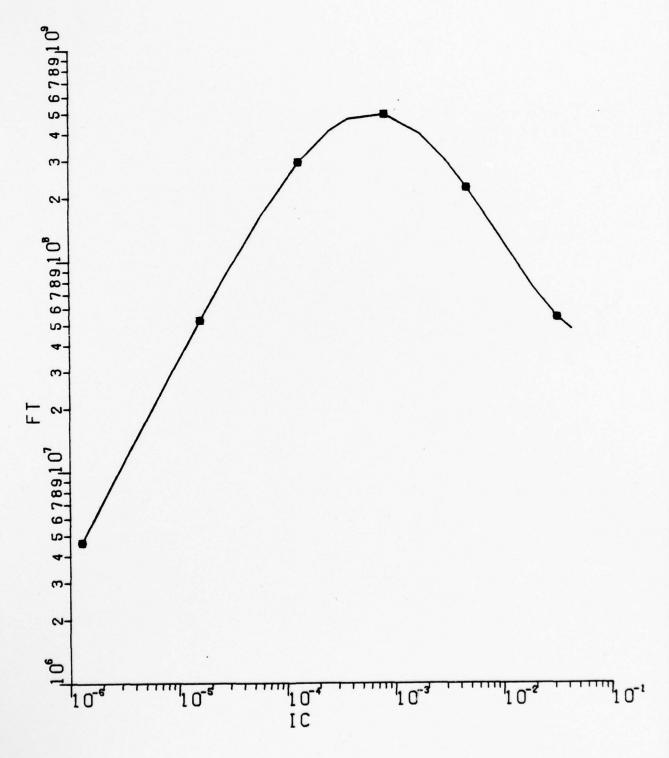
# -- PHIC = 4.000E-01
+ -- PHIC = 1.200E+00
0 -- PHIC = 7.000E-01

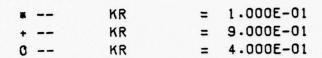


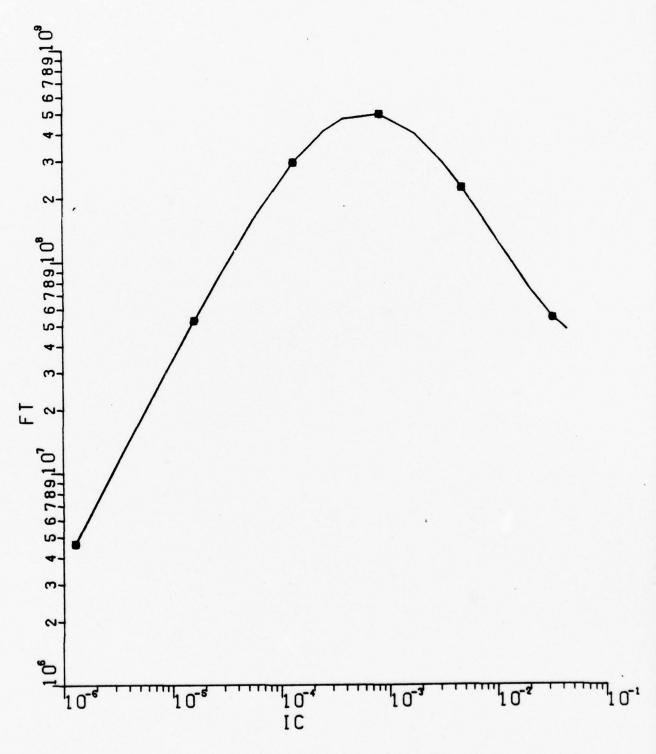




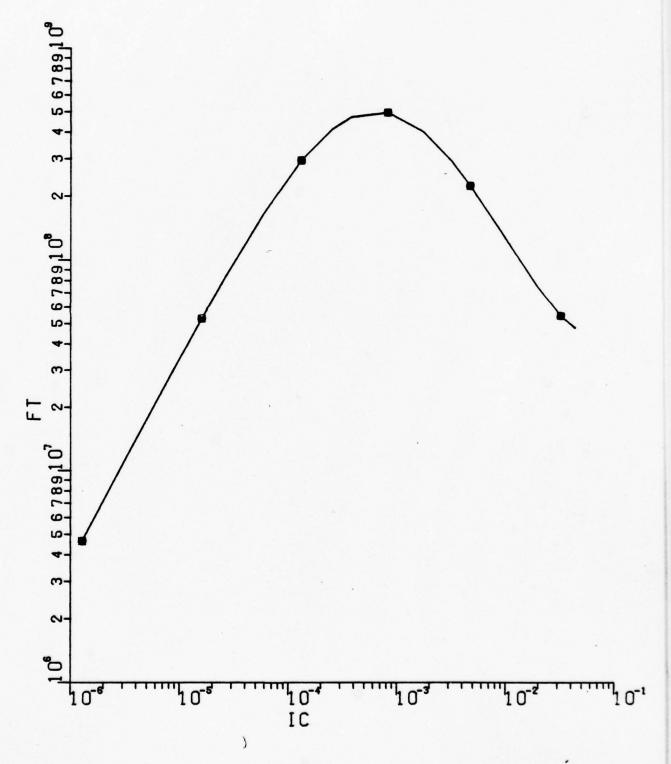




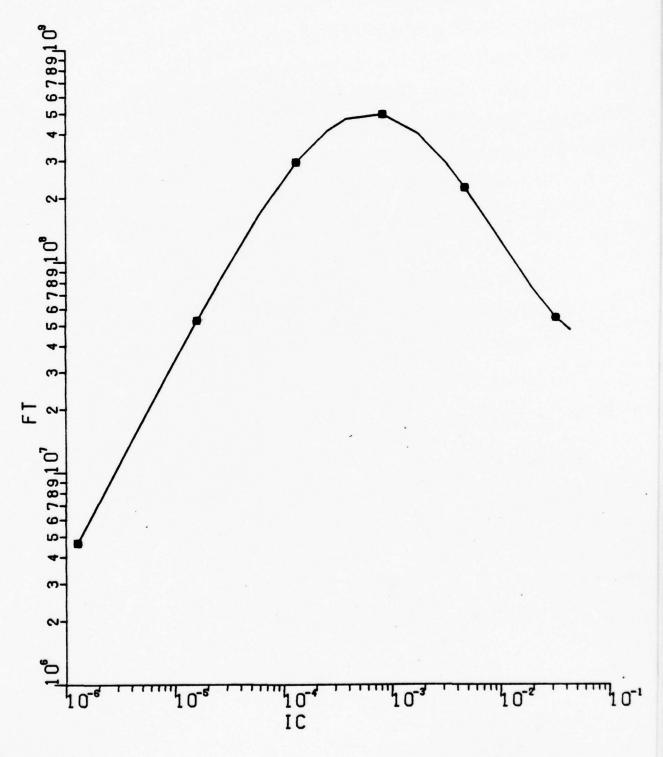




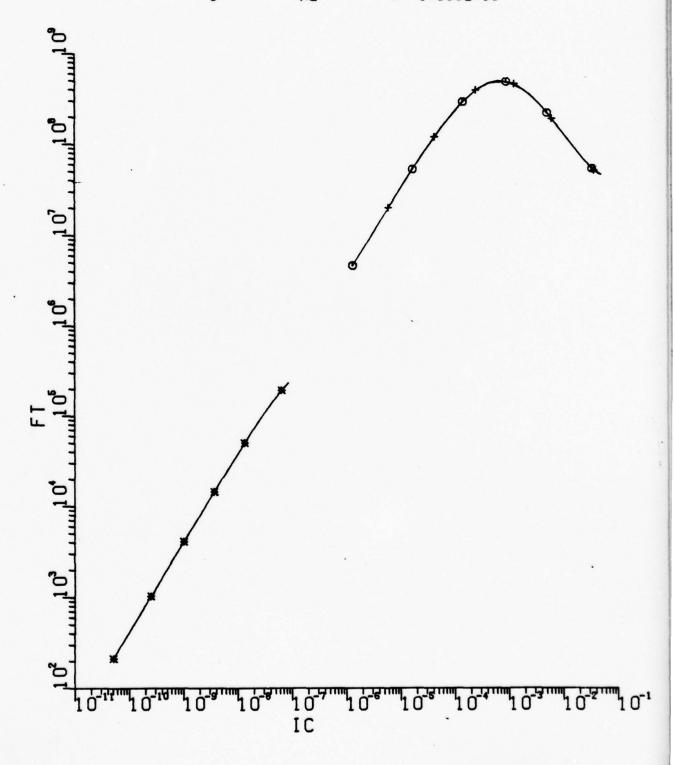
TD = 2.000E-12 + -- TD = 2.000E-10 0 -- TD = 2.000E-11



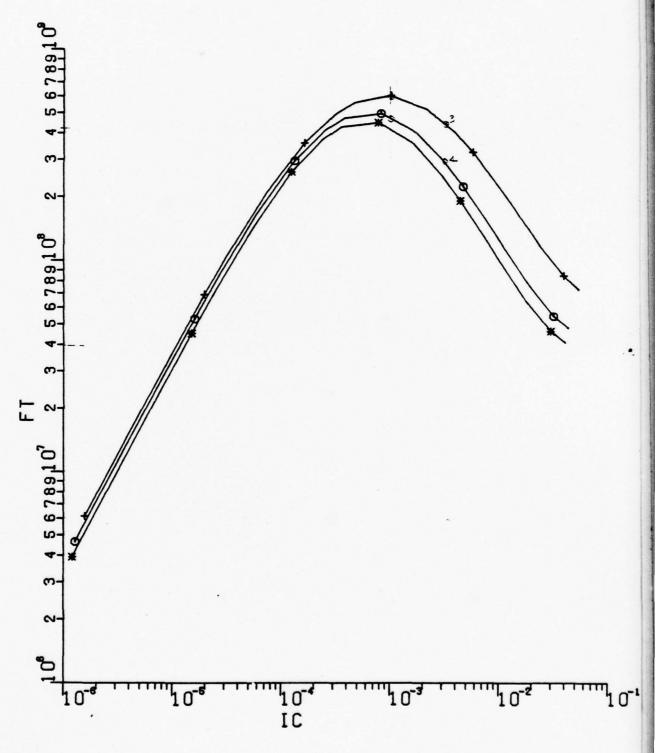




NE = 8.000E-01
+ -- NE = 2.300E+00
0 -- NE = 1.500E+00



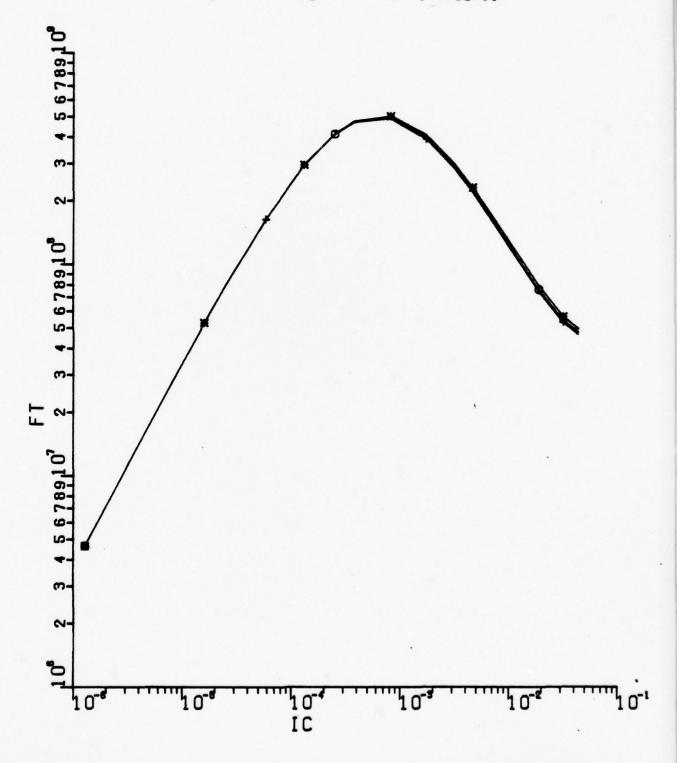




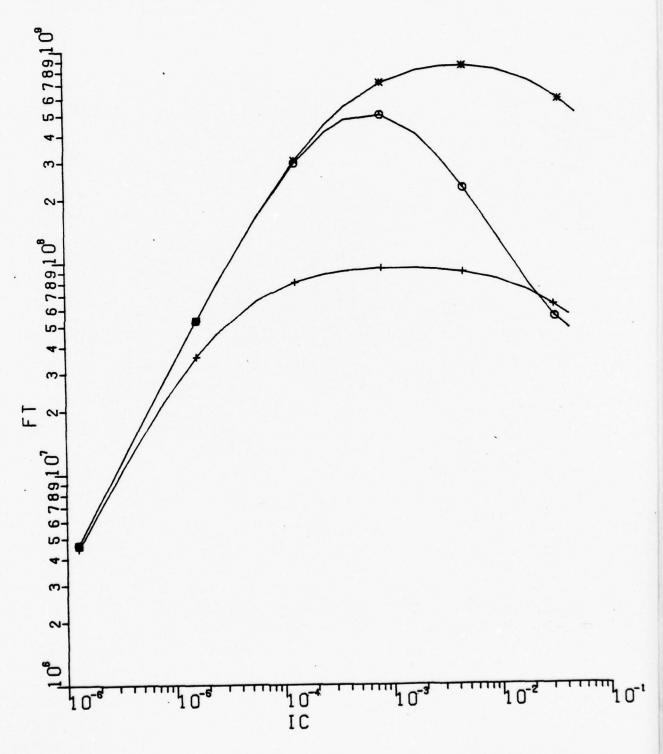
 # - VB
 = 1.500E+01

 + - VB
 = 2.000E+01

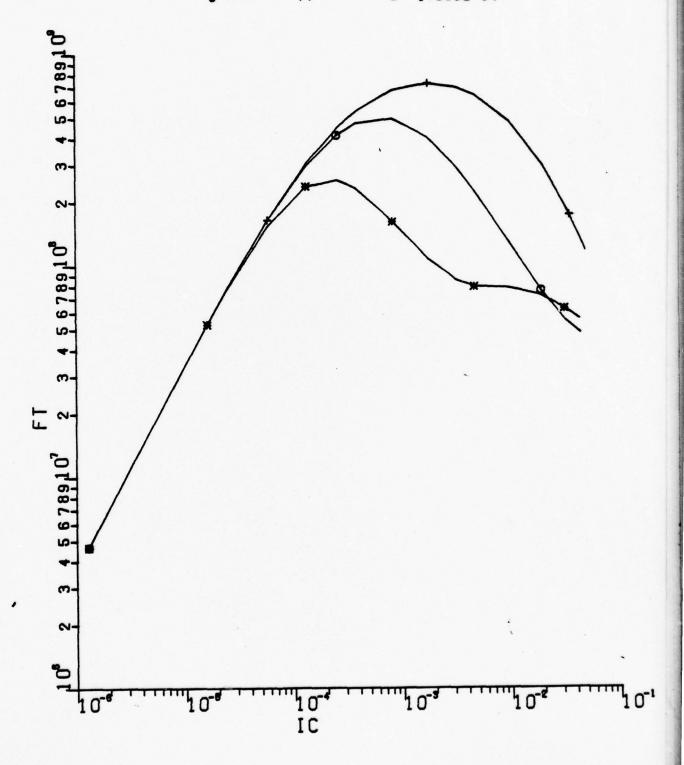
 0 - VB
 = 1.779E+01



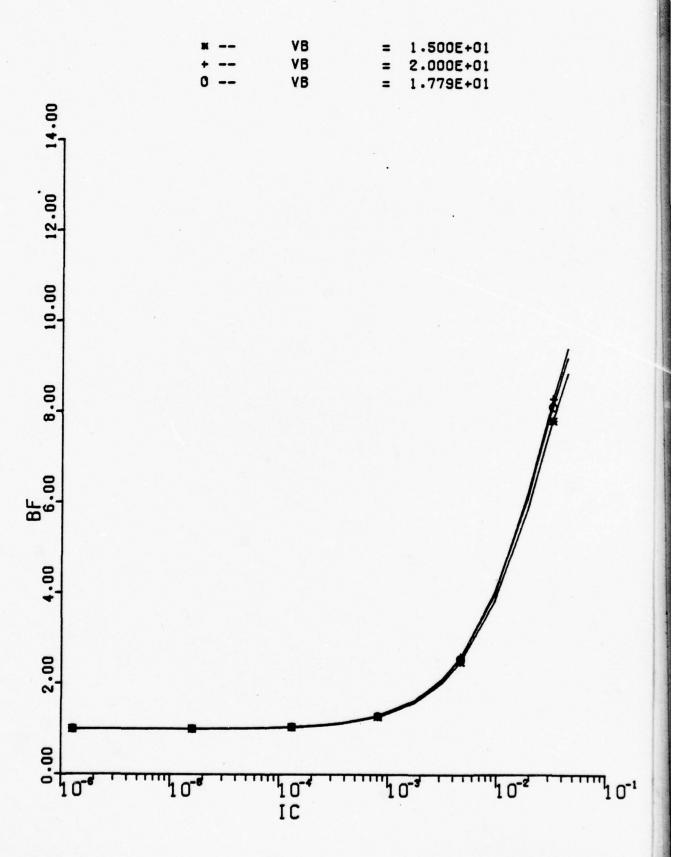




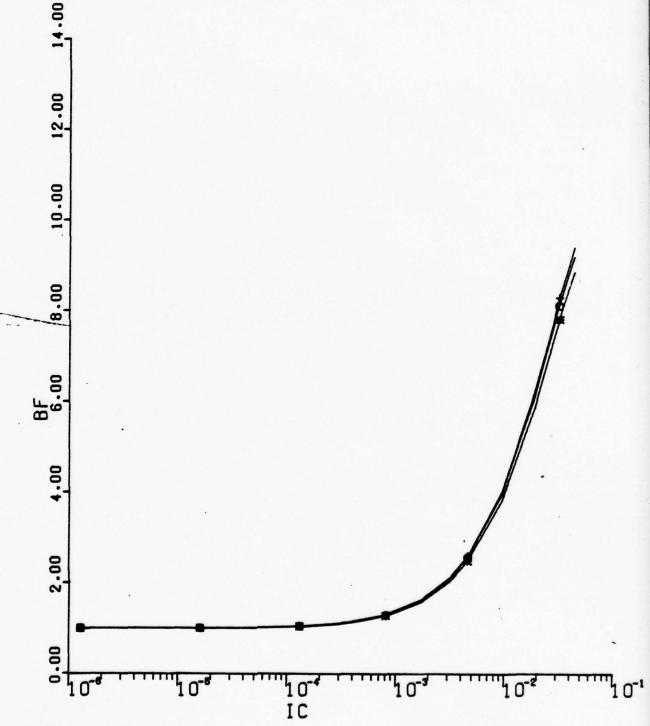
# -- FF = 1.336E-13 + -- FF = 1.336E-15 0 -- FF = 1.336E-14

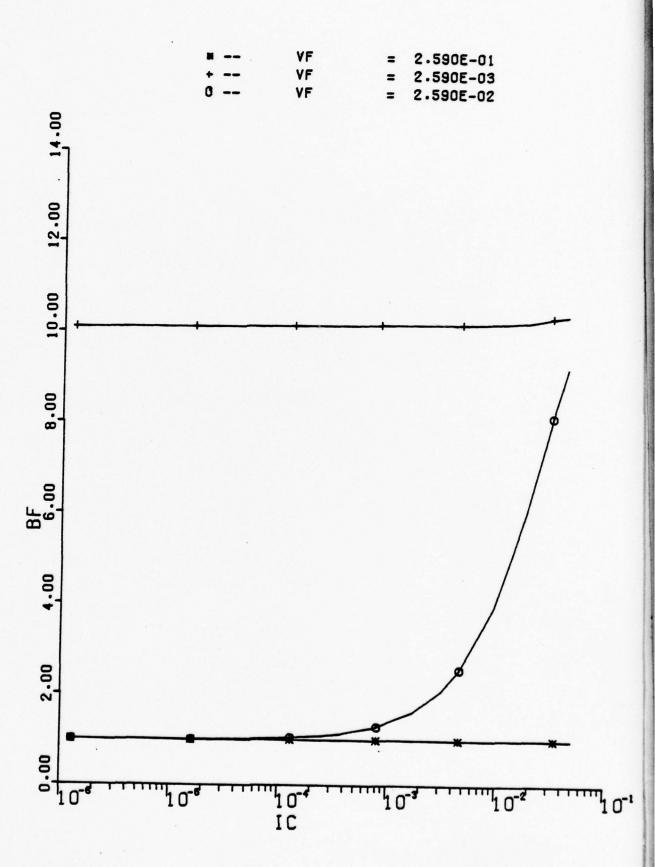


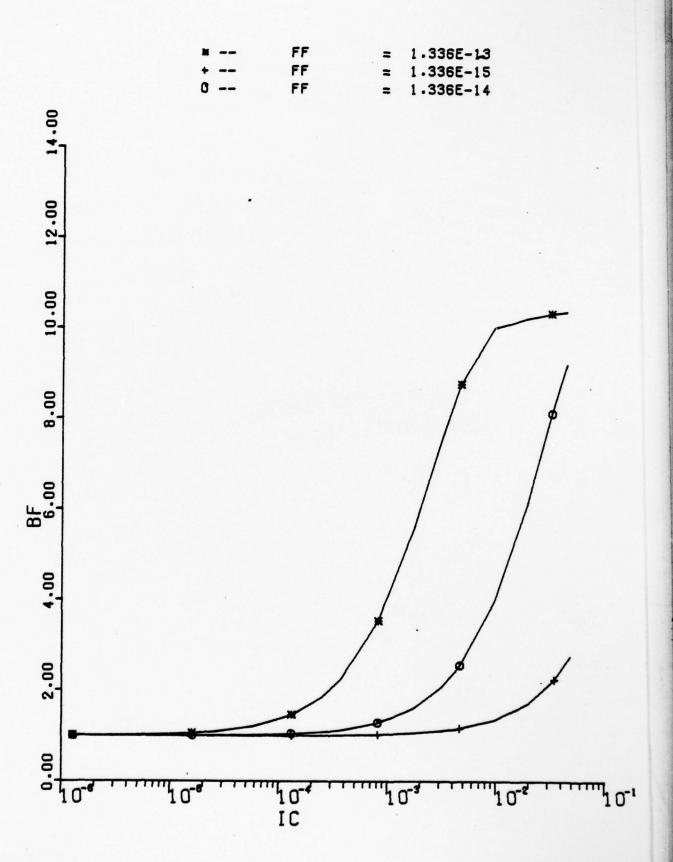
B.5 B<sub>F</sub> vs. I<sub>C</sub> Curves (BASE PUSHOUT PARAMETERS)











## APPENDIX C

PAREV (PARAMETER EVALUATION) CODE LISTING

PROGRAM PAREVILINDUTAUTPUTATAPES INPUTATAPES TAPETATAPEZ) DIMENSION AIB(20) AIC(20) VBE(20) SCR(100), P(3) DIMENSION V(2,20) 2(3,20) PI(3), P0(3) COMMON/PARAVVIZ IERANE IS BETADARBUARBAIBBAREARCCABFZ + VB FF PVER VCR COMMON/OPPT/XIB XICAXVBE COMMON/TEMP/THETA	EXTER OTHEN REAL	CONTINUINININININININININININININININININI	C THIS CODE IS AN INTERACTIVE CODE WHICH ACCEPTS ELECTRICAL AND CECEDIFY DATA AND GENERATES PARAMETERS FOR THE CHOMA BAT MODEL. THE PARAMETERS ARE USED FOR SIMULATION IN GENERAL CAD PROGRAMS COR IN THE COMPANION CODE FOR SIMULATION SPECIFICALLY FOR THE STMULATION OF THE CHOMA MODEL. SEVERAL MODULES IN THE CODE MAY BE BYPASSED VIA DEFAULT OR USER OVERRIDE. HOWEVER, SOME MODULES ARE VITAL TO SUBSEQUENT	### MOD WE EXFCUTION I CONSULT DOCUMENTATION FOR DETAILS.  E. MOCK TRW SYSTEMS/DSSG SEPT. 1977  ***LILLI	AUF HIS C'V C'V E'X E-3	+ 2.7E-49.8E-797.90./ DATA TPISRe IERANGE ISPBETADE IBRANCABETAR / + 27.93E-1293.2E-1491.593.2E-16.55.93E-1391.801.5 / DATA VFR-VCR-60F-60R-9FN-9CN-7AUFD-7AURAR / + 20.0600.90.90.1.5E-16.1.5E-16.1.59E-10.1.27E-9.5 / DATA CSDCEDD.CCD.PHISS.PHICD.PHIED / + 2.6-12.9.6-13.8.8.E-13.9.9.91.2 /
00100 00110 00120 00130 00140 00150		00200	1 1	00310 00320 00340 00350 00350	004100	00440 00440 00440 00540 00510

06300	4 .024001.11.44.17.700240.1.2245_14.8 .1 205_4.
00240	
00260	READ IN THE NAMELIST DATA
U (	
00200	
00900	READ(1, INDATA)
00610 C	*******
00630 C	PERFORM THE DEFAULT CALCULATIONS
9	
٠	
00000	80.F0.0.)
07000	The state of the s
00000	CC. FO. 0. 1
00000	IFCKR. NEC. 10-AE/(AE+AB))
	i
3 01/00	•
200700	ENIER THE TIERALIVE LOUP
200700	
	20
09200	ACCE
00770	GD TO (50, 100, 200, 300, 400, 500, 600, 700, 800, 900, 1100) MASTER
U	*******
	C-CODE FOR JUNCTION CAPACITANCES PARAMETER, FITTING
00810 C	
00820 C	ASSIGN ESTINATE AND BOUNDS
00830 C	2
00840	100 P(1)*2,F-12
00820	PO(1)=1.E=14
00800	P1(1)-1E-10
00810	P(2) 0.9
00880	PO(2) 4
06800	2
00600	P(3)*1.E-12
00600	3
00450	P1(3)-1E-11
06000	200

06600		
00420	TOL = F S= 1 E= 5	
09600	ND=1	
00600	TXX T	
00600	lik•1	
06600	74.0	
9		
Ü		
01010	1(1)1) 0(16 1(5)1) 0(16	
01040	ACCEPT(1) NPT	
01050	1-1.N	
01060		
01070	FIN	
01080	DISPLAY PEMITTER JUNCTION CAPACITANCE FITS	
01000	CALL MINIZITY SCR. P.NV. NPT. NP1. TOL. EPS. EE. NO.	
01100	TALLKAPOAPIA INAFUNEJ)	
01110	DISPLAY CCED EST. PHIED EST. PARASIT. EST.	NSE+
01120	AY PILLEIZIEPIBLEE	
01130	3	
01110		
01150		
	::	
•	C-READ IN THE COLLECTOR CAPAC. CURVE VALUES	
9	1011111	
_		
01200	Ξ'	
01710	Idvet ) no	
01220		
01230	ann i no	
01540	DISPLAT WOLLECTUR JUNCTION CAPACITANCE FITS	7.334
01270	CALL MANAGES DAME AND	• 360
01710		
01280		The second section of the second seco
01200		
01300	1010101110	
01310	60 TO 50	
U	******	
01330 C		
01340 C	READ IN THE SUBSTRATE CAPAC. CURVE VALUES	
01350 C	11-635	

01370 C	
	1
01390	DO 6 1-1
1400	ACCEPT(1) Y(1,1),Y(2,1)
01410	
1420	AY SUBSTRATE JUNCTION CAPACITANCE FITS
1430	
01410	
01450	
1460	
1470	
01480	(2)d=SS   Hd
1490	60 10 50
	********
1	C VER AND VCR EVALUATION CODE
	*****
	400 DISPLA
1560	
1570	~
1580	VCR=XNUM/(G0F+(RIZ+G0R+VBEX))
01200	VER.XNUM/(GOR+(RI1+GOF+VBCX))
1000	
1610	
01620 C	•
1630	C-RELABLE THE INPUT FOR 18-VBE FIT FOR IS/BETAD. TER.NE
	:::::::::::::::::::::::::::::::::::::::
01660	SOO ACCEPT (1) NCTF-VCEX
01670	1=1,
01680	ACCE
01690	10 CONTINUE
01100	
01710	
01730	21.2007 11.20 CONTRACTOR OF THE PROPERTY OF TH
01740	-SET UP
01750	
01760	P(1) = 16 = 14
	12-12)

010799 PO111-PO131-0. 01080 PO121-20	01780	
PO(2)=20 PO(2)=20 NV=2 NV=2 NV=2 NV=2 NV=2 NV=2 NV=2 NV=2	01790	
P1(1)=P1(3)=10  P1(2)=40  NV=2  NP=3  NP=3  NP=4  T01=EPS=1E-5  ND = 1  TV=10	01800	P0(2) = 20
NV-2  NV-2  NV-2  NV-2  NV-2  NV-2  NV-1	01810	P1(1) eP1(3) e1.
NW*2  NW*2  NP3-4  NP1-4  TOL = EPS-1E-5  ND 1  NETT - SOLUTION TO THE ESTINATE FROM IC DATA  NCTM -	01820	P1(2)#40
NP=3	01830	NVe2
TOLEPS-1E-5  ND1  HXIT-80  LLK-1  DISPLAY *POINT ID POINT TEMPERATURE (DEG. C)*  DO 20 I=1,NCTM1  Joint THI-1594.2*Y(2,J1-Y(2,11/4L0G(ALC(J))(ALC(J)))  ZO DISPLAY *IMPUT TEMP. ESTIMATE (DEG. C)*  ACCEPT(5) T  THETA-11594.2/THETA  COUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT  COUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT  COUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT  COUTPUT THE STIMATE  COUTPUT THE SESTIMATE  CONTINUE  C-ACCEPT THE IS ESTIMATE  DISPLAY *POINT  SALIC(1)/FXP(VBE(1)/VT)  DISPLAY *POINT  SALIC(1)/FXP(VBE(1)/VT)  DISPLAY *POINT  CONTINUE  C-ACCEPT THE IS ESTIMATE  DISPLAY *POINTS AND INCREASE NO. OF POINTS IN FIT  CSTART THE 3-PARAMETER MOVING FIT  CSTART THE 3-PARAMETER NOING FIT  CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT	01840	E-d-Z
TOL-EPS-1E-9 NOT-1 NXIT-90 LLK-1 TW-0 LLK-1 TW-0 C	01850	Tolat.
ND=1	01860	Š
CCALCULATE TEMPERATURE ESTIMATE FROM IC DATA  NCTMI-NCTF-1  DISPLAY *POINT TO POINT TEMPERATURE (DEG. C)*  DO 20 I=1.NCTM1  J=-273+11294.2*(1/2.1)-7(2.1)(1+273.)  20 DISPLAY *DIDUT TEMP. ESTIMATE (DEG. C)*  ACC EPT (3) T  THETA-11394.2*(1/2.1)-7(2.1)(1+273.)  C **	01870	NO.1
C	01880	MXIT=90
C	01890	LIKel
CCALCULATE TEMPERATURE ESTINATE FROM IC DATA  NCTM1-NCTF-1  01SPLAY *POINT TO POINT TEMPERATURE (DEG. C)*  00 SO 1=1,NCTM1  J=13594.2*(1/2.4)-7(2.1)/(1+273.)  20 DISPLAY *INPUT TEMP. ESTIMATE (DEG. C)*  ACCEPT(5) THE 13594.2*(1/2.4)-7(2.1)/(1+273.)  THETA-13594.2*(1/2.4)-7(2.1)/(1+273.)  COUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT COUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT COUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT COUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT COUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT COUTPUT THE IS ESTIMATE  CACCEPT THE 1S ESTIMATE  DISPLAY *ENTER THE 1S ESTIMATE*  CACCEPT THE 3-PARAMETER MOVING FIT CSTART THE 3-PARAMETER MOVING FIT CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT	01000	TV0
NCTM1=NCTF-1  NCTM1=NCTF-1  DISPLAY *POINT IO POINT TEMPERATURE (DEG. C)*  DO 20 1=1,NCTM1  J=12  TH=13594.24Y(2,1)/(1+273.)  20 DISPLAY *IMPUT TEMP. ESTIMATE (DEG. C)*  ACCEPT(5) T  THETA-13594.24Y(2,1)/(1+273.)  CONTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT  C VALUE OF TEMPERATURE  C VALUE OF POINTS IN FIT  C VALUE OF POINTS IN FIT IN FIT IN FIT  C VALUE OF POINTS IN FIT IN F	01910 C-	ALC UL AT
DO 20 I=1,NCTN1  J=174  DO 20 I=1,NCTN1  J=174  T=-273+11594,241(2,41-Y(2,11)/ALQG(ALG(4)(4)(1))  20 DISPLAY  THI=13594,24Y(2,1)-Y(2,11)/(1+273.)  20 DISPLAY  THETA-11594,2/17+273.)  VT=1,THETA  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C OUTPUT THE CALCULATED VALUES OF IS BASED ON AN INPUT  C VALUE OF TEMPERATURE  C V	01920	•
DO 20 I=1,NCTH1  J=141  TH=11594,24Y(2,1)/(T+273.)  ZO DISPLAY I,J/1  ACCEPT(3) T  THETA-11594,2/(T+273.)  THETA-11594,2/(T+273.)  THETA-11594,2/(T+273.)  THETA-11594,2/(T+273.)  THETA-11594,2/(T+273.)  C OUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT  C VALUE OF TEMPERATURE  C++  DISPLAY +POINT IS+  DO 25 I=1,NCTF  IS-AIC(1)/FXP(VBE(1)/VT)  DISPLAY +POINT IS+  C-ACCEPT THE IS ESTIMATE  DISPLAY -ENTER THE IS ESTIMATE  DISPLAY -ENTER THE IS ESTIMATE  ACCEPT IS  CSTART THE 3-PARABETER MOVING FIT  CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT  CSTART WITH 4 PUINTS AND INCREASE NO. NE	01930	>
10-14-1 10-17-11594,24(Y(2,4)-Y(2,1))(ALGG(ALG(ALG(ALG(AL))) 20	01940	-
20 DISPLAY 1.994.247(2.1).(14273.) 20 DISPLAY 1.9.7 20 DISPLAY 1.9.7 20 DISPLAY INDUT TEMP. ESTIMATE(DEG. C19 40 CEPT(3) THE ACCEPT(3) THETA-11994.2/(14273) THETA-11994.2/(1427	01950	
20 DISPLAY 1.J.T 20 DISPLAY 1.J.T ACCEPT(5) THE LIPPUT TEMP. ESTIMATE(DEG. C1* ACCEPT(5) THETA-11994.2/(T+273) C+ C OUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT C VALUE OF TEMPERATURE C+ DISPLAY +POINT IS* DO 29 1=1,NCTF IS-AIC(1)/FXP(VBE(1)/VT) DISPLAY 1.IS C-ACCEPT THE 1S ESTIMATE* ACCEPT THE 3-PARAMETER MOVING FIT CSTART THE 3-PARAMETER MOVING FIT CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT	01960	273+
20 DISPLAY 1, J,T ACCEPTIST THETA-11994.2/(T+273) C++ COUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT C VALUE OF TEMPERATURE C++ DISPLAY +POINT IS+ DO 29 1=1, NCTF IS-AIC(1) /FXP(VBE(1)/VT) DISPLAY 1, IS CONTINUE CACCEPT THE IS ESTIMATE DISPLAY VENTER THE IS ESTIMATE ACCEPT THE 3-PARAMETER MOVING FIT CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT CSTART WITH 4 POINTS AND INCREASE NO. NE	01970	-
01990 DISPLAY * INPUT TEMP, ESTIMATE(DEG, C)* 02000 ACCEPT(S) T 02010 THETA-11594,2/(T+273) 02010 VT-1./THETA 02010 VT-1./THETA 02010 VT-1./THETA 02010 VT-1./THETA 02010 CONTOUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT 02010 CONTOUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT 02010 CONTOUT THE CALCULATED VALUE OF 1S BASED ON AN INPUT 02010 CONTOUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT 02100 CONTOUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT 0210 CONTOUT THE SETTIMATE 0210 ACCEPT IS 0210 CONTOUT THE SETTIMATE 0210 ACCEPT IS 0210 ACCEPT IS 0210 CONTOUT THE SETTIMATE 0210 ACCEPT IS 02	01980	_
02000 ACCEPT(9) T 02010 THETA=11594.2/(T+273) 02020 VT=1.7THETA 02030 C++ 02040 C OUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT 02040 C OUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT 02050 C VALUE OF TEMPERATURE 02060 C++ 02070 DT 29 1=1,NCTF 02080 DT 29 1=1,NCTF 02100 DT 29 1=1,NCTF 0210 C-ACCEPT THE IS ESTIMATE 02120 C-ACCEPT THE IS ESTIMATE 02130 ACCEPT IS 02150 CSTART THE 3-PARAMETER MOVING FIT 02150 CSTART THE 3-PARAMETER MOVING FIT 02150 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT 02150 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT 02150 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT	01600	_
02010 THETA=11994.2/(T+273) 02020 VT=1./THETA 02020 C OUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT 02050 C VALUE OF TEMPERATURE 02050 C VALUE OF TEMPERATURE 02050 C VALUE OF TEMPERATURE 02050 D D 29 1=1,NC TF 02100 C CONTINUE 02100 C CONTINUE 02100 C CONTINUE 02100 C C CONTINUE 02100 D D SPLAY FER THE 1S ESTIMATE 02100 ACCEPT IS 02100	05000	2
02020 02030 C++ 02030 C++ 02040 C QUIPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT 02050 C VALUE OF TEMPERATURE 02050 C VALUE OF TEMPERATURE 02050 C++ 02070 D1SPLAY +POINT IS+ 02080 D1SPLAY IP IS 02100 CACCEPT THE IS ESTIMATE+ 02100 D1SPLAY +ENTER THE IS ESTIMATE+ 02100 ACCEPT IS 02150 CSTART THE 3-PARAMETER MOVING FIT 02150 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT 02150 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT 02150 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT 02150 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT 02150 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT 02150 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT	02010	
02030 C++ 02040 C QUIPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT 02050 C VALUE OF TEMPERATURE 02050 C++ 02070 DTSPLAY +PDINT IS+ 02080 DD 25 1=1,NCTF 02090 IS-AIC(1)/FXP(VBE(1)/VT) 02100 DISPLAY ID IS 02110 Z5 CONTINUE 02120 CACCEPT THE IS ESTIMATE 02130 ACCEPT IS 02140 ACCEPT IS 02150 CSTART THE 3-PARAMETER MOVING FIT 02150 CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT 02170 CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT 02150 ACCEPT S IER NO. OF POINTS IN FIT 02150 ACCEPT S IER NO. OF POINTS IN FIT 02150 ACCEPT S IER NO. OF POINTS IN FIT 02150 ACCEPT S IER NO. OF POINTS IN FIT	- 1	>
02040 C GUTPUT THE CALCULATED VALUES OF 1S BASED ON AN INPUT 02050 C VALUE OF TEMPERATURE 02060 DTSPLAY *POINT IS* 02090 IS*AIC(1)/EXP(VBE(1)/VT) 02100 DISPLAY IP.IS 02110 25 CONTINUE 02120 C-ACCEPT THE IS ESTIMATE* 02130 DISPLAY *ENTER THE IS ESTIMATE* 02140 ACCEPT IS 02150 C************************************		
02050 C VALUE OF TEMPERATURE 02060 C++ 02060 D		TPUT THE
02060 C++ 02070 DISPLAY +PDINT IS+ 02080 DD 29 I=1,NCTF 02080 DS PI=1,NCTF 02080 DISPLAY IP IS 02100 DISPLAY IP IS 02110 25 CONTINUE 02120 CACCEPT THE IS ESTIMATE+ 02130 ACCEPT IS 02140 ACCEPT IS 02150 CSTART THE 3-PARAMETER MOVING FIT 02160 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT 02170 CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT 02190 ACCEPT SAME TERMOVING FIT 02110 ACCEPT SAME TERMOVING FIT		UE OF TEM
02070 DISPLAY *PDINT IS* 02080 DD 25 1=1,NCTF 02090 DISPLAY IS* 02100 DISPLAY IS* 02110 25 CONTINUE 02130 DISPLAY *ENTER THE IS ESTIMATE* 02150 C************************************		
02080 DD 29 1=1,NCTF 02090 IS=AIC(I)/FXP(VBE(I)/VI) 02100 DISPLAY ID IS 02120 C—ACCEPT THE IS ESTIMATE 02130 DISPLAY **ENTER THE IS ESTIMATE** 02140 ACCEPT IS 02150 C==-START THE 3-PARAMETER MOVING FIT 02170 C==-START WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT 02190 ***MARE*** NE	02020	IN *POINT
02090 IS=AIC(I)/EXP(VBE(I)/VI) 02100 DISPLAY ID IS 02120 CACCEPT THE IS ESTIMATE 02130 DISPLAY FENTER THE IS ESTIMATE 02150 CSTART THE 3-PARAMETER MOVING FIT 02150 CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT 02170 CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT 02190	02080	_
CSTART THE 3-PARAMETER MOVING FIT CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT CSTART WEER MOVING FIT CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT CSTART WEER MOVING FIT CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT	05030	IS-AIC(1)/FXP(VBE(1)/VI)
CSCEPT THE IS ESTIMATE  CACCEPT THE IS ESTIMATE  ACCEPT IS  CSTART THE 3-PARAMETER MOVING FIT  CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT  CSTART WERE  A MEE	02100	DISPLAY I. IS
CACCEPT THE IS ESTIMATE  DISPLAY **ENTER THE IS ESTIMATE **  ACCEPT IS  CSTART THE 3-PARAMETER MOVING FIT  CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT  OISPLAY **NPTS** IER**  * * * ****  ****  ****  ****  ***  ***  ***  ***  ***  **  ***  ***  ***  ***  ***  ***  ***  ***  ***  ***  **		5 CONTINUI
DISPLAY FENTER THE IS ESTIMATE **  ACCEPT IS  COCCOCCOCCOCCOCCOCCOCCOCCOCCOCCOCCOCCO	Ü	-ACCEPT THE
Cossosososososososososososososososososo	02130	PLAY *ENTER THE IS
CSTART THE 3-PARAMETER MOVING FIT CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT OISPLAY WHOTS IER  • • MEE*	05170	T IS
CSTART THE 3-PARAMETER MOVING FIT CSTART WITH 4 POINTS AND INCREASE NO. OF POINTS IN FIT DISPLAY 4NPTS IER A B MSE6	U	***
CSTART WITH 4 PUINTS AND INCREASE NO. OF POINTS IN FIT OISPLAY 4NPTS IER NE NE NE NEE NEE NEE NEE NEE NEE NEE		THE
DISPLAY +NPTS IER NE	_	TART WITH
+	02180	DISPLAY DADIS TER
	02190	- MAR-

02200	06 00
02210	
02220	+ 669
02230	
02250	30 DISPLA
05260	
02270	ACCEPT 60 T9
05290	
02300	C BUALLIATE THE DELEBOR DIONE FLANDACTED COTTON AND TO TO DELABORATED
02320	
02330	
02340	900 ACC
02350	17 00
02360	ACCEPT
02370	71 CONTINUE
02380	CSTART THE 3
02390	CSTART WITH
02400	DISPLAY
02410	+ • MSE•
02420	00 01
02430	כעור
05420	
06420	
02400	NC-THE
02480	-
02490	DISPLA
02500	ACCEPT
02510	1 09
02520	C++++
02530	J
02540	C-EVALUATE T
02550	
02560	200
02570	ACCEPT
02580	PNST
05220	IBBORNST
00970	DISPLAT PINPUT
02910	2

02020	
2000	
03040	
03050	2.
03060	TOLeEPsife-5
03070	1
03080	TX11-00
03090	
03100	0
03110	DISPLAY 40FN FIT+
03120	DISPLAY *NPIS** GEN ESI** MEAN SQ.ERROR** MSE/NPIS*
03130	IC#2, NCTF
03140	CALL MINIZ(2,SCR,P,NV,NNC,NP1,TOL,EPS,
03150	
03160	
03170	DISPLAY NNC. P(1). EE.RHEP
03180	-
03190	DISPLAY + INPUT GFN ESTIMATE+
03200	0
03210	20
03220	****
03240	C CALCULATE TAUFD FROM MAXIMUM ET DATA POINT
1	
03260	
03270	900 CJE-CEO/(1,-,7/PHJED)++.5
03280	•
03290	÷
03300	-
03310	SPLAY
03320	60 17 50
03330	•••••••
03340	
03350	C TAUR CALCULATION FROM SATURATION TIME DATA
03360	
03370	
03380	1000 DISPLAY GINPUT 18F, 18F, 1CF, 15+
03390	ACCEPT RB1, RB2,
03400	75
03410	TAUR=TSAT+(1.+1./BETAR-BETAD/(BETAO+1.))-
03450	
03430	
03440	
•	

03450	60.10 50
03460	***********
03470	
03480	C DUIPUT THE VALUES OF MODEL PARAMETERS TO UNIT (2)
03480	
03500	C*************************************
03510	1100 CONTINUE
03520	VF=VT
03530	BS=BC=1,/3.
03240	BF • 0 • 5
03220	PC-RHDC
03260	WRITE(25TRANS)
03570	ENO
03580	SUBROUTINE ZFIND(Z, FZ)
03200	CFIND CALCULATES Z GIVEN 18,188 IN CONJ. WITH GRIM ROOTFINDER
03600	COMMON/PARAM/SKIP(7), 108/SKIP2(7)
03610	COMMON/OPPT/18,5KIP3(2)
03620	REAL 180¢18
08960	•
03640	RETURN
03650	END
03960	SUBRO
03670	CSUBR
03680	
03690	
03700	C REMAINING DATA IS FROM THE PREVIOUSLY INPUT
03710	C IC, IB, VBE TABLES AND MODEL PARAMETERS
03720	COMMON/PARAM/VI, IER, NE, IS, BETAD, RBB, IBB, RE, RCC, BF2,
03730	+ VB. FF. VER. VCR
03740	
03750	EXTERNAL ZFIND
03760	REAL
03770	╸.
03,480	TALL CALL
09/40	C-FIND OB BY INTERPOLATION
0000	200
03000	
03830	TOTAL SELECTION TO SOLUTION TO
03840	
03850	VEDT=XVBE-X1B+(RBD+RB)+RB+1ER+EXP((XVBE-X1B+RBD)/(NE+VT))
03860	VEOT-VEOT-RE+(XIB+RICIN)

03990 C-LINEAR 03920 C-LINEAR 03920 C-LINEAR 03920 C-LINEAR 03950 10 11	PC_AMAXIOO.ACC+61SQC_CXF-1.)/BF2)) VCOT-VEDT+(XIB+RICIN)*RE-***********************************	
	C-EXACT VALUE IN TABLES  20 DUTP-Y(I)  RETURN  C-INTERPOLATION REQUIRED  30 DUTP-Y(I-1)+(VAL-X(I-1))+(Y(I)-Y(I-1))/(X(I)-X(I-1))  RETURN  RETURN  100 FORMAT(*INTERP ROUTINE IS OFF THE LOW END FOR INPUT-*,E10.4,  * *AND MIN TABLE VALUE-*,E10.4)  110 FORMAT(*INTERP ROUTINE IS OFF THE HIGH END FOR INPUT-*,E10.4,  * *AND MAX TABLE VALUE-**E10.4)  END	
04120 04130 04140 04150 04160	ST S	68 5 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1
04180 04190 C 04200 C 04220 04230	10UT=IN BEGIN <u>Loop to find n roots</u> Do 170 L=L1¢L2 T(1)=G(L)	
04240 04250 04260 04270 04280 100	H=-1.0 IF(T(1).NE.O.) H=-T(11/10. D=5 J1=MAXIT MAXJ=J1	GRTH 13 GRTH 14 GRTH 15 GRTH 16

IFCA		GRIM	19
T(2)		GRIM	50
T(3)-T(	-T(1)+H	GRIM	21
3		GRIM	22
BEGIN	N ITERATIONS TO FIND ONE ROOT	GRIM	23
		GRIM	24
00	150 J-15 MAXJ	GRIM	25
CALL	AUX(RT, FRT)	GRIM	26
FPRT-FR1	-FBT	GRIM	27
IFIL	.Fo.1) 60 TO 110	GRIM	28
J		GRIM	59
EVALUAT	UATE FP(RT)	GRIM	30
		GRIM	3.1
00 10	162.1	GRIM	32
DEN-RT		GRIM	33
TF( J.6T	67.3) 60 TO 105	GRIM	36
71.7		GPTM	
TELAB	BELDEN, GE 1. E-201 GD TD 10K	1100	
		E 1 4 9	000
		E 25	31
RESTARI	ITERATION WHEN ANY OF THE THREE	GRIM	36
	VALUES ARE TOO CLOSE TO A PREVIOUS ROOT	GRIM	39
		GRTH	0,
TF(10UT	DUT.EQ.O) WRITE(6,500)L	GRIM	41
T(1)•T(	-T(1)+,001	GRIM	42
	Tn 100	GRIM	43
105 FPRT-FPE	•FPRT/DEN	GRIM	**
	DUT. E0.0) WRITE(6,500)L.RT.FRT.FRT	GRIM	45
		GRIM	40
C RT 15	S A ROOT IF F(RT) AND FP(RT) ARE BOTH SHALL	GRIM	4.7
		GRIM	48
IF	(ABS(FRT).LT.1.E-20).AND.(ABS(FPRT).LT.1.E-20)) 60 TO 160	GRIM	64
IFCJ		GRIM	20
CCX	PT	GRIM	51
TECJ	J.E0.3) 60 TO 114	GRIM	52
RTeT	=	GRIM	53
60 10	0 150	GRIM	34
114 RTeT(	=	GRIM	55
09	-	GRIM	26
115 IFCAB	S	GRIM	57
		GPTM	. e

		T as	19
		2100	1 2
		E-80	70
-		GRIA	60
	061 11 130	GRIM	•
		GRTH	69
	MOVE DATA	GRTM	99
		GRIM	29
20	×	GRIM	99
	X(2)•X(1)	GRIM	69
	X110sfPRT	GRIM	20
	10-0	GRIM	11
		GRIM	72
	CALCULATE NEXT APPROXIMATION	GRIM	73
		GRIM	74
25	A=1.0+D	GRTM	75
	B-X(3)+0+0-X(2)+A+A+X(1)+(A+D)	GRTM	92
	DEN-88-8-(4.0+x(1)+0+A)+(x(3)+0-x(2)+A+x(1))	GRIM	11
	DI2.0+x(1)+A	GRIN	78
1	1	GATA	42
	TIOEN	GRIN	90
		GATA	81
	35	GRIM	62
130	IFCTOUT.EG.O) WRITE(6,500)L,DEN	GRIM	83
	.0.1 60	GRIM	94
	DFN-B	GRIM	85
35	DI-01/0FN	GRIM	96
9	H-01+H	GRIM	87
	AT-RT+H	GRTM	98
	TF (ABS(H).LE.EPS1*ABS(RT)) GO TO 155	GRTM+1	-
2	CONTINUE	GRIM	90
25	CALL AUX(RT, FRT)	GRTM	16
9	C(L)*RT	GRIM	92
	IF(TOUT.EQ.O) WRITE(6,501)L,RT,FRT	GRIM	63
170	CONTINUE	GRIM	*6
	IF(IOUT.EQ.0) WRITE(6,502)	GRIM	95
	RETURN	GRIM	96
8	FORMAT(13, 3£20,8)	GRIM	44
201	HAT	GRIM	86
20	FORMAT(1H1)	GRIM	66

PNDPHXTT-LLK,POPPLIMFUN) SUBRDUTINE HINIZ FOR BHDO7R APRIL 4, 1966 FENSION A(NPL,NPL),X(NV,NC),P(100),FP(100),D(100),PO(100),P1(100) (10)				The state of the s		15xe1P9612.4/(22xc1P9612.4))	E12.4/(22x,1P9E12.4))	ROR PARAMETERS*/14x, *MEAN*/14x, *SOUARE			CHARLES OF COLUMN COLUM						P. LLK1		
101	05210 C	ANN	05260 DO 557 I 10 NP	o	05290 EPS1=5.+EPS+1. 05300 WRITE(6,23)(PO(I),1=1,NP)	WRITE (6.22) (PICE	22 FORMATI COMAXIMA	195	194 (11.5	5	÷	6 00	926	29 00	On 45	82	כ כעון	CALL FUNCO, FF	05530 IF(IN.NF.01/C*XIIV.)

						CONVERGING)  DEVIATIONS OF THE PARAMETERS*//	
	6 JF(EE, LE.FO) GO TO 871  NAIL-NAIL+1  JENNAIL+1  OF A72 T-1 NO B TO 871	872 D(1)=PPO(1)+D(1)/2. 672 D(1)=D(1)/2. 60 T0 29	29 61	TECE TECE TECE TECE TECE TECE		00 WRITE (6,201) 01 FORMAT(30HOTHE PROCESS IS NOT 60 TO 470 00 [LL-LLL-] 1 FFLLL-NE.0) 60 TO 121 70 DO 778 I=1.0 NP 778 FP(1) = SQRT(ABS(A(1,1)*EE)) WRITE (6,779) (FP(1)*I=1.0 NP) 779 FN(1)*SQRT(ABS(A(1,1)*EE))	
05550 05560 05570 05580	05620		00 T 6	05740 05750 05760 05760 05780	7 ~	20 - 4	09650

									•			
	TERS+)			TO 843			VARIABLES		The state of the s		1(100), P(100)	
	TI+OASYMPTOTIC CORRELATION MATRIX OF THE PARAMETERS+)			11).Eq.0.0.0B.FP(J).Eq.0.0.0B.A(I.J).Eq.0.0)GD TD 843	TARREST CONTRACTOR OF THE PROPERTY OF THE PROP		Y-F STANDARD		to the state of th	J. T. L. MVJ.	DUTINE STEP(A,D,P,PO,PI,NPI,TDL) SUBROUTINE STEP FOR BHDOTR NSION A(NPI,NPI),V(100),D(100),IN(100),P0(100),P(100))	
-SGRT(ABS(A(1,1))) 1 J-1,1 1-4(1,1)	NPTOTIC CORRELATION	NO(NP.11+10) (6,972)(1,1-10,11)		0.0.0R. FP(J).E0.0.0.		[6,374)[,(8(J),J=1,K) [7(14,2%,10F12.5) [.NE.NO)GO TO 37]	F(6,353) AT(*1CASE F Y-F (************************************	1 Lelanc acao FUNCOFPEX(12L)	JelonP #SDEST+FP(1)+FP(J)+A(1,J) #SORT(SDEST+EE)	ND.L.)-0 (6,592)L.0 Q.TT.SDESTP (X(10L)) 1912.NV) .T(1X)14,10F12.5/(41X)7F12.5)	TEP (As D.P. PO.P I.S NPI.S NE STEP FOR BHDOTR NPI.S NPII.S V(100), D(10	
PP(1)=SORT(AL DO 841 J=1.1 B41 A(J.1)=CA(I.J.1) WRITE(6.774)	LI OR HA	22.0	0 ×	K-K+1 IF(FP(I).EQ.O.	2 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0		515 WRITE(6,553) 553 FORMAT(+1CAS) 1+/30x,+DEVIA	SOEST-0.0 CALL FUNCO, F	**		SURROUTINE S SUBROUTI	NP=NP1-1
	06010 77			06100		06160	-	06220 06230 06240 06250			06340 06350 C 06350 C	06370 C

	VID-SORTIALID
00049	
06410	50
6420	2
6430	A(1,2)-A(1,2)/(V(1)+V(1))
9440	2
	ACION
6460 20	CONTIN
	00 2 1
A 0044	
6510	-
6520	IF(IN(I).NE.3 .DR. A(I.1).LE.TOL) 60 TO 1
6530	
9940	
5550	
9290	0.0
1 0259	CONTINUE
5580	IF(K, E0.0) 60 T0 3
2590	
\$ 0099	DN 5 101pK
2610	D(1)-A(K,1)
\$ 0299	=
2630	PPOCKI
2640	00 6 I-K,NP1
9650	=
9 09990	ACTAN 1-0
2670	ņ
9899	IN(K)-IN(K)+1
06990	00 7 T-1/NP1
0029	-
6710	2
6720	60 TO 25
92 06730	Y=0.
	00 7 3
2 0929	A(1,1)-A(1,1)-0(1)+Y
6770	2
6780 3	# · 1 ·
9449	-
6800	IF(IN(I).NE.1) 60 TO 9

	111707(11						VATIVE VALUES FOR 6-VBE VALUES	
06820 DII)-AINPL/I)+V(NPI)/V(I) 06830 C 06840 C	2 2	2	06970	A(1,	1	10	07150 SUBROUTINE FUNIBE(F,D,P,X) 07160 C 07170 C-THIS ROUTINE PROVIDES FUNCTION AND DERIVATIVE VALUES FOR 07180 C-THE EVALUATION OF IER,NE,IS/80 USING 16-V8E VALUES 07200 C 07200 C	07220 DIMENSION X(1),D(1),P(1)

1) + (EXP(P(2) + X(2)) + 1 + 1 + 1 + 1 + 1 + X(2) + X(2) + 1 + 1 + X(2)
=
INE PROVIDES FUNCTION AND DERIVATIVES FOR EVALUATION ND PHTEO VIA C.JE STI. F. MOCK TRU SYSTEMS AND 1977
1/2 PERMANENT ASSIGNMENT
SION D(1),X(1),P(1) 1/(1,-X(2)/P(2))**,5)*P(3) F/P(1)
16.5
UTINE FUNCJIF, D, P, X)
INE PROVIDES FUNCTION AND DERIVATIVES FOR EVALUATION ND PHICO VIA CJC FIT. E. MOCK TRW SYSTEMS AUG.
1/3 PERMANENT ASSIGNMENT
1/((1x(2)/P(2))++.333331+P(3)
10 N
UTINE FUNCFN(F,D,P,X) M/Param/VI, Ier,Ne, IS, Betad, Rbd, Rbb, Ibb, Re, Rcc, Bf2,Vb, Fr,Vcr
SURPOUTINE PROVIDES FUNCTION AND DERIVATIVES FOR ATION OF QFN FOR COLLECTOR CURRENT FIT BY ROUTINE E. MOCK TRW SYSTEMS AUG., 1977
SION D(1),P(1),X(1)

FFX FF + (FX   1ER + NE   1S + 1BB + LNB 2   LNB 2	•	),(TV))	B2+xSqR) +xSqR)																											4.147.146.140.	
			D(1) - (-RK1+8F+EXP(X(2)/VI))/(2.+(LHB2+XSQR)+XSQR)	A C C C C C C C C C C C C C C C C C C C	ROUTINE	DIMENSION A(1)	LOGICAL BLANK	BLANK TRUE.	S-1.0	NUMB-0	TEN-1.0	•	00 10 TeleN	L.INTCHR(A)I)	IF(L.E0.36) GD TD 10	OLANK OF PALSE	. NE . 30 J	60 10 10	3	TEN=10.0	01 09	DIV-DIV+TEN	9 CONTINUE	u	TECOLANK PRETURN	F-S+FLOAT (NUMB)/DIV	PETURN	END	FUNCTION INTCHRESTRING'N	DATA SECTION SECTIONS THE THE THE THE THE THE	A PUT A DUTA DUTA DUTA DUTA DUTA DUTA DUTA

06070	ZHUZZHVZZHWZZHYZZHYZZHYZZH ZZH÷ZH÷Z	
08080	X 1H/e lH(s lH)e lHeelHeelHeelHeelHeelHeelHeelHeelHeelH	
06080	DATA EBCD/IH+.1H(.1H).1H+.1H-/	
00100	CALL GETCHR (STRINGON) CHR	
08110	IF (CHR.NE.SEQ(37)) GO TO 2	
08120	INTCHR . 36	
08130	<b>60 T0 10</b>	
08140	2 DD 1 I=1548	
08190	IF(SEQ(1).EQ.CHR) 60 10 9	
08160	1 CONTINUE	
08170	[-5]	
06160	TFEBCOLL . FO CLR . LO 30	
06190	TF(F9CD(2), Eq.CHR) 1-52	
00290	IF(EBCD(3),EQ.CHR) 1-43	
08210		
08250	TF(EBCD(5).Eq.CHR) [-47	
08230	9 INTCHR-1-1	
08240	10 RETURN	
06250	ONU	
08260	SUBROUTINE GETCHR(A,N,C)	
09270	DIMENSION A(N)	
08280	DIMENSION FMAT(6)	
08280	D4TA FMAT/4H(A1) 27H(1x2A1) 27H(2x2A1) 27H(3x2A1) 27H(4x2A1) 2	
00890	+ 7H(5x,A1)/	
01690	DEC DO E (10, FMAT (N), A) C	
08320	RETURK	
06330	END	

## APPENDIX D

TRANSIENT RADIATION RECOVERY CHARACTERISTICS OF BIPOLAR JUNCTION TRANSISTORS

## TRANSIENT RADIATION RECOVERY CHARACTERISTICS OF BIPOLAR JUNCTION TRANSISTORS

Dr. John Choma, Jr. Senior Member, IEEE

## ABSTRACT

This paper addresses the transient recovery characteristics of a bipolar junction transistor subjected to a saturating dose-rate radiation event. An estimate of the maximum recovery time is derived in terms of parameters typically listed on device specification sheets, physical device characteristics, and circuit parameters. Performance degradation due to high-injection phenomena is also addressed.

#### 1.0 INTRODUCTION

Saturation is invariably incurred by a bipolar junction transistor exposed to an ionizing dose rate event, particularly if the transistor in question conducts large quiescent collector current or if the collector of the device is terminated in a large resistance. Since a saturated transistor behaves as a short-circuit in the sense that it is incapable of responding to input signal excitation, the time required by the device to recover from the saturating excitation is clearly a critical measure of circuit performance. For example, this recovery certainly influences the frequency response of a linear system and in digital circuitry, it degrades the nominal propagation delay. If the transistor in question is in a power supply, excessive recovery requirements may well lead to catastrophic circuit or system failure.

This paper addresses the problem of generating realistic worst case estimates of recovery time in saturated transistors. The equations are cast as functions of device and circuit parameters. Although the results are predicated on simple modeling techniques, a semi-quantitative account of complex high-injection phenomena is offered to refine the first-order results obtained.

## 2.0 PHOTOCURRENT RESPONSE

Figure (1) depicts a simplified schematic diagram of a bipolar junction transistor (BJT) subjected to a gamma dose rate or gamma dot ( $\dot{\gamma}$ ) event. The gamma dot event is modeled electrically by a collector-to-base photocurrent source,  $i_{pp}(t)$ . As shown in Figure (2), photocurrent  $i_{pp}(t)$  is a simple current pulse of amplitude  $I_{pp}$  and width  $I_{p}$ .

Let it be assumed that for all steady-state conditions corresponding to  $i_{pp}(t) = 0$ , the effective source current,  $I_S$ , is of a magnitude appropriate to operation of the BJT in the linear active region of its static characteristic curves. Thus, for time  $t < T_p$ , the collector current is  $i_c(t) = I_{CA}$ . The corresponding base current,  $i_b(t)$ , is

$$I_{BA} = I_{CA}/h_{FE}, \qquad (1)$$

where h is the static forward current transfer ratio of the bipolar device.

At  $t = -T_p$ , the collector current,  $i_c(t)$ , rises toward its saturated value, say  $I_{CS}$ , in response to abrupt photocurrent excitation. By inspection of Figure (1), this saturated value of collector current is

$$I_{CS} = \frac{V_{CC} - V_{CSAT}}{R_C} - I_{PP}, \qquad (2)$$

where  $V_{\rm CSAT}$  is the collector-to-emitter saturation voltage. The corresponding base current is not linearly related to  $I_{\rm CS}$  by an expression of the form of (1) since  $h_{\rm FE}$  is defineable only for the case of transistor operation in its linear active region. Accordingly, if constant collector current in the amount of  $I_{\rm CS}$  is achieved in the time frame,  $-T_{\rm p} < t < 0$ , the base current,  $I_{\rm RS}$ , is

$$I_{BS} = I_{S} + I_{PP} - \frac{V_{BSAT}}{R_{BB}}.$$
 (3)

In (3),  $V_{\rm BSAT}$  is the value of base-emitter voltage,  $v_{\rm BE}(t)$ , under steady-state saturated conditions.

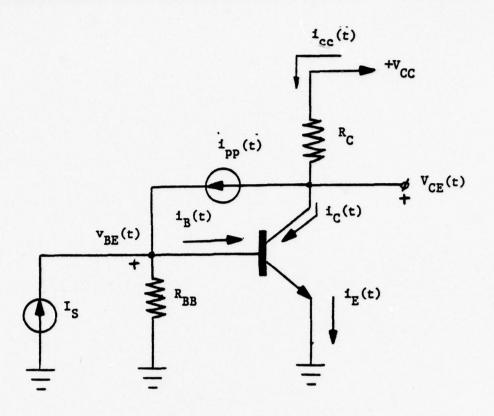


Figure (1). Schematic Diagram of Circuit Used For Radiation Storage Time Analysis

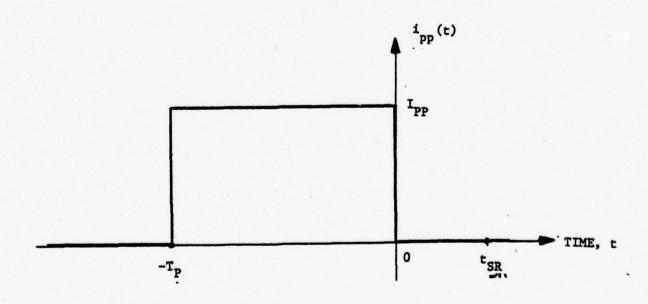


Figure (2). Time Domain Variation of Photocurrent Source, ipp(t)

Assuming that the transistor in Figure (1) is indeed saturated in the neighborhood of t < 0, both device junctions are forward biased. Resultantly, a substantial amount of charge is stored in the base. The presence of excess base charge, which is viewed herewith as the difference between total base charge and the charge that accrues if only the base-emitter junction is forward biased, precludes an immediate response of the collector current to the instantaneous termination of the photocurrent pulse. The time required to reduce the excess charge to zero or equivalently, the time required to re-establish non-forward bias across the base-collector junction, is termed the "radiation storage time," t<sub>CP</sub>.

The concept of radiation storage time is best appreciated by an investigation of Figure (3), which conceptually plots the distribution of charge injected into the base under varying bias conditions. Figure (3a) pertains to a BJT operating in its linear active region. Charge  $Q_{RO}$ , which is symbolized by the area of the cross-hatched region, is the equilibrium minority carrier distribution evidenced in the base under zero bias conditions. If the base width (W) is small, the charge profile arising out of forward injection across the base-emitter junction can be approximated by a straight line. This line has negative slope, since the reverse bias existing across the base-collector junction promotes removal of charges from base-to-collector. To the extent that carrier transport resulting from removal of base charge can be explained in terms of diffusion mechanics alone, the collector current is directly proportional to the magnitude of the slope of injected charge profile. Thus, collector current increases are in one-to-one correspondence with increasing magnitude of charge profile slope.

Under saturated conditions, the collector current is fixed at the level defined by (2), despite the fact that photocurrent flow gives rise to continuing injection of charge into the base region. Since collector current, and therefore charge profile slope, is constant, the resultant charge distribution must mirror Figure (3b). The excess charge alluded to earlier is  $Q_{\rm BS}$ , which materializes primarily from injection across a forward biased base-collector junction. In short, both junctions are

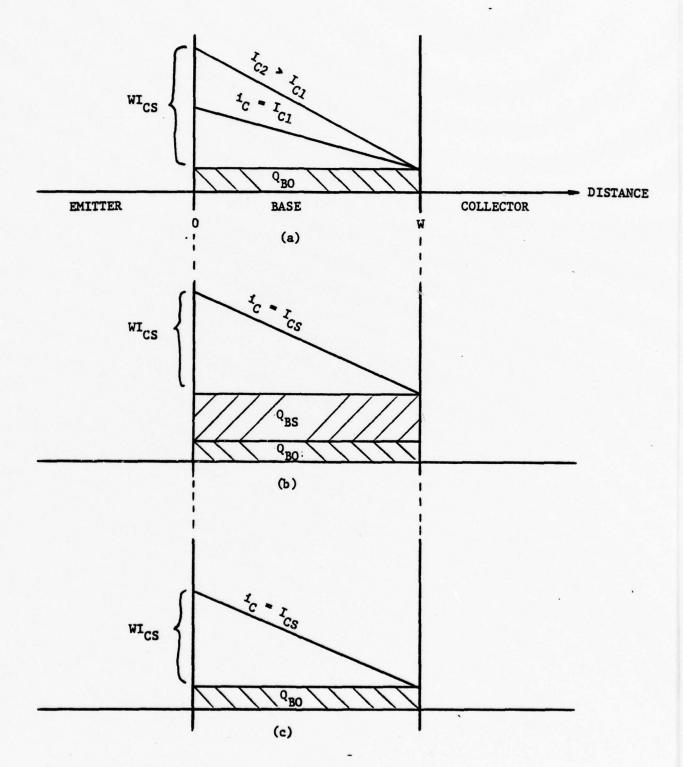


Figure (3). Charge Profiles In Base Region

- (a) Normal Active Operation
- (b) Saturated Conditions
- (C) Threshold of Saturated-Normal Active Operation

forward biased in saturated regimes, and this condition is properly reflected in Figure (3b) by the fact that the charge distribution is not in equilibrium at either junction.

Upon termination of the radiation event, the collector current remains nominally constant for a time,  $t_{SR}$ . This situation can be visualized as spacially uniform removal of  $Q_{BS}$ , as suggested in Figure (3c). When  $t_{SR}$  seconds have elapsed, the collector current commences its response to photocurrent pulse discontinuation, since further removal of charge mandates a decrease in profile slope magnitude. Note that at the instant when  $Q_{BS}$  = 0, the BJT is at the threshold of re-entering its linear active operating regime.

Figure (4) offers plots of the collector and base current responses to the photocurrent excitation defined in Figure (2). Observe that the base current decreases almost instantaneously at t = 0 to a level,  $I_{BO}$ , despite the fact that the collector current remains constant in the time interval,  $0 < t < t_{SR}$ . The magnitude of this change in base current is a function only of  $I_{pp}$  and resistive loading in the base circuit, since the voltage across the strongly charged base-emitter junction cannot change instantaneously. To the extent that most of the excess charge,  $Q_{BS}$ , is swept primarily into either emitter or collector regions, the base current level of  $I_{BO}$  is essentially a constant up to time  $t_{SR}$ . This statement reflects the tacit assumptions that (1) removal of  $Q_{BS}$  does not perturb the charge stored as a result of forward injection from the emitter and (2) the base current component attributed to a forward biased base-collector junction is much smaller than the base current attributed solely to a forward-biased base-emitter junction.

Time  $t_R$  is termed the "recovery time" of the BJT. It is conventionally defined as the time required for collector current decrease to within 5% of quiescent value, following a saturating event.

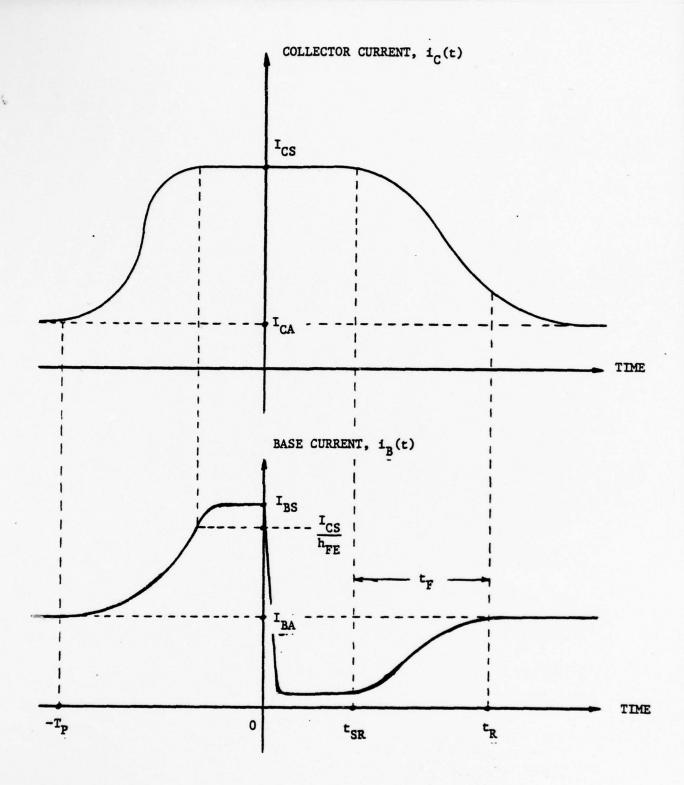


Figure (4). Collector And Base Current Responses
To Photocurrent Excitation

- (a) Collector Current Waveform
- (b) Base Current Waveform

#### 3.0 STORAGE TIME ANALYSIS

3.0

The low injection static equations which characterize BJT performance in all operating regimes are

$$I_{E} = I_{S} \left[ \epsilon^{V_{E}^{\prime}/V_{T}} - 1 \right] - I_{S} \left[ \epsilon^{V_{C}^{\prime}/V_{T}} - 1 \right] + \frac{I_{S}}{h_{FE}} \left[ \epsilon^{V_{E}^{\prime}/V_{T}} - 1 \right]$$
(3)

$$I_{C} = I_{S} \left[ \epsilon^{V_{E}^{\prime}/V_{T}} - 1 \right] - I_{S} \left[ \epsilon^{V_{C}^{\prime}/V_{T}} - 1 \right] - \frac{I_{S}}{h_{FR}} \left[ \epsilon^{V_{C}^{\prime}/V_{T}} - 1 \right], \qquad (4)$$

where  $V_E'$  and  $V_C'$  are intrinsic voltages across base-emitter and base-collector junctions, respectively,  $V_T$  is the voltage equivalent of absolute junction temperature (25.9mV at 300°K), and  $I_S$  is termed a reference transport current. The static values of emitter current  $I_E$  and collector current  $I_C$  flow in the polarity directions defined in Figure (1), and voltages  $V_E'$  and  $V_C'$  are positive when the respective junctions are forward biased. Finally, parameter  $h_{FE}$  is the forward gain exploited in (1), while  $h_{FR}$  is the static reverse current transfer ratio; i.e.,  $h_{FR}$  is the gain that materializes if the roles of emitter and collector are interchanged.

An inspection of (3) and (4) readily indicates that both emitter and collector currents are superpositions of the effects of charge injection across base-emitter and base-collector junctions. Since the terms involving  $V_E$  are dominant for forward or linear active modes of operation, while the terms in  $V_C$  predominate under inverted operational circumstances, it is reasonable to decompose  $I_E$  and  $I_C$  into forward and inverse current components. Accordingly, let

$$I_{E} \stackrel{\Delta}{=} I_{EF} - I_{EI} \tag{5}$$

$$I_{C} \stackrel{\Delta}{=} I_{CF} - I_{CI}, \tag{6}$$

where

$$I_{EF} = \frac{I_S}{\alpha_F} (\epsilon^{V_E^2/V}_T - 1), \qquad (7)$$

$$I_{EI} = I_{S} (\epsilon^{V}C^{/V}T - 1), \qquad (8)$$

$$I_{CF} = \alpha_F I_{FF}, \tag{9}$$

$$I_{CI} = I_{EI}/\alpha_{R}.$$
 (10)

In (7) through (10),

$$\alpha_{\rm F} = \frac{h_{\rm FE}}{h_{\rm FE} + 1} \tag{11}$$

$$\alpha_{R} = \frac{h_{FR}}{h_{FR} + 1} \quad (12)$$

The decomposition of (3) and (4) into the form of (5) and (6) is significant in the sense that the constituents of the BJT current can be identified in terms of the charge storage components which materialize in the base due to injection across either junction. In particular, if  $Q_{\rm BS}$  in Figure (3) can be attributed solely to charges which traverse a forward-biased ( $V_{\rm C}^{\prime}$  > 0) base-collector junction, both  $I_{\rm EI}$  and  $I_{\rm CI}$  are non-zero only if the time rate of change of  $Q_{\rm BS}$  is non-zero. It follows that upon termination of the photocurrent pulse, a solution for radiation storage time,  $t_{\rm SR}$ , embodies an investigation of the time at which  $I_{\rm EI}$  and  $I_{\rm CI}$  vanish.

Although the foregoing solution tact is meaningful, (5) through (10) are incapable of delivering time-domain response information unless they are modified to account for charge dynamics in the base. These dynamics are typically modeled by currents flowing across nonlinear junction capacitances which charge or discharge upon application or removal of the photocurrent pulse. A first order account of dynamical charge effects is accomplished by replacing constants  $\alpha_F$  and  $\alpha_R$  in (11) and (12) by the respective complex frequency transforms,

$$\alpha_{F}(s) = \frac{\alpha_{F0}}{1 + s/\omega_{F}} \tag{13}$$

$$\alpha_{\rm R}(s) = \frac{\alpha_{\rm RO}}{1 + s/\omega_{\rm F}} \quad . \tag{14}$$

In the above relationships,  $\omega_F$  is defined as the forward gain-bandwidth product, while  $\omega_R$  is the inverted analog of  $\omega_F$ . Specifically,  $\omega_F$  is the radial gain bandwidth product of forward common-emitter current gain, and  $\omega_R$  is the radial gain bandwidth product of inverse common-emitter current gain.

In the time interval,  $0 \le t \le t_{SR}$ , it may be argued that the inverse components of emitter and collector currents are expressible in the form,

$$i_{ET}(t) = i_{EI}(0) + \Delta i_{EI}(t)$$
 (15)

$$i_{CI}(t) = i_{CI}(0) + \Delta i_{CI}(t)$$
 (16)

In (15) and (16),  $i_{EI}(0)$  and  $i_{CI}(0)$  are steady-state inverse currents which exist immediately prior to termination of photocurrent excitation. On the other hand,  $\Delta i_{EI}(t)$  and  $\Delta i_{CI}(t)$  are changes incurred in the inverse emitter and collector currents during the interval of radiation storage time. Clearly, radiation storage time,  $t_{SR}$ , can be defined implicitly as the largest of the two times,  $t_{SRE}$  and  $t_{SRC}$ , such that

$$\Delta i_{EI}(t_{SRE}) = i_{EI}(0). \tag{17}$$

$$\Delta i_{CI}(t_{SRC}) = i_{CI}(0). \tag{18}$$

It is necessary to stipulate

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$$t_{SR} = MAX(t_{SRE}, t_{SRC}), \qquad (19)$$

since the transistor re-enters its active operational region only when all inverse current flow ceases.

From (5) through (10), it is easily shown that

$$I_{EI}(s) = \left[\frac{\alpha_{R}(s)}{1 - \alpha_{F}(s)\alpha_{R}(s)}\right] \left[\alpha_{F}(s)I_{E}(s) - I_{C}(s)\right]$$
(20)

$$I_{CI}(s) = \left[\frac{\alpha_{F}(s)I_{E}(s) - I_{C}(s)}{1 - \alpha_{F}(s)\alpha_{R}(s)}\right]$$
(21)

whence, by (13) and (14)

$$i_{EI}(0) \equiv I_{EI} = \left(\frac{\alpha_{RO}}{1 - \alpha_{FO}\alpha_{RO}}\right) \left(\alpha_{FO}I_E - I_C\right)$$
 (22)

$$i_{CI}(0) = I_{CI} = \frac{\alpha_{FO}I_E - I_C}{1 - \alpha_{FO}\alpha_{RO}}$$
 (23)

In (22) and (23), it is to be understood that  $I_E$  and  $I_C$  represent the steady state irradiated values of emitter and collector current, respectively.

Since the collector current is nominally constant for  $0 \le t \le t_{SR}$ , the transforms of the incurred changes in inverse currents are, from (20) and (21),

$$\Delta I_{EI}(s) = \begin{bmatrix} \alpha_{F}(s)\alpha_{R}(s) \\ 1 - \alpha_{F}(s)\alpha_{R}(s) \end{bmatrix} \Delta I_{E}(s)$$
 (24)

$$\Delta I_{CI}(s) = \Delta I_{EI}(s)/\alpha_{R}(s). \tag{25}$$

From Figures (1) and (5), the net change in emitter current during the storage time interval is

$$I_{E}(s) = \frac{1}{s} \left[ \left( I_{BO} + I_{CS} \right) - \left( I_{BS} + I_{CS} \right) \right] = -\frac{1}{s} I_{BS} - \left( I_{BS} - I_{BO} \right) \right] ,$$
 (26)

where  $I_{BO}$  is the nominal base current flowing in the time interval,  $0 < t < t_{SR}$ . By virtue of the charge storage arguments presented earlier,

$$I_{BO} = I_{BS} - \left(\frac{R_{BB}}{R_{BB} + R_{B}}\right) I_{PP}, \tag{27}$$

where  $R_B$  is the ohmic resistance of the active base region. Substitution of (13), (14), and (26) into (24) leads to the second-order transform,

$$\Delta I_{EI}(s) = -\frac{\alpha_{FO} \alpha_{RO} \omega_{F} \omega_{R} (I_{BS} - I_{BO})/s}{s^{2} + (\omega_{F} + \omega_{R})s + (1 - \alpha_{FO} \alpha_{RO}) \omega_{F} \omega_{R}}.$$
 (28)

Assuming  $\omega_F >> \omega_R^{(1)}$ , (28) reduces to

$$\Delta I_{EI}(s) = -\frac{\alpha_{FO}\alpha_{RO}\omega_{R}}{s(s+s_{o})} (I_{BS} - I_{BO}), \qquad (29)$$

where

$$s_{o} \stackrel{\Delta}{=} \frac{(1 - \alpha_{FO} \alpha_{RO}) \omega_{R}}{\omega_{R} + \omega_{F}} \simeq (1 - \alpha_{FO} \alpha_{RO}) \omega_{R}. \tag{30}$$

The inverse transform of (29) is

$$\Delta i_{EI}(t) = -\left(\frac{\alpha_{FO}\alpha_{RO}}{1 - \alpha_{FO}\alpha_{RO}}\right) (I_{BS} - I_{BO}) (1 - \epsilon^{-s} o^{t}), \qquad (31)$$

for  $0 \le t \le t_{SRE}$ . An analogous tact produces

$$\Delta i_{CI}(t) = -\left(\frac{\alpha_{FO}}{1 - \alpha_{FO}\alpha_{RO}}\right)(I_{ES} - I_{BO})(1 - \alpha_{FO}\alpha_{RO} \epsilon^{-s} o^{t})$$
 (32)

in the interval,  $0 \le t \le t_{SRC}$ , for the change in the inverse component of collector current.

<sup>(1)</sup> The inverse bandwidth is much smaller than the forward bandwidth owing to the reduced transport efficiency incurred as a result of charge injection from a lightly doped collector and charge collection by a heavily doped emitter.

An expression for  $t_{SRE}$  evolves straight-forwardly if (17), (22), and (31) are combined. A satisfying form for this expression rests on the observation that in (22),

$$\alpha_{FO}I_E - I_C = \alpha_{FO}(I_{CS} + I_{BS}) - I_{CS} = \alpha_{FO}(I_{BS} - I_{CS}/h_{FE}).$$

In particular, it can be shown that

$$t_{SRE} = \frac{1}{S_0} \ln \left( \frac{1}{1 - N_0} \right) , \qquad (33)$$

with

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$$N_o \stackrel{\Delta}{=} \frac{I_{BS} - I_{CS}/h_{FE}}{I_{BS} - I_{RO}} . \tag{34}$$

Similarly,

$$t_{SRC} = \frac{1}{S_o} \ln \left( \frac{\alpha_{FO} \alpha_{RO}}{1 - N_o} \right). \tag{35}$$

Since  $\alpha_{FO}\alpha_{RO}$  is less than unity,  $t_{SRE} > t_{SRC}$  and thus,

$$t_{SR} = t_{SRE}$$
 (36)

It might be noted in passing that the inequality,  $t_{\rm SRE}$  >  $t_{\rm SRC}$  is a reasonable result. This assertion follows from the fact that since inverse current flow reflects charge transport from collector-to-emitter, the time difference,  $t_{\rm SRE}$  -  $t_{\rm SRC}$ , represents delays associated with charge transport across the base layer of the BJT.

There is considerable engineering significance to the parameter,  $N_{\rm O}$ , defined by (34). Observe that the ratio,  $I_{\rm CS}/h_{\rm FE}$ , is the minimum base current required for BJT saturation in the steady state. This is to say that a static base current of  $I_{\rm CS}/h_{\rm FE}$  is commensurate with the onset of inverse injection across base-collector junction. A slightly larger base drive delivers  $Q_{\rm BS} > 0$ , and the BJT is forced into saturation. It follows that since  $I_{\rm BS}$  is the actual steady state base current for the saturated BJT,  $(I_{\rm BS} - I_{\rm CS}/h_{\rm FE})$  in the numerator on the right hand side of (34) represents the excess base current commensurate with assurance

of a saturated condition. Recalling that  $I_{BO}$  is the base current evidenced immediately after cessation of the photocurrent,  $(I_{BS} - I_{BO})$  in (34) can be viewed as the driving current serving to effect recovery of the BJT. Since  $I_{BO}$  must be smaller than  $I_{CS}/h_{FE}$  if the BJT is indeed to recover,  $N_{O}$  is always less than one. Accordingly,  $N_{O}$  might be termed a "recovery" factor in the sense that it compares the base current needed for device recovery to the excess base current required to ensure device saturation. As is inferred by Figure (5), minimal radiation storage time is in one-to-one correspondence with small  $N_{O}$ . In turn, small  $N_{O}$  demands a small  $R_{D}/R_{BB}$  ratio, since from (3), (27), and (34),

$$N_o = \left(1 + \frac{R_B}{R_{BB}}\right) \left(1 - \frac{V_{BSAT}/R_{BB} + I_{CS}/h_{FE} - I_S}{I_{PP}}\right)$$
 (37)

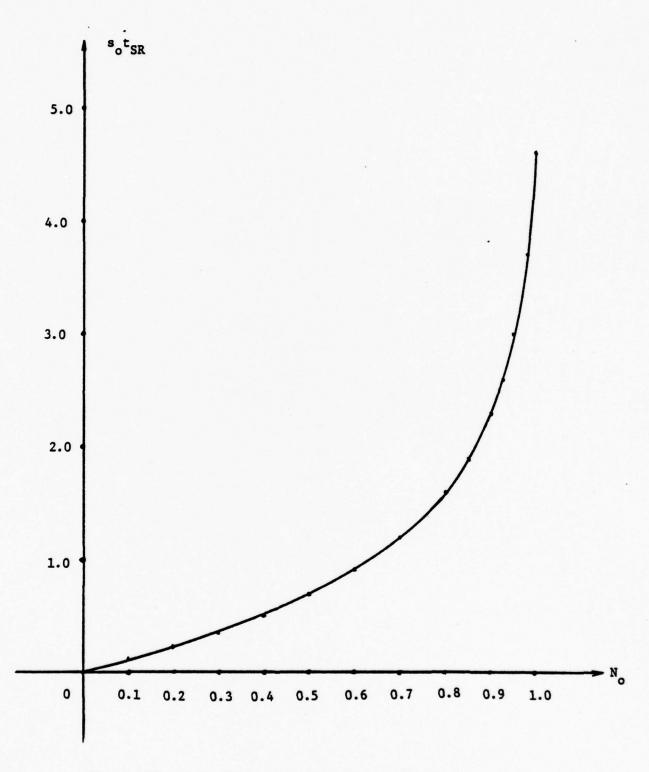


Figure (5). Normalized Radiation Storage Time As A Function of Recovery Factor,  $\mathbf{N}_{\mathbf{O}}$ 

#### 4.0 RECOVERY TIME

From Figure (4), the recovery time,  $t_R$ , is the superposition of radiation storage time,  $t_{SR}$ , and the so-called fall time,  $t_F$ . In the fall-time interval,  $t_{SR} < t < t_R$ , the BJT is in its linear active region of operation and if pole dominance can be assumed, the collector current waveform may be approximated by a single time constant expression of the form,

$$i_{C}(t) = \left[I_{CA} - (I_{CA} - I_{CS})e^{-(t - t_{SR})/T_{P}}\right] \mu(t - t_{SR}).$$
 (38)

In (38),  $\mu$ (t) is the unit step function, and  $T_p$  is the first time moment of the small-signal model used to simulate the collector current response. Taking  $i_C(t_R) = 1.05I_{CA}$ , (38) delivers,

$$t_{F} = T_{P} \left[ 3.0 + Ln \left( \frac{I_{CS}}{I_{CA}} - 1 \right) \right] . \tag{39}$$

Figure (6) depicts the small-signal model alluded to in the preceeding paragraph [1]. In this model,  $r_{\pi}$  is the diffusion resistance of the base-emitter junction,  $r_{0}$  is the incremental output resistance,  $C_{\mu}$  is the transition capacitance of the reverse-biased base-collector junction,  $C_{\pi}$  is the diffusion capacity of the base-emitter junction, and  $\beta_{0}$  is the small-signal common-emitter current gain of the BJT. The first time moment, which equates to the inverse of the dominant pole frequency, is the sum of the time constants individually attributed to  $C_{\pi}$  and to  $C_{\mu}$ , respectively. By straight-forward circuit analysis,

$$T_{p} = R_{i} \left[ C_{\pi} + \left( 1 + \frac{\beta_{o} R_{L}}{r_{\pi}} \right) C_{\mu} \right] + R_{L} C_{\mu},$$
 (40)

where

$$R_{i} = \frac{r_{\pi}(R_{B} + R_{BB})}{r_{\pi} + R_{B} + R_{BB}},$$
 (41)

$$R_{L} = \frac{r_{o}R_{C}}{r_{o} + R_{C}}.$$
 (42)

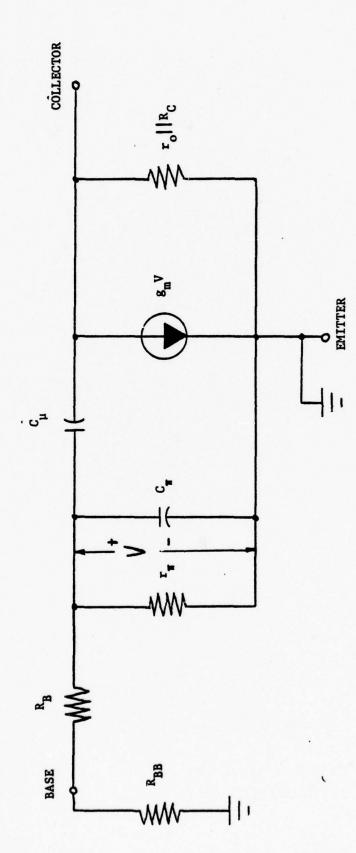


Figure (6). Small-Signal Bipolar Transistor Model Used in the Computation of Fall Time

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Assuming  $\beta_0 >> r_{\pi}/R_L$ , (40) can be cast into the form

$$T_{\mathbf{P}} \approx \left(\frac{R_{\mathbf{B}} + R_{\mathbf{B}\mathbf{B}}}{R_{\mathbf{B}} + R_{\mathbf{B}\mathbf{B}} + r_{\pi}}\right) \left(\frac{1}{\omega_{\mathbf{F}}} + R_{\mathbf{L}}C_{\mu}\right) \beta_{\mathbf{0}}, \tag{43}$$

where

$$\omega_{\mathbf{F}} = \frac{\beta_{\mathbf{O}}}{\mathbf{r}_{\pi}(\mathbf{C}_{\pi} + \mathbf{C}_{\mathbf{U}})} \tag{44}$$

is the short-circuit common-emitter gain-bandwidth product, as defined in conjunction with (13).

Care must be exercised in the utilization of (43) and (44), since these results derive from a linearized model which may inappropriately simulate the effects of large collector current excursions. In particular, the model does not incorporate the variations in small-signal parameters incurred by large changes in device currents and voltages. Accordingly, one can rationalize the tact of conservatively estimating  $t_F$  by choosing a worst case maximum value of  $T_p$ . From (43) and its companion equations, it is clear that a realistic maximum value for  $T_p$  is

$$T_{\text{PMAX}} = \beta_{\text{oMAX}} \left[ \frac{1}{\omega_{\text{FMIN}}} + R_{\text{C}}^{\text{C}}_{\text{ob}} \right], \tag{45}$$

with the understanding that  $\beta_{OMAX}$  and  $\omega_{FMIN}$  are extrema pertinent to the current excursion,  $I_{CS}^{-}$ to- $I_{CA}^{-}$ , and the corresponding collector-emitter voltage variation,  $V_{CSAT}^{-}$ to- $(V_{CC}^{-}$   $I_{CA}^{R}$ C). Furthermore,  $C_{Ob}^{-}$ , the common-base output capacitance for zero collector-base bias, is a commonly specified device parameter which exceeds, but closely approximates, small-signal capacitance  $C_{U}^{-}$ .

## 5.0 EXAMPLE CALCULATION

Assume that the circuit of Figure (1) has  $R_C = 500$  ohms,  $R_{BB} = 5,000$  ohms,  $V_{CC} = 5$  volts, and  $I_S = 200$  micro-amperes. Assume further that the BJT is characterized by the following parameters:

h<sub>FE</sub> = 50;

 $f_{TMTN} = 80MHz$ 

fp = 3MHz (inverse gain-bandwidth);

 $C_{ob} = 2.5pF;$ 

 $\beta_{\text{oMAX}} = 150;$ 

 $R_R = 800 \text{ ohms}$ 

 $h_{FR} = 0.4;$ 

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V<sub>BSAT</sub> = 2.3 volts (includes ohmic drop in R<sub>B</sub>);

V<sub>CSAT</sub> = 0.55 volts.

Finally, let the BJT be exposed to a gamma dose-rate event which gives rise to a photocurrent pulse amplitude of 2 milli-amperes.

From (2), the collector saturation current is

$$I_{CS} = \frac{5 - 0.55}{500} - (2)(10^{-3}) = 6.9 \text{mA},$$

while (3) yields for saturated base current,

$$I_{BS} = (200)(10^{-6}) + (2)(10^{-3}) - \frac{2.3}{5000} = 1.74 \text{mA}.$$

Taking  $V_{BE} = 0.78$  volt for linear active operation, Figure (1) shows that the pre-irradiated value of base current is

$$I_{BA} = I_{S} - \frac{V_{BE}}{R_{BB}} = (200)(10^{-6}) - \frac{0.78}{5000} = 44\mu A;$$

hence, from (1), the corresponding collector current is

$$I_{CA} = (50)(44) = 2.2 \text{mA}.$$

Finally, (27) produces

$$I_{BO} = (1.74)(10^{-3}) - (\frac{5000}{5800})(2)(10^{-3}) = 15.9\mu A.$$

Using (11) and (12),

$$\alpha_{\rm FO} = \frac{50}{51} = 0.98$$

and

$$\alpha_{RO} = \frac{0.4}{1.4} = 0.29.$$

Then by (30),

$$s_0 = \left[1 - (0.98)(0.29)\right] \left[2\pi(3)(10^6)\right] = 13.5 \text{MRPS}.$$

Equation (34) gives for the the recovery factor,

$$N_o = \frac{(1.74)(10^{-3}) - (6.9)(10^{-3})/50}{(1.74)(10^{-3}) - (15.9)(10^{-6})} = 0.929,$$

whence (33) and (36) deliver

$$t_{SR} = \frac{1}{(13.5)(10^6)}$$
 Ln  $\left[\frac{1}{1-0.929}\right]$  = 196nsec

for the estimated radiation storage time.

Using (21),

$$T_{\text{PMAX}} = 150 \left[ \frac{1}{2\pi (80) (10^6)} + (500) (2.5) (10^{-12}) \right] = 486 \text{nsec.}$$

Then from (39),

$$t_{\text{FMAX}} = (486)(10^{-9})\left[3.0 + \text{Ln}\,\frac{6.9}{2.2} - 1\right] = 1.83 \mu \text{sec.}$$

Thus, the estimated worst case BJT recovery time is

$$t_{RMAX} = t_{FMAX} + t_{SR} = 2.0 \mu sec.$$

Observe in the foregoing example that the storage time is in the order of 10% of the maximum recovery time. This result generally reflects situations actually encountered in realistic radiation test environments. Accordingly, it follows that BJT recovery time is largely determined by time moment parameter  $T_{\rm PMAX}$ . In turn,  $T_{\rm PMAX}$  is minimal for devices which exude low collector-base capacitance, large gain bradwidth product, and low small-signal current gain.

#### 6.0 CONCLUSIONS

An analysis of the radiation recovery transient in a bipolar junction transistor subjected to a saturating dose rate event has been developed. The analysis, which is supplemented by exemplifying calculations of recovery time, reveals device parameters and performance attributes that must be critically scrutinized if recovery time is to be constrained to a minimum. Although the general nature of sensitive device parameters is properly uncovered, it is crucial that approximations invoked during the course of analysis be understood, particularly if the primary objective of utilizing the analysis in a given application is realistic absolute estimation of recovery time.

Undoubtedly, the most significant assumption is that BJT currents flow by diffusion mechanisms alone. In simplified terminology, it has been presumed that (1) the injected charge profile in the base is linearly graded, (2) the collector current is proportional to the magnitude of profile gradient, and (3) the excess stored charge is removed prior to removal of charge accumulated from forward injection across the base-emitter junction. In reality, drift components of BJT currents are at least comparable to, if not larger than, the diffusion components in saturated devices. These drift components loom increasingly more important if the actual base current is allowed to increase significantly over the minimum base current ( $I_{CS}/h_{FE}$ ) commensurate with the threshold of saturation.

Aside from compromising the foregoing three presumptions, the immediate effect of non-negligible current transport via high-field drift mechanisms is high-current, low-voltage degradation in the gain-bandwidth product,  $\omega_F[2]$ . The maximum amount of degradation is proportional to the square of the ratio of epitaxial layer length-to-neutral base width. Depending on the type of BJT addressed, the extent of this degradation can be in the range of 15-to-20. Similar attenuations are evidenced with respect to both the static and small-signal current gains, although the factor by which these gains degrade is not as severe as the degradations quoted for gain-bandwidth product. The upshot of the matter is that if manufacturer's gain-bandwidth product and gain specifications

are used to compute T<sub>PMAX</sub>, it is conceivable that the estimated maximum recovery time might be low by factors in the range of 5-to-10, since device specifications are invariably applicable to operational regions other than high-injection regimes.

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It must also be remembered that (38) and (39) presume device linearity and pole dominance for  $t_{\rm SR}$  < t <  $t_{\rm R}$ . The effects of inherent device nonlinearities are especially pronounced if the saturated-to-quiescent collector current ratio,  $I_{\rm CS}/I_{\rm CA}$ , is larger than four or five. Ringing in the collector and base current responses is indicative of either circuit nonlinearity or a non-dominant pole population. Indeed, if non-linearity is particularly severe, the analytical concept of a pole cannot be exploited in the computation of fall time  $t_{\rm F}$ . It can be envisaged that the impropriety of either linearity or pole dominance is conducive to recovery time estimates that are low by 20%-to-40%.

In summary, it is suggested that the foregoing equations are useful in ascertaining device characteristics which strongly influence the radiation recovery response of a BJT. They are also useful when applied to the problem of bracketing a realistic estimate of maximum recovery time. The lowest anticipated maximum recovery time derives from direct employment of manufacturer's nominal specifications for key device parameters. The largest estimate for maximal recovery time is obtained by allowing gain and gain-bandwidth product to degrade in accordance with specified, or estimated high-injection device performance characteristics.

## 7.0 REFERENCES

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## APPENDIX E

CIRCUIT MODEL FOR NEUTRON-INDUCED
PERFORMANCE DEGRADATION IN BIPOLAR JUNCTION TRANSISTORS

# CIRCUIT MODEL FOR NEUTRON-INDUCED PERFORMANCE DEGRADATION IN BIPOLAR JUNCTION TRANSISTORS

J. Choma, Jr., Senior Member
IEEE

J. M. Oberholtzer

### ABSTRACT:

This paper proposes the use of an existing large-signal circuits model to simulate bipolar transistor performance degradation induced by neutron irradiation. It is shown that only three radiation parameters are required to establish a model for neutron effects over the entire practical current and voltage ranges of bipolar transistor operation.

#### 1.0 INTRODUCTION

The predominant effects of neutron displacement damage in silicon are an increase in carrier generation rate and a decrease in majority carrier concentration. In a bipolar junction transistor (BJT), these effects combine to produce current gain attenuation, degradation in gain-bandwidth product, enhanced base-collector leakage current, increased collector breakdown voltage, and increased collector saturation voltage [1]. Additionally, neutron displacement damage in a BJT may significantly perturb the nominal thermal characteristics of a bipolar circuit. Depending on circuit configuration, the nominal thermal behavior can either be aggravated or improved by neutron bombardment [2]-[3].

Traditionally, the performance of bipolar circuits immersed in a neutron environment is simulated by exploiting various damage constants which are enumerated from experimental device characterizations. In particular, gain, leakage current, and gain-bandwidth product are measured prior to and immediately following controlled neutron irradiation of a given bipolar device. This data is used to extrapolate simple relationships between pre-irradiated and irradiated values of various device parameters. An example of such a relationship is the Messenger-Spratt equation [4],

$$\frac{1}{\beta_{\phi}} = \frac{1}{\beta_{o}} + K_{\phi} \tau_{FO} \Phi, \qquad (1)$$

where  $\beta_{\varphi}$  is the irradiated value of static short-circuit current transfer ratio,  $\beta_{0}$ ,  $\tau_{FO}$  is the transit time of minority carriers across the neutral base,  $\Phi$  is neutron fluence, and  $K_{\varphi}$  is the pertinent damage constant. Other device performance indices, such as gain-bandwidth product, breakdown voltage, and leakage current, are related to device neutron fluence exposure by expressions whose forms are analogous to (1).

Although neutron-induced damage assessment equations are useful both for simulating the dependence of irradiated electrical parameters on their pre-irradiated counterparts and for establishing a basis of neutron hardness screening, at least two fundamental objections can be raised if these equations are used exclusively in evaluating bipolar response to neutron irradiation.

First, the damage constant methodology can only simulate measured results. It cannot be applied to the problem of predicting the radiation response of bipolar transistors, because the damage constants are empirical performance barometers. These constants are not well-defined in terms of intrinsic physical or geometrical device properties but instead, they derive purely from experimental data.

A second shortcoming is that most of the damage constants invoked in a neutron radiation response exercise are sensitive functions of device injection level. For example,  $K_{\phi}$  in (1) varies with the voltage and current at which gain  $\beta_0$  is measured. The upshot of the matter is that most damage assessment equations are applicable only to the injection regime in which they are formulated. Attempts made to invoke these same equations to damage assessment at significantly lower or higher injection levels generally result in substantial analytical errors. It follows that the methodology in question conveys little, if any, understanding of the sensitivity of neutron radiation response with respect to the possibly wide injection range which a bipolar device encounters in a prescribed circuit application.

This paper presents a new bipolar neutron damage model which is uniformly applicable to the low-through-high injection range of BJT operation. The model is an extension of an existing circuits model which is capable of predicting large-signal transient responses to forcing inputs that can drive the transistor from low-to-high injection regimes. In addition to the model parameters pertinent to simulation and prediction of traditional electrical responses, only three additional parameters are required for an evaluation of the effects of neutron radiation. These parameters are independent of voltage and current and moreover, they can be related to damage constants which have been enumerated from a damage assessment appropriate to mid-injection operational regimes.

# 2.0 FUNDAMENTAL TRANSISTOR MODEL

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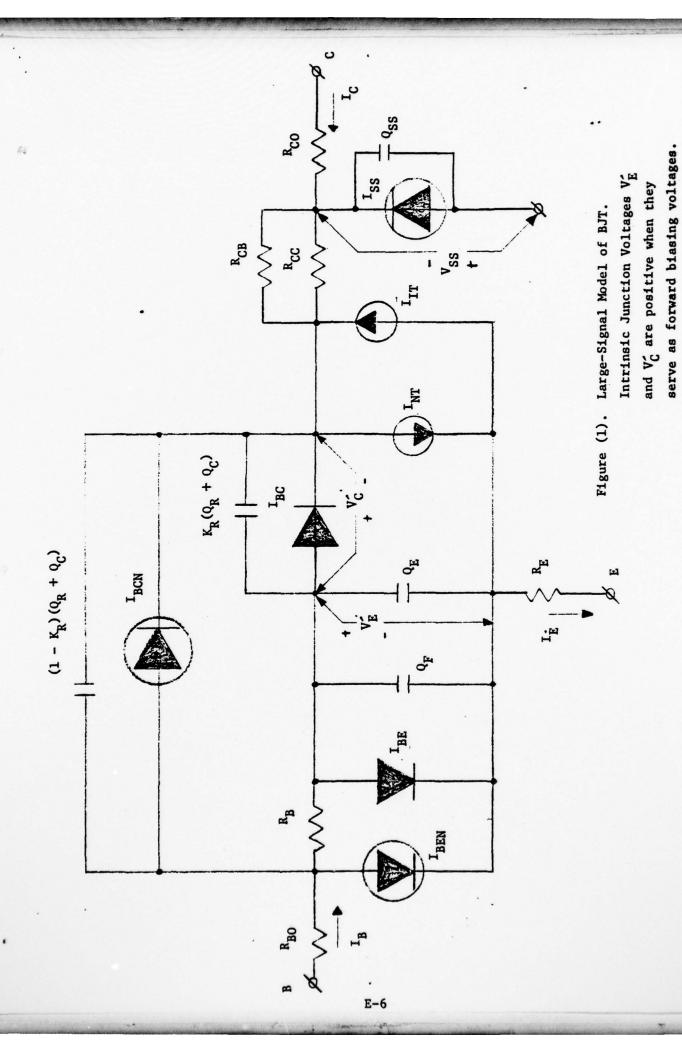
The bipolar transistor circuit model diagrammed in Figure (1) is proposed as the foundation for neutron radiation response studies. (1) The model possesses all of the attributes of the classical Ebers-Moll [5] and Gummel-Poon [6] topologies and in addition, it provides the following simulation capabilities.

- (1). The effects of base-collector depletion layer width modulations arising out of collector junction voltage changes are incorporated. This phenomenon, known as the Early Effect [7], gives rise to current gain dependence on collector voltage. Analogously, an account is made of Inverse Early Effect, which produces inverse beta dependence on baseemitter voltage.
- (2). Low-injection forward current gain attenuation due to non-ideal diode currents flowing across the base-emitter junction is replicated by the model. In Figure (1), the lumped equivalent to the non-ideal baseemitter diode current is

$$I_{EFN} = I_{ER} (\epsilon^{V_E^{\prime}/N} E^{V} TE - 1), \qquad (2)$$

where  $V_{TE}$  is junction thermal voltage, and  $I_{ER}$  and  $N_{E}$  relate implicitly to emitter efficiency, degree of emitter sidewall injection, and the recombination velocity at the base-emitter interface. Parameter  $N_{E}$  is always greater than or equal to unity. The diode current,  $I_{BCN}$ , mathematically parallels the form of (2) and allows for low-level inverse current gain attenuation for transistors operated in inverted or saturated modes.

<sup>(1)</sup> The model, which was developed by the author at the University of Pennsylvania, is embedded in the ISPICE computer-aided design program and is available on a timeshare network marketed by National CSS, Inc., Norwalk, Connecticut.



4. 4.4.

(3). High-injection forward current gain degradation caused by conductivity modulation in both the base and quasi-saturated collector regions is incorporated in the model of Figure (1). The element germane to this simulation capability is

$$I_{NT} = \left(\frac{Q_{BO}}{Q_B}\right) I_{S}(\epsilon^{V_E^2/V}TE - 1), \qquad (3)$$

where  $Q_B$ , the net majority charge in the effective base, has a value of  $Q_{BO}$  under zero bias ( $V_E^* = V_C^* = 0$ ) conditions. The charge function,  $Q_B$ , which is the superposition of charge  $Q_E$  stored in the base-emitter transition layer, charge  $Q_C$  stored in the base-collector depletion region, and the charges ( $Q_F$  and  $Q_R$ ) required to preserve space-charge neutrality in the face of charge injection across either junction, is given by

$$\frac{Q_B}{Q_{BO}} = (1 + \lambda_E) + \frac{\tau_F I_S (\epsilon^{V_E/V} TE - 1)}{Q_B} + \frac{\tau_R I_S (\epsilon^{V_C/V} TC - 1)}{Q_B}$$
 (4)

In (4),

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$$\lambda_{E} = \frac{V_{E}^{\prime}}{V_{ER}} + \frac{V_{C}^{\prime}}{V_{CR}} \quad , \tag{5}$$

where  $V_{ER}$  and  $V_{CR}$  are inverse and forward Early effect parameters, and  $\tau_F$  ( $\tau_R$ ) is the forward (inverse) transit time of minority carriers in the field-neutral base. Because of base pushout phenomena, these transit times vary with device current and voltage [8]. Parameter  $I_S$  in (3) and (4) is directly proportional to the product of intrinsic carrier concentration and carrier diffusivity in the base, and it relates inversely to the total majority charge in the base region.

(4). High-injection gain-bandwidth product ( $f_T$ ) degradation caused by base pushout and conductivity modulation in the collector layer is a model feature. The model also addresses the related nonlinearity in effective collector resistance, ( $R_{CC} + R_{CO}$ ).

- (5). The effects of static emitter crowding on the effective base resistance,  $(R_{BO} + R_{BB})$ , are embodied in the model. Excess phase and both static and dynamic conduction across the collector substrate interface are also included.
- (6). A first order approximation of the distributed effects of high-frequency excitation is introduced by way of the semi-empirical parameter,  $K_R$  ( $K_R \le 1$ ).

#### 3.0 LOW-INJECTION CURRENT GAIN DEGRADATION

Under static forward bias operating conditions, all charge functions vanish, the substrate diode current,  $I_{SS}$ , is very nearly zero, and currents  $I_{BCN}$ ,  $I_{BC}$ , and  $I_{IT}$ , which are exponential functions of the intrinsic base-collector junction voltage, are essentially zero. Accordingly, Figure (1) verifies that base current  $I_{R}$  is given by

$$I_{B} = I_{BE} + I_{BEN}, \tag{6}$$

where IREN is defined by (2) and

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$$I_{BE} = \frac{I_{S}}{\beta_{O}} \left[ \epsilon^{V_{E}^{\prime}/V_{TE}} - 1 \right]. \tag{7}$$

In (7),  $\beta_0$  is the maximum value of the static transfer ratio,

$$h_{FE} = \frac{I_C}{I_R} , \qquad (8)$$

under the special condition of zero bias  $(V_C^* = 0)$  across base-collector junction.

Returning to Figure (1), it is clear that under the conditions stipulated above, the collector current,  $I_{\mathbb{C}}$ , is identical to the current,  $I_{\mathbb{NT}}$ . From (3) and (4) it can be shown that

 $I_{C} = \frac{I_{S}(\varepsilon^{V_{E}^{*}/V_{TE}-1)}}{\left(\frac{1+\lambda_{E}}{2}\right) + \left\{\left(\frac{1+\lambda_{E}}{2}\right)^{2} + Q_{FN}(\varepsilon^{V_{E}^{*}/V_{TE}-1}) + Q_{RN}(\varepsilon^{V_{C}^{*}/V_{TC}-1})\right\}^{1/2}}, \quad (9)$ 

where

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$$Q_{FN} \stackrel{\Delta}{=} \frac{\tau_{F}^{I} s}{Q_{BO}}$$
 (10)

$$Q_{RN} \stackrel{\Delta}{=} \frac{\tau_{R}^{I} S}{Q_{BO}}$$
 (11)

The term involving  $Q_{\rm RN}$  in (9) is essentially zero, provided that the BJT is operated in its linear active domain. Moreover, for low-to-moderate injection levels, the term containing  $Q_{\rm FN}$  is negligible, so that (9) collapses to

$$I_{C} = \left(\frac{I_{S}}{1 + \lambda_{E}}\right) \epsilon^{V_{E}^{*}/V_{TE}}.$$
 (12)

In (12), it is assumed that the exponential term is much larger than unity. If this same reasonable assumption is invoked on (2), (8) can be shown to deliver

$$h_{EE} \simeq \frac{\beta_{o}/(1+\lambda_{E})}{1+EXP\left[\left(\frac{N_{E}-1}{N_{E}}\right)\left(\frac{V_{K}-V_{E}^{*}}{V_{TE}}\right)\right]},$$
(13)

with

$$V_{K} \stackrel{\Delta}{=} \left(\frac{N_{E}}{N_{E}-1}\right) V_{TE} Ln \left(\frac{\beta_{o} I_{ER}}{I_{S}}\right) . \tag{14}$$

Figure (2) portrays the dependence of  $h_{FE}$  and  $V_E^*$ , and hence the dependence of  $h_{FE}$  on the natural logarithm of collector current  $I_C$ . Recalling (5), the effect of parameter  $\lambda_E$ , as depicted in the figure, is to simulate the sensitivity of current gain with respect to the reverse biasing voltage applied at the collector. Observe then that for increasing  $V_E^*$ ,  $h_{FE}$  correctly approaches a constant at fixed collector junction voltage, provided that the contribution of  $V_E^*$  to  $\lambda_E$  is insignificant in comparison to the contribution of  $V_C^*$ . For decreasing forward bias,  $h_{FE}$  attenuates in such a way that at  $V_E^* = V_K$ ,  $h_{FE}$  is one-half of the current gain at moderate current levels. Notice that  $V_K$  is dependent on both  $I_{ER}$  and  $N_E$ , which implicitly relate to the characteristics of the base-emitter interface. Finally,  $N_E$  fixes the slope of the gain curve at  $V_E^* = V_K$ . If  $N_E = 1$ , which infers only idealized diode current flow across the base-emitter junction,  $h_{FE}$  is constant over the addressed range of  $V_E^*$ .

The preceding discussion applies to a BJT operated in a traditional, non-nuclear, electrical environment. When this environment is perturbed by significant neutron fluence, one can postulate that additional base current, say  $\Delta I_B$ , must be supplied in order to maintain the collector current at its pre-irradiated level. Let this base current enhancement be expressed as

1. 1.3.

$$\Delta I_{B} = I_{\phi} \left[ \epsilon^{\nabla_{E}^{2}/V} TE - 1 \right] + I_{\phi \circ} \left[ \epsilon^{\nabla_{E}^{2}/N} E^{V} TE - 1 \right] , \qquad (15)$$

where it is understood that  $I_{\phi}$  is a function of the neutron energy imparted to the active region of the base-emitter junction. On the other hand,  $I_{\phi\,o}$  is parametrically related to the neutron energy absorbed by the emitter sidewalls, the emitter-base depletion region, and the surface of the depleted layer associated with the base-emitter junction.

Clearly, the induced increment in base current can be modeled by incorporating two additional current sources in the base-emitter circuit of Figure (1). Current  $I_{\rm BEN}$  is placed in parallel with a source whose current is defined by the second term on the right hand side of (15), while the first term in this relationship is a diode in shunt with current  $I_{\rm BE}$  in Figure (1). Notice, however, that these modifications infer a mere redefinition of parameters  $\boldsymbol{\beta}_{\rm O}$  and  $I_{\rm ER}$ . Specifically,  $\boldsymbol{\beta}_{\rm O}$  and  $I_{\rm ER}$  in the traditional model can be replaced by irradiated parameter values,  $\boldsymbol{\beta}_{\rm O}$  and

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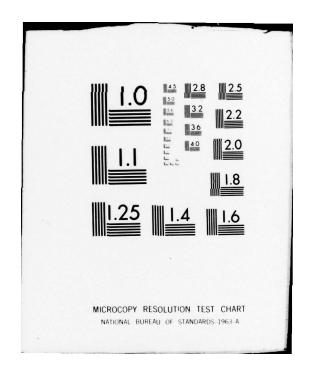








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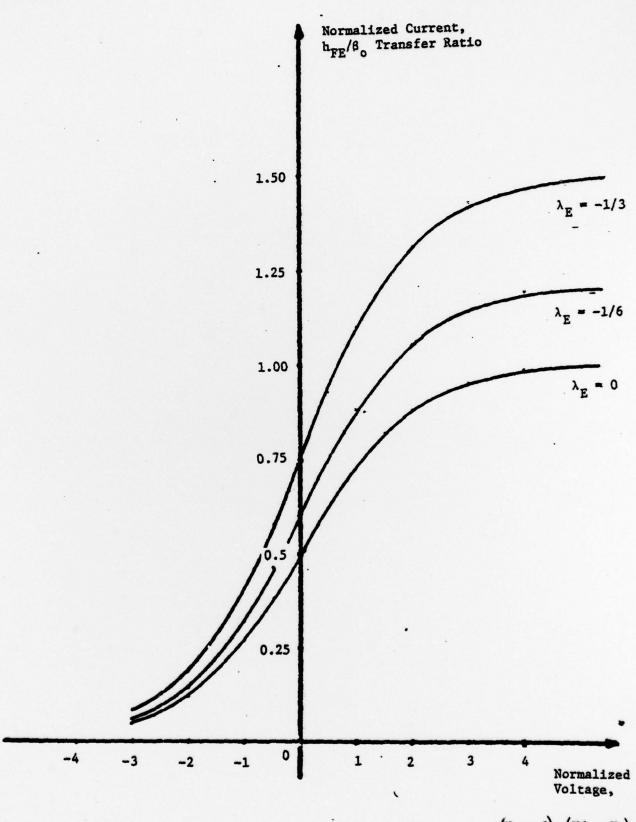


Figure (2). Forward Large-Signal, Static Current Transfer Characteristic Predicted by Model of Figure (1).

 $\left(\frac{N_E-1}{N_E}\right)\left(\frac{V_E^*-V_K}{V_T}\right)$ 

I oR, respectively, such that

$$\beta_{\phi} = \frac{\beta_{o}}{1 + \beta_{o} I_{\phi} / I_{S}} \tag{16}$$

$$I_{\phi R} = I_{ER} + I_{\phi Q}. \tag{17}$$

A number of noteworthy comments can be offered at this juncture. First, only two constant parameters,  $I_{\phi}$  and  $I_{\phi 0}$ , are required to establish neutron-induced changes in the current gain response over the range of very low-to-moderate injection levels. In particular, parameter  $I_{\phi}$  serves to attenuate the maximum current transfer ratio at moderate current levels, while  $I_{\phi 0}$  influences the nature of  $h_{FE}$  in low injection regimes. The latter point becomes clear when one replaces  $I_{ER}$  and  $\beta_0$  in (14) by  $I_{\phi R}$  and  $\beta_{\phi}$  to obtain

$$V_{K\phi} = V_K + \left(\frac{N_E}{N_E - 1}\right) V_{TE} Ln \left[\frac{1 + I_{\phi o}/I_{ER}}{1 + \beta_o I_{\phi}/I_S}\right]$$
 (18)

Thus, the base-emitter junction bias commensurate with an h\_FE that is one-half of the maximum current transfer ratio at moderate current levels shifts as a function of the radiation parameters,  $I_{\phi}$  and  $I_{\phi o}$ . Note, however, that while the maximum current transfer ratio always decreases in response to neutron irradiation, the half-gain junction bias level does not necessarily change, since if  $I_{\phi o}/I_{ER} = \beta_o I_{\phi}/I_S$ ,  $V_{K\phi} \equiv V_K$ . In effect, the voltage change,  $V_{K\phi} - V_K$ , monitors the relative deposition of energy between the active and inactive portions of the base-emitter junction.

A second interesting point is that (16) can be written as

$$\frac{1}{\beta_{\phi}} = \frac{1}{\beta_{o}} + \frac{I_{\phi}}{I_{S}} \quad , \tag{19}$$

which is precisely of the form of the classical Messenger-Spratt relationship. Indeed, a comparison of (1) and (19) infers the  $\mathbf{I}_{\phi}$  relates to neutron damage constant  $\mathbf{K}_{\phi}$  and fluence  $\Phi$  in accordance with

$$I_{\phi} = K_{\phi} \tau_{FO} I_{S}^{\phi}. \tag{20}$$

It is critical to note that since  $I_{\phi}$  is a constant model parameter,  $\tau_{FO}$  is necessarily a constant. This is to say that  $\tau_{FO}$  is the minority carrier transit time across a field-neutral base. Since a field-neutral base is ensured at moderate currents only under the condition of a strongly backbiased collector-base junction, and since  $\tau_{FO}$  is inversely proportional to common-emitter gain bandwidth product  $f_{T}$ , (20) is meaningfully expressed as

$$I_{\phi} = \left(\frac{K_{\phi}}{2\pi f_{\text{TMAX}}}\right) I_{S}^{\phi}, \qquad (21)$$

where  $f_{\underline{TMAX}}$  is understood to be the value of  $f_{\underline{T}}$  appropriate to linear active operation of a BJT subjected to strong collector-base back bias.

### 4.0 HIGH-INJECTION CHARACTERISTICS

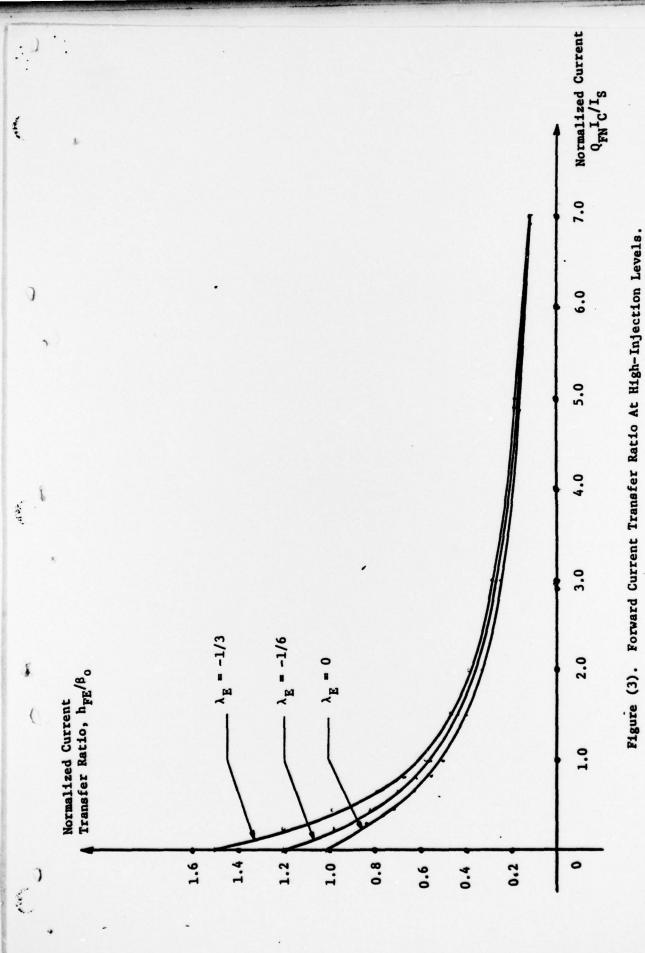
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When the base-emitter junction is strongly forward biased, the term involving  $Q_{FN}$  in (9) cannot be ignored. Moreover, since  $N_E$  is larger than unity,  $I_{BE}$  is much greater than  $I_{BEN}$  and thus, the net base current is closely approximated by the right hand side of (7). It follows that for  $V_C^{\prime}$  << 0, (7) and (9) combine to produce

$$\frac{h_{FE}}{\beta_o} \simeq \frac{1}{(1+\lambda_E) + \frac{Q_{FN}I_C}{I_c}}.$$
 (22)

Equation (22) is plotted in Figure (3), wherein it is understood that zero collector current ( $I_C = 0$ ) is in one-to-one correspondence with the moderate current levels associated with regions on the far right hand side in the graph of Figure (2). Observe in (22) that the high-injection current transfer ratio attenuates to one-half of its moderate current value at a collector current, say  $I_K$ , given by

$$I_{K} = \frac{(1 + \lambda_{E})Q_{BO}}{\tau_{F}} . \qquad (23)$$



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In arriving at (23), (10) is utilized to express  $I_K$  in terms of base region transit time,  $\tau_F$ . It should be pointed out that the actual dependence of  $h_{FE}$  on collector current is more complicated than suggested by (22) and (23), since base pushout at high-injection levels incurs a dependence of  $\tau_F$  (and hence,  $Q_{FN}$ ) on current  $I_C$  and the voltage parameter,  $\lambda_E$ .

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The development of a model to simulate neutron-induced changes in the high-injection current transfer characteristics derives from two observations. First, parameter  $\beta_0$  attenuates to  $\beta_{\dot{\varphi}}$ , as defined by (16). Thus,  $\beta_0$  in (22) is replaced by  $\beta_{\dot{\varphi}}$ , thereby incurring a uniform downward shift in the  $h_{FE}$  curves portrayed in Figure (3). Second, the majority carrier concentration in the base decreases from  $Q_{BO}$  to  $Q_{B\dot{\varphi}}$  in accordance with the semi-empirical relationship [9]

$$Q_{B\phi} = \frac{Q_{BO}}{1 + K_B \Phi} , \qquad (24)$$

where  $K_B$  can be thought of as charge damage constant which reflects neutron-induced decreases in majority carrier concentration. Upon replacement of  $Q_{BO}$  in (10) by  $Q_{B\phi}$ , one concludes that the irradiated value of  $Q_{FN}$  is

$$Q_{FN\phi} = Q_{FN}(1 + K_B\phi), \qquad (25)$$

which directly supplants pre-irradiated parameter  $Q_{\overline{FN}}$  in (22). Note, however, that such a replacement of variables is equivalent to allowing the transit time to assume an irradiated value,  $\tau_{F\phi}$ , of

$$\tau_{F\phi} = \tau_{F}(1 + K_{B}\Phi).$$
 (26)

There are at least two important ramifications of (26). First and foremost, the equation allows for the simulation of neutron-induced perturbations in the gain-bandwidth product, as well as neutron-induced changes in the forward current transfer ratio. This statement exploits the simple fact that for collector currents in excess of a critical value, say  $I_0(V_C^*)$ ,

$$f_{T} = \frac{1}{2\pi\tau_{F}}, I_{C} > I_{O}.$$
 (27)

In (27),  $\tau_F$  is a function of both collector current  $I_C$  and collector junction bias voltage,  $V_C^*$ . This functional relationship, which is portrayed graphically by Figure (4), is mathematically defined by a so-called "base-pushout function,"  $B_F(I_C, V_C^*)$ , such that

$$\tau_{\rm F} = \tau_{\rm FO} B_{\rm F} (I_{\rm C}, V_{\rm C}) = \frac{B_{\rm F} (I_{\rm C}, V_{\rm C})}{2\pi f_{\rm TMAX}}$$
 (28)

The current parameter,  $I_0(V_C)$  which defines the onset of BJT entry into high-injection operating regimes, and the pushout function,  $B_F(I_C, V_C)$ , can be related to such device processing characteristics as epitaxial layer doping, base and epitaxial layer geometries, minority carrier diffusivities, and the like [8], [10]. Alternatively, the behavior of  $\tau_F$  in high-injection domains can be modeled by empirical equations whose parameters derive conveniently from measured  $f_T$ -characteristics [11]. The upshot of the matter is that (26) allows for the simulation of  $f_T$  over both moderate and high injection regimes since if  $\tau_F$  in (27) is replaced by  $\tau_{F\phi}$ , (26) through (28) infer an irradiated value of  $f_T$  that is given by

$$f_{T\phi} = \frac{1}{2\pi\tau_{FO}(1 + K_B\phi)B_F(I_C, V_C)} = \frac{f_T}{1 + K_B\phi}, \qquad (29)$$

for  $I_C > I_O(V_C)$ .

It is useful to interject that for  $I_C \leq I_O(V_C^*)$ , (26) remains applicable to the problem of simulating gain-bandwidth product characteristics. For collector currents not exceeding  $I_O(V_C^*)$ , the base pushout function is unity and

$$f_T = \frac{1}{2\pi [\tau_{FO} + v_{TE} c_T / I_C]},$$
 (30)

where  $C_T$ , the sum of the transition capacitances of both junctions, is nominally insensitive to neutron fluence. Clearly, the neutron irradiated value of low-current  $f_T$  is obtained by replacing  $\tau_{F0}$  in (30) by  $\tau_{F\phi}$ , as defined by (26). Unlike the uniform neutron-incurred degradation of  $f_T$  at high collector currents, the effect of neutrons on low-current  $f_T$  is seen to be non-uniform, owing to the relatively small influence which neutron fluence exerts on transition layer storage dynamics.

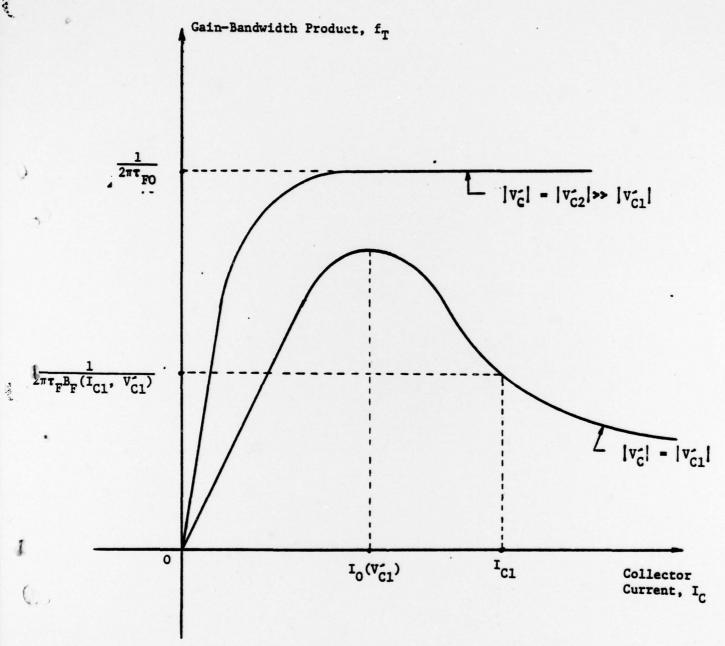


Figure (4). Common-Emitter Gain-Bandwidth Product as a Function of Voltage and Current.

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A second ramification of (26) derives directly from the observation that this equation is predicated on the fundamental assumption that the irradiated value of base region majority charge is linearly related to its pre-irradiated counterpart by an expression of the form of (24). Since both  $Q_{RN}$  in (11) and  $Q_{FN}$  in (10) are inversely proportional to  $Q_{RN}$ , it is clear that the irradiated inverse transit time behaves in a manner that is identical to the forward transit time behavior prescribed by (26).

### 5.0 CONCLUSIONS

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The fundamental premise of this paper is that an existing and reasonably accessible electrical model of a bipolar junction transistor is easily modified to incorporate the cognate effects of BJT exposure to a given neutron fluence, . The foundation for the required modifications is the fact that (1) additional base current must be supplied if the collector current of an irradiated BJT is to be maintained at its pre-irradiated value, and (2) the majority charge in the active base decreases in proportion to the neutron fluence. The additional base current and charge perturbation are stipulated by (15) and (24), respectively, which in turn define three neutron-related modeling parameters; namely, I, I, and KB.

Parameter  $I_{\phi}$  can be determined in either of two ways. One way exploits (19), which entails a comparison of the pre-irradiated and post-irradiated values of the gain parameter  $\beta_{\phi}$ , as defined in Figure (2). A second technique is to relate  $I_{\phi}$  to a known current gain damage constant in accordance with (20) or (21). Parameter  $I_{\phi \phi}$  requires an investigation of current transfer ratio evidenced at low current levels. In particular, the base-emitter voltage (which closely approximates the base-emitter junction bias at low currents) is measured prior and subsequent to neutron irradiation to deduce  $I_{\phi \phi}$  by way of (18). Finally,  $K_B$  is most easily estimated by monitoring pre-irradiated and irradiated  $f_T$  at fixed current and voltage appropriate to operation in high-injection regimes. The equation of interest is (29).

It is important to note that while parameters  $I_{\phi}$ ,  $I_{\phi o}$ , and  $K_B$  are constants, their correct utilization enables the simulation of the damaging effects of neutron bombardment over all forward active regimes of BJT operation. Indeed, parameter  $K_B$  is ostensibly applicable to gain-bandwidth

product studies of bipolar devices operated in inverted or saturated regimes. Although inverse BJT operation is not specifically addressed in this work, parameters analogous to  $\mathbf{I}_{\phi}$  and  $\mathbf{I}_{\phi o}$  can presumably be found to simulate inverted current transfer ratio characteristics. The only problem inherent in this presumption is the difficulty encountered in monitoring pre-irradiated inverted performance without incurring permanent damage at the base-emitter interface.

It is believed that the theory of neutron effects modeling developed in this paper is applicable to a broad range of BJT simulation and prediction problems. For example, over a wide range of practical bipolar device operation, the dependence of collector current on temperature at fixed bias voltage is given by [12]

$$I_C = AT^M \epsilon^{-V} GO^{/V} TE, \qquad (31)$$

where V<sub>GO</sub> is the bandgap voltage and constants A and M are determined by device fabrication processes. Clearly, M is critical to the determination of the thermal sensitivity of collector current. It turns out that published data related to the dependence of minority carrier mobility on base region impurity concentration infers [2] that M behaves approximately as

$$M \simeq Ln(Q_{BO}/Q_{MO}) - M_{O}, \qquad (32)$$

where  $Q_{MO}$  and  $M_O$  are constants. The irradiated value of M evolves from replacement of  $Q_{BO}$  by  $Q_{B\phi}$ , whence by (31), the neutron-induced thermal sensitivity of collector current is easily discerned.

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APPENDIX F

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